

Soft-Collinear Effective Theory

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T. Becher, A. Broggio, A. Ferroglia,
Introduction to Soft-Collinear Effective Theory,
arXiv:1410.1892

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The Soft-Collinear Effective Theory,
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Problems with widely separated scales

- ▶ Method of regions
- ▶ Effective field theories

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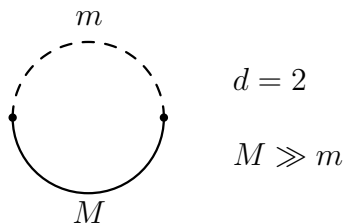
Problems with widely separated scales

- ▶ Method of regions
- ▶ Effective field theories

Plan

- ▶ Scalar φ^3
- ▶ QCD
- ▶ Some applications

Method of regions



$$\begin{aligned} I &= \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{(M^2 - k^2 - i0)(m^2 - k^2 - i0)} \\ &= \int \frac{d^d k}{\pi^{d/2}} \frac{1}{(k^2 + M^2)(k^2 + m^2)} \end{aligned}$$

Exact solution

$$d = 2 - 2\varepsilon$$

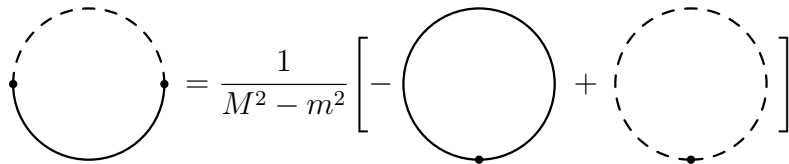
$$I = \frac{1}{M^2 - m^2} \int \frac{d^d k}{\pi^{d/2}} \left[-\frac{1}{k^2 + M^2} + \frac{1}{k^2 + m^2} \right]$$

$$\text{Diagram} = \frac{1}{M^2 - m^2} \left[-\text{Diagram}_1 + \text{Diagram}_2 \right]$$

Exact solution

$$d = 2 - 2\varepsilon$$

$$I = \frac{1}{M^2 - m^2} \int \frac{d^d k}{\pi^{d/2}} \left[-\frac{1}{k^2 + M^2} + \frac{1}{k^2 + m^2} \right]$$


$$= \frac{1}{M^2 - m^2} \left[-\text{Solid Circle} + \text{Dashed Circle} \right]$$

$$\begin{aligned} I &= -\Gamma(\varepsilon) \frac{M^{-2\varepsilon} - m^{-2\varepsilon}}{M^2 - m^2} \rightarrow \frac{\log \frac{M^2}{m^2}}{M^2 - m^2} \\ &= \frac{1}{M^2} \log \frac{M^2}{m^2} \left[1 + \frac{m^2}{M^2} + \frac{m^4}{M^4} + \dots \right] \end{aligned}$$

Method of regions

$$I = I_h + I_s$$

hard $k \sim M$

soft $k \sim m$

Hard region

$$k \sim M$$

$$I_h = \int \frac{d^d k}{\pi^{d/2}} T_h \frac{1}{(k^2 + M^2)(k^2 + m^2)}$$

$$T_h \frac{1}{(k^2 + M^2)(k^2 + m^2)} = \frac{1}{k^2 + M^2} \frac{1}{k^2} \left[1 - \frac{m^2}{k^2} + \frac{m^4}{k^4} - \dots \right]$$

$$I_h = -\frac{M^{-2\varepsilon}}{M^2} \Gamma(\varepsilon) \left[1 + \frac{m^2}{M^2} + \frac{m^4}{M^4} + \dots \right]$$

Infrared divergence

Taylor series in m

Loop integrals with a single scale $M \Rightarrow M^{-2\varepsilon}$

Soft region

$$k \sim m$$

$$I_s = \int \frac{d^d k}{\pi^{d/2}} T_s \frac{1}{(k^2 + M^2)(k^2 + m^2)}$$

$$T_s \frac{1}{(k^2 + M^2)(k^2 + m^2)} = \frac{1}{M^2} \frac{1}{k^2 + m^2} \left[1 - \frac{k^2}{M^2} + \frac{k^4}{M^4} - \dots \right]$$

$$I_s = \frac{m^{-2\varepsilon}}{M^2} \Gamma(\varepsilon) \left[1 + \frac{m^2}{M^2} + \frac{m^4}{M^4} + \dots \right]$$

Ultraviolet divergence

Taylor series in $1/M$

Loop integrals with a single scale $m \Rightarrow m^{-2\varepsilon}$

Proof

$$m \ll \Lambda \ll M$$

$$\begin{aligned} I &= \int_{\mathbf{k} > \Lambda} \frac{d^d \mathbf{k}}{\pi^{d/2}} \frac{1}{(k^2 + M^2)(k^2 + m^2)} \\ &\quad + \int_{\mathbf{k} < \Lambda} \frac{d^d \mathbf{k}}{\pi^{d/2}} \frac{1}{(k^2 + M^2)(k^2 + m^2)} \\ &= \int_{\mathbf{k} > \Lambda} \frac{d^d \mathbf{k}}{\pi^{d/2}} T_h \frac{1}{(k^2 + M^2)(k^2 + m^2)} \\ &\quad + \int_{\mathbf{k} < \Lambda} \frac{d^d \mathbf{k}}{\pi^{d/2}} T_s \frac{1}{(k^2 + M^2)(k^2 + m^2)} \\ &= I_h + I_s - \Delta I \end{aligned}$$

$$\begin{aligned}
\Delta I &= \int_{\mathbf{k} < \Lambda} \frac{d^d \mathbf{k}}{\pi^{d/2}} T_h \frac{1}{(k^2 + M^2)(k^2 + m^2)} \\
&\quad + \int_{\mathbf{k} > \Lambda} \frac{d^d \mathbf{k}}{\pi^{d/2}} T_s \frac{1}{(k^2 + M^2)(k^2 + m^2)} \\
&= \int_{\mathbf{k} < \Lambda} \frac{d^d \mathbf{k}}{\pi^{d/2}} T_s T_h \frac{1}{(k^2 + M^2)(k^2 + m^2)} \\
&\quad + \int_{\mathbf{k} > \Lambda} \frac{d^d \mathbf{k}}{\pi^{d/2}} T_h T_s \frac{1}{(k^2 + M^2)(k^2 + m^2)} \\
&= \int \frac{d^d \mathbf{k}}{\pi^{d/2}} T_s T_h \frac{1}{(k^2 + M^2)(k^2 + m^2)}
\end{aligned}$$

$$\begin{aligned} T_s T_h \frac{1}{(k^2 + M^2)(k^2 + m^2)} &= T_h T_s \frac{1}{(k^2 + M^2)(k^2 + m^2)} \\ &= \frac{1}{M^2 k^2} \left[1 - \frac{m^2}{k^2} + \frac{m^4}{k^4} - \dots \right] \left[1 - \frac{k^2}{M^2} + \frac{k^4}{M^4} - \dots \right] \\ \Delta I &= 0 \end{aligned}$$

No scale

Form factor

Massless φ^3 theory $d = 4$

$$L = \frac{1}{2} (\partial_\mu \varphi) (\partial^\mu \varphi) - \frac{g}{3!} \varphi^3$$

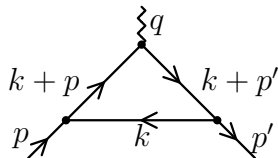
$$[L] = d, [\varphi] = (d - 2)/2, [g] = (6 - d)/2$$

Form factor

Massless φ^3 theory $d = 4$

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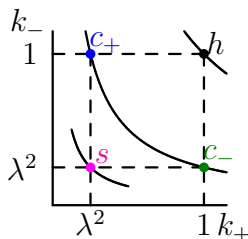
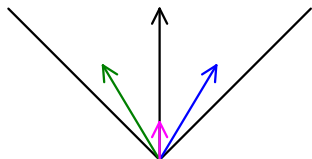
$$I = \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{[-k^2 - i0][-(k + p)^2 - i0][-(k + p')^2 - i0]}$$

$-p^2, -p'^2 \ll -q^2$

Regions

$$\lambda^2 \sim \frac{-p^2}{-q^2} \sim \frac{-p'^2}{-q^2} \ll 1$$

hard	h	$k \sim (1, 1, 1)Q$
collinear	c_+	$k \sim (1, \lambda^2, \lambda)Q$
collinear	c_-	$k \sim (\lambda^2, 1, \lambda)Q$
soft	s	$k \sim (\lambda^2, \lambda^2, \lambda^2)Q$

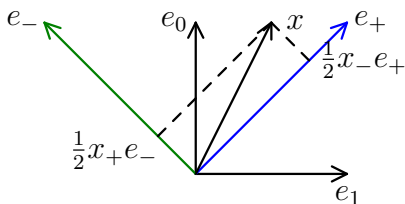


Light-front components

$$e_{\pm} = e_0 \pm e_1 \quad e_{\pm}^2 = 0 \quad e_+ \cdot e_- = 2$$

$$x_{\pm} = x \cdot e_{\pm} = x^0 \mp x^1$$

$$x = \frac{1}{2}(x_- e_+ + x_+ e_-) + x_{\perp} = (x_+, x_-, \vec{x}_{\perp})$$

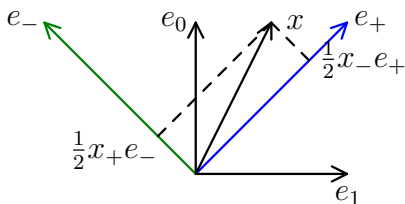


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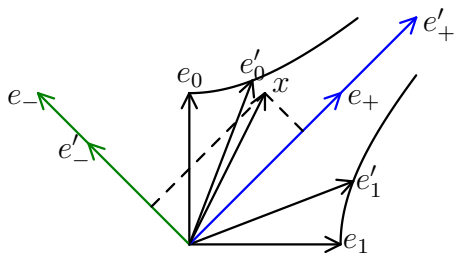


$$x \cdot y = \frac{1}{2}(x_+ y_- + x_- y_+) - \vec{x}_{\perp} \cdot \vec{y}_{\perp}$$

$$x^2 = x_+ x_- - \vec{x}_{\perp}^2$$

Light-front components

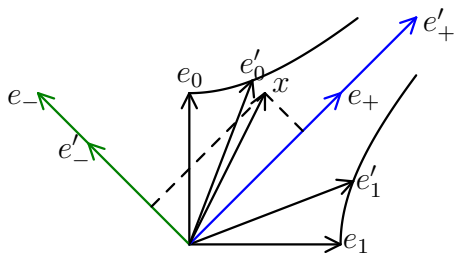
Lorentz transformation



$$x'_\pm = x_\pm e^{\pm\varphi}$$

Light-front components

Lorentz transformation



$$x'_\pm = x_\pm e^{\pm\varphi}$$

$$d^d x = \frac{1}{2} dx_+ dx_- d^{d-2} \vec{x}_\perp$$

Light-front components

$$\begin{aligned}df(x) &= \frac{\partial f}{\partial x^0} dx^0 + \frac{\partial f}{\partial x^1} dx^1 + \frac{\partial f}{\partial \vec{x}_\perp} \cdot d\vec{x}_\perp \\ &= \partial^0 f dx^0 - \partial^1 f dx^1 - \vec{\partial}_\perp \cdot d\vec{x}_\perp \\ \partial^0 &= \frac{\partial}{\partial x^0} \quad \partial^1 = -\frac{\partial}{\partial x^1} \quad \vec{\partial}_\perp = -\frac{\partial}{\partial \vec{x}_\perp}\end{aligned}$$

Light-front components

$$df(x) = \frac{\partial f}{\partial x^0} dx^0 + \frac{\partial f}{\partial x^1} dx^1 + \frac{\partial f}{\partial \vec{x}_\perp} \cdot d\vec{x}_\perp$$

$$= \partial^0 f dx^0 - \partial^1 f dx^1 - \vec{\partial}_\perp \cdot d\vec{x}_\perp$$

$$\partial^0 = \frac{\partial}{\partial x^0} \quad \partial^1 = -\frac{\partial}{\partial x^1} \quad \vec{\partial}_\perp = -\frac{\partial}{\partial \vec{x}_\perp}$$

$$df(x) = \frac{\partial f}{\partial x_+} dx_+ + \frac{\partial f}{\partial x_-} dx_- + \frac{\partial f}{\partial \vec{x}_\perp} \cdot d\vec{x}_\perp$$

$$= \frac{1}{2} \partial_+ f dx_- + \frac{1}{2} \partial_- f dx_+ - \vec{\partial}_\perp \cdot d\vec{x}_\perp$$

$$\partial_+ = 2 \frac{\partial}{\partial x_-} \quad \partial_- = 2 \frac{\partial}{\partial x_+}$$

Light-front components

$$df(x) = \frac{\partial f}{\partial x^0} dx^0 + \frac{\partial f}{\partial x^1} dx^1 + \frac{\partial f}{\partial \vec{x}_\perp} \cdot d\vec{x}_\perp$$

$$= \partial^0 f dx^0 - \partial^1 f dx^1 - \vec{\partial}_\perp \cdot d\vec{x}_\perp$$

$$\partial^0 = \frac{\partial}{\partial x^0} \quad \partial^1 = -\frac{\partial}{\partial x^1} \quad \vec{\partial}_\perp = -\frac{\partial}{\partial \vec{x}_\perp}$$

$$df(x) = \frac{\partial f}{\partial x_+} dx_+ + \frac{\partial f}{\partial x_-} dx_- + \frac{\partial f}{\partial \vec{x}_\perp} \cdot d\vec{x}_\perp$$

$$= \frac{1}{2} \partial_+ f dx_- + \frac{1}{2} \partial_- f dx_+ - \vec{\partial}_\perp \cdot d\vec{x}_\perp$$

$$\partial_+ = 2 \frac{\partial}{\partial x_-} \quad \partial_- = 2 \frac{\partial}{\partial x_+}$$

$$dx_+ = dx^0 - dx^1 \quad dx_- = dx^0 + dx^1$$

$$\partial_+ = \partial^0 - \partial^1 \quad \partial_- = \partial^0 + \partial^1$$

Kinematics

$$p = \frac{1}{2} \left(p_- e_+ - \frac{-p^2}{p_-} e_- \right)$$

$$p' = \frac{1}{2} \left(-\frac{-p'^2}{p'_+} e_+ + p'_+ e_- \right)$$

$$-q^2 = p_- p'_+ - p^2 - p'^2 + \frac{(-p^2)(-p'^2)}{p_- p'_+}$$

Kinematics

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$$-q^2 = p_- p'_+ - p^2 - p'^2 + \frac{(-p^2)(-p'^2)}{p_- p'_+}$$

$$I = I_h + I_{c_+} + I_{c_-} + I_s$$

Hard

$$k \sim (1, 1, 1)Q \quad p \sim (\lambda^2, 1, 0)Q \quad p' \sim (1, \lambda^2, 0)Q \quad d^d k \sim \mathcal{O}(1)$$

$$-k^2 = -k_+ k_- + \vec{k}_\perp^2 = \mathcal{O}(1)$$

$$\begin{aligned} -(k+p)^2 &= -(k_+ + p_+)(k_- + p_-) + \vec{k}_\perp^2 \\ &= \underbrace{-k^2 - p_- k_+}_{\mathcal{O}(1)} - \underbrace{p_+(k_- + p_-)}_{\mathcal{O}(\lambda^2)} \end{aligned}$$

$$\begin{aligned} -(k+p')^2 &= -(k_+ + p'_+)(k_- + p'_-) + \vec{k}_\perp^2 \\ &= \underbrace{-k^2 - p'_+ k_-}_{\mathcal{O}(1)} - \underbrace{p'_-(k_+ + p'_+)}_{\mathcal{O}(\lambda^2)} \end{aligned}$$

$$I_h = \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{(-k^2 - i0)(-k^2 - p_- k_+ - i0)(-k^2 - p'_+ k_- - i0)}$$

On-shell external legs

Hard

Single scale $-q^2 = p_- p'_+ + \mathcal{O}(\lambda^2)$

$$k_- = p_- \mathbf{k}_- \quad k_+ = p'_+ \mathbf{k}_+ \quad \vec{k}_\perp = \sqrt{p_- p'_+} \vec{k}_\perp$$

$$I_h = (-q^2)^{d/2-3}$$

$$\times \int \frac{d^d \mathbf{k}}{i\pi^{d/2}} \frac{1}{(-\mathbf{k}^2 - i0)(-\mathbf{k}^2 - \mathbf{k}_+^2 - i0)(-\mathbf{k}^2 - \mathbf{k}_-^2 - i0)}$$

Hard

Single scale $-q^2 = p_- p'_+ + \mathcal{O}(\lambda^2)$

$$k_- = p_- k_- \quad k_+ = p'_+ k_+ \quad \vec{k}_\perp = \sqrt{p_- p'_+} \vec{k}_\perp$$

$$I_h = (-q^2)^{d/2-3}$$

$$\times \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{(-k^2 - i0)(-k^2 - k_+ - i0)(-k^2 - k_- - i0)}$$

Feynman parametrization

$$\frac{1}{D_1 D_2 D_3} = \int d\alpha_1 d\alpha_2 d\alpha_3 e^{-\alpha_1 D_1 - \alpha_2 D_2 - \alpha_3 D_3}$$

Hard

Single scale $-q^2 = p_- p'_+ + \mathcal{O}(\lambda^2)$

$$k_- = p_- k_- \quad k_+ = p'_+ k_+ \quad \vec{k}_\perp = \sqrt{p_- p'_+} \vec{k}_\perp$$

$$I_h = (-q^2)^{d/2-3}$$

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Feynman parametrization

$$\frac{1}{D_1 D_2 D_3} = \int d\alpha_1 d\alpha_2 d\alpha_3 e^{-\alpha_1 D_1 - \alpha_2 D_2 - \alpha_3 D_3} d\eta \delta(\alpha_1 + \alpha_2 + \alpha_3 - \eta)$$

Hard

Single scale $-q^2 = p_- p'_+ + \mathcal{O}(\lambda^2)$

$$k_- = p_- \mathbf{k}_- \quad k_+ = p'_+ \mathbf{k}_+ \quad \vec{k}_\perp = \sqrt{p_- p'_+} \vec{k}_\perp$$

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Feynman parametrization

$$\frac{1}{D_1 D_2 D_3} = \int d\alpha_1 d\alpha_2 d\alpha_3 e^{-\alpha_1 D_1 - \alpha_2 D_2 - \alpha_3 D_3} d\eta \delta(\alpha_1 + \alpha_2 + \alpha_3 - \eta)$$

$$\alpha_i = \eta x_i$$

Hard

Single scale $-q^2 = p_- p'_+ + \mathcal{O}(\lambda^2)$

$$k_- = p_- \mathbf{k}_- \quad k_+ = p'_+ \mathbf{k}_+ \quad \vec{k}_\perp = \sqrt{p_- p'_+} \vec{k}_\perp$$

$$I_h = (-q^2)^{d/2-3}$$

$$\times \int \frac{d^d \mathbf{k}}{i\pi^{d/2}} \frac{1}{(-\mathbf{k}^2 - i0)(-\mathbf{k}^2 - \mathbf{k}_+^2 - i0)(-\mathbf{k}^2 - \mathbf{k}_-^2 - i0)}$$

Feynman parametrization

$$\frac{1}{D_1 D_2 D_3} = \int d\alpha_1 d\alpha_2 d\alpha_3 e^{-\alpha_1 D_1 - \alpha_2 D_2 - \alpha_3 D_3} d\eta \delta(\alpha_1 + \alpha_2 + \alpha_3 - \eta)$$

$$\alpha_i = \eta x_i$$

$$= \int dx_1 dx_2 dx_3 \delta(x_1 + x_2 + x_3 - 1) d\eta \eta^2 e^{-\eta(x_1 D_1 + x_2 D_2 + x_3 D_3)}$$

Hard

Single scale $-q^2 = p_- p'_+ + \mathcal{O}(\lambda^2)$

$$k_- = p_- \mathbf{k}_- \quad k_+ = p'_+ \mathbf{k}_+ \quad \vec{k}_\perp = \sqrt{p_- p'_+} \vec{k}_\perp$$

$$I_h = (-q^2)^{d/2-3}$$

$$\times \int \frac{d^d \mathbf{k}}{i\pi^{d/2}} \frac{1}{(-\mathbf{k}^2 - i0)(-\mathbf{k}^2 - \mathbf{k}_+^2 - i0)(-\mathbf{k}^2 - \mathbf{k}_-^2 - i0)}$$

Feynman parametrization

$$\frac{1}{D_1 D_2 D_3} = \int d\alpha_1 d\alpha_2 d\alpha_3 e^{-\alpha_1 D_1 - \alpha_2 D_2 - \alpha_3 D_3} d\eta \delta(\alpha_1 + \alpha_2 + \alpha_3 - \eta)$$

$$\alpha_i = \eta x_i$$

$$= \int dx_1 dx_2 dx_3 \delta(x_1 + x_2 + x_3 - 1) d\eta \eta^2 e^{-\eta(x_1 D_1 + x_2 D_2 + x_3 D_3)}$$

$$= \Gamma(3) \int \frac{dx_1 dx_2 dx_3 \delta(x_1 + x_2 + x_3 - 1)}{(x_1 D_1 + x_2 D_2 + x_3 D_3)^3}$$

Hard

Set $-q^2 = 1$

$$I_h = \Gamma(3) \int dx_1 dx_2 dx_3 \delta(x_1 + x_2 + x_3 - 1) \frac{d^d k}{i\pi^{d/2}} \frac{1}{D^3}$$

$$\begin{aligned} D &= -x_1 k^2 - x_2(k^2 + k_+) - x_3(k^2 + k_-) - i0 \\ &= -k^2 - (x_2 e_+ + x_3 e_-) \cdot k - i0 = -k'^2 + x_2 x_3 - i0 \end{aligned}$$

$$k' = k + \frac{1}{2}(x_2 e_+ + x_3 e_-)$$

Hard

Set $-q^2 = 1$

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$$k' = k + \frac{1}{2}(x_2 e_+ + x_3 e_-)$$

$$I_h = \Gamma(3 - d/2) \int dx_2 dx_3 (x_2 x_3)^{d/2-3}$$

Hard

Set $-q^2 = 1$

$$I_h = \Gamma(3) \int dx_1 dx_2 dx_3 \delta(x_1 + x_2 + x_3 - 1) \frac{d^d k}{i\pi^{d/2}} \frac{1}{D^3}$$

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$$k' = k + \frac{1}{2}(x_2 e_+ + x_3 e_-)$$

$$I_h = \Gamma(3 - d/2) \int dx_2 dx_3 (x_2 x_3)^{d/2-3}$$

$$x_2 = zx \quad x_3 = z(1-x)$$

$$I_h = \Gamma(3 - d/2) \int_0^1 dz z^{d-5} \cdot \int_0^1 dx [x(1-x)]^{d/2-3}$$

Hard

Set $-q^2 = 1$

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$$k' = k + \frac{1}{2}(x_2 e_+ + x_3 e_-)$$

$$I_h = \Gamma(3 - d/2) \int dx_2 dx_3 (x_2 x_3)^{d/2-3}$$

$$x_2 = zx \quad x_3 = z(1-x)$$

$$\begin{aligned} I_h &= \Gamma(3 - d/2) \int_0^1 dz z^{d-5} \cdot \int_0^1 dx [x(1-x)]^{d/2-3} \\ &= \frac{\Gamma(3 - d/2)\Gamma^2(d/2 - 2)}{\Gamma(d - 3)} \end{aligned}$$

Hard

$$I_h = \frac{\Gamma(1 + \varepsilon)}{\Gamma(1 - 2\varepsilon)} \Gamma^2(-\varepsilon) (-q^2)^{-1-\varepsilon}$$

Infrared and collinear divergences

Single scale $-q^2$

Collinear c_+

$$k \sim (\lambda^2, 1, \lambda)Q \quad p \sim (\lambda^2, 1, 0)Q \quad p' \sim (1, \lambda^2, 0)Q$$
$$d^d k \sim \mathcal{O}(\lambda^d)$$

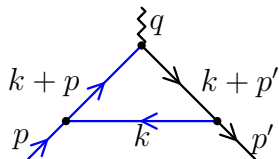
$$-k^2 = -k_+ k_- + \vec{k}_\perp^2 = \mathcal{O}(\lambda^2)$$

$$-(k+p)^2 = -(k_+ + p_+)(k_- + p_-) + \vec{k}_\perp^2 = \mathcal{O}(\lambda^2)$$

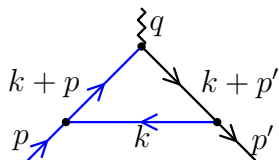
$$\begin{aligned} -(k+p')^2 &= -(k_+ + p'_+)(k_- + p'_-) + \vec{k}_\perp^2 \\ &= \underbrace{-p'_+ k_-}_{\mathcal{O}(1)} - \underbrace{(k^2 + p'_+ p'_-)}_{\mathcal{O}(\lambda^2)} - \underbrace{p'_- k_+}_{\mathcal{O}(\lambda^4)} \end{aligned}$$

$$I_{c_+} = \frac{1}{p'_+} \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{(-k^2 - i0)(-(k+p)^2 - i0)(-k_- - i0)}$$

Collinear c_+



Collinear c_+



Single scale $-p^2 = -p_+p_-$

$$k_+ = (-p_+)k_+ \quad k_- = p_-k_- \quad \vec{k}_\perp = \sqrt{-p_+p_-}\vec{k}_\perp \quad p = (-1, 1, \vec{0})$$

$$I_{c_+} = \frac{(p_+p_-)^{d/2-2}}{p_-p'_+} \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{(-k^2 - i0)(-(k+p)^2 - i0)(-k_- - i0)}$$

Collinear c_+

Feynman parametrization

$$I_{c_+} = \Gamma(3) \int dx_1 dx_2 \delta(x_1 + x_2 - 1) dy \frac{d^d k}{i\pi^{d/2}} \frac{1}{D^3}$$

$$D = -x_1 k^2 - x_2 (k + p)^2 - y k_- - i0$$

$$= -k^2 - (2x_2 p + y e_-) \cdot k + x_2 - i0$$

$$= -k'^2 + x_2 (1 - x_2 + y) - i0$$

$$k' = k + x_2 p + \frac{1}{2} y e_-$$

$$I_{c_+} = \Gamma(3 - d/2) \int dx dy [x(1 - x + y)]^{d/2 - 3}$$

Collinear c_+

$$y = (1 - x)z$$

$$\int = \int_0^\infty dz (1 + z)^{d/2-3} \cdot \int_0^1 dx x^{d/2-3} (1 - x)^{d/2-2}$$

$$I_{c_+} = \frac{\Gamma(2 - d/2)\Gamma(d/2 - 2)\Gamma(d/2 - 1)}{\Gamma(d - 3)}$$

Collinear c_+

$$y = (1 - x)z$$

$$\int = \int_0^\infty dz (1 + z)^{d/2-3} \cdot \int_0^1 dx x^{d/2-3} (1 - x)^{d/2-2}$$

$$I_{c_+} = \frac{\Gamma(2 - d/2)\Gamma(d/2 - 2)\Gamma(d/2 - 1)}{\Gamma(d - 3)}$$

$$I_{c_+} = \frac{\Gamma(d/2 - 1)}{\Gamma(d - 3)} \Gamma(2 - d/2)\Gamma(d/2 - 2) \frac{(-p_+ p_-)^{d/2-2}}{p_- p'_+}$$

$$I_{c_+} = \frac{\Gamma(1 - \varepsilon)}{\Gamma(1 - 2\varepsilon)} \Gamma(\varepsilon)\Gamma(-\varepsilon) \frac{(-p^2)^{-\varepsilon}}{-q^2}$$

Soft

$$k \sim (\lambda^2, \lambda^2, \lambda^2)Q \quad p \sim (\lambda^2, 1, 0)Q \quad p' \sim (1, \lambda^2, 0)Q$$

$$d^d k \sim \mathcal{O}(\lambda^{2d})$$

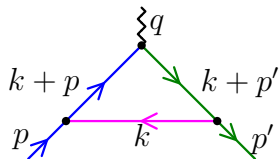
$$-k^2 = -k_+ k_- + \vec{k}_\perp^2 = \mathcal{O}(\lambda^4)$$

$$\begin{aligned} -(k+p)^2 &= -(k_+ + p_+)(k_- + p_-) + \vec{k}_\perp^2 \\ &= \underbrace{-p_-(k_+ + p_+)}_{\mathcal{O}(\lambda^2)} - \underbrace{(k^2 + p_+ k_-)}_{\mathcal{O}(\lambda^4)} \end{aligned}$$

$$\begin{aligned} -(k+p')^2 &= -(k_+ + p'_+)(k_- + p'_-) + \vec{k}_\perp^2 \\ &= \underbrace{-p'_+(k_- + p'_-)}_{\mathcal{O}(\lambda^2)} - \underbrace{(k^2 + p'_- k_+)}_{\mathcal{O}(\lambda^4)} \end{aligned}$$

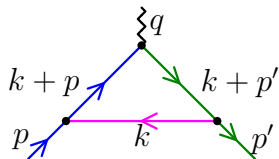
$$I_s = \frac{1}{p_- p'_+} \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{(-k^2 - i0)(-k_+ - p_+ - i0)(-k_- - p'_- - i0)}$$

Soft



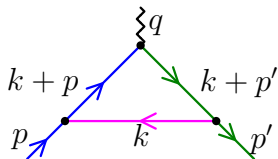
Off-shell collinear legs. Ultraviolet divergence.

Soft



Off-shell collinear legs. Ultraviolet divergence.
Single scale

$$\frac{(-p^2)(-p'^2)}{-q^2} = \underbrace{p_+ p'_-}_{\mathcal{O}(\lambda^4)} + \mathcal{O}(\lambda^6)$$



Off-shell collinear legs. Ultraviolet divergence.
Single scale

$$\frac{(-p^2)(-p'^2)}{-q^2} = \underbrace{p_+ p'_-}_{\mathcal{O}(\lambda^4)} + \mathcal{O}(\lambda^6)$$

$$k_+ = (-p_+)k_+ \quad k_- = (-p'_-)k_- \quad \vec{k}_\perp = \sqrt{(-p_+)(-p'_-)}\vec{k}_\perp$$

$$I_s = \frac{(p_+ p'_-)^{d/2-2}}{p_- p'_+} \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{(-k^2 - i0)(-k_+ + 1 - i0)(-k_- + 1 - i0)}$$

Feynman parametrization

$$I_s = \Gamma(3) \int dy_1 dy_2 \frac{d^d k}{i\pi^{d/2}} \frac{1}{D^3}$$

$$D = -k^2 + y_1(-k_+ + 1) + y_2(-k_- + 1) - i0$$

$$= -k^2 - (y_1 e_+ + y_2 e_-) \cdot k + y_1 + y_2 - i0$$

$$= -k'^2 + y_1 + y_2 + y_1 y_2 - i0$$

$$k' = k + \frac{1}{2}(y_1 e_+ + y_2 e_-)$$

$$I_s = \Gamma(3 - d/2) \int dy_1 dy_2 (y_1 + y_2 + y_1 y_2)^{d/2-3}$$

$$\int = \int_0^\infty dy y^{d/2-2} \int_0^1 dx [1 + yx(1-x)]^{d/2-3}$$

$y_1 = yx$ $y_2 = y(1-x)$

$$\begin{aligned} & y_1 = yx \quad y_2 = y(1-x) \\ \int &= \int_0^\infty dy y^{d/2-2} \int_0^1 dx [1 + yx(1-x)]^{d/2-3} \\ & yx(1-x) = z \\ &= \int_0^\infty dz z^{d/2-2} (1+z)^{d/2-3} \cdot \int_0^1 dx [x(1-x)]^{1-d/2} \end{aligned}$$

$$\begin{aligned} & y_1 = yx \quad y_2 = y(1-x) \\ \int &= \int_0^\infty dy y^{d/2-2} \int_0^1 dx [1 + yx(1-x)]^{d/2-3} \\ & yx(1-x) = z \\ &= \int_0^\infty dz z^{d/2-2} (1+z)^{d/2-3} \cdot \int_0^1 dx [x(1-x)]^{1-d/2} \\ & 1+z = 1/u \\ &= \int_0^1 du u^{3-d} (1-u)^{d/2-2} \cdot \int_0^1 dx [x(1-x)]^{1-d/2} \end{aligned}$$

$$\begin{aligned}
 & y_1 = yx \quad y_2 = y(1-x) \\
 \int &= \int_0^\infty dy y^{d/2-2} \int_0^1 dx [1 + yx(1-x)]^{d/2-3} \\
 & yx(1-x) = z \\
 &= \int_0^\infty dz z^{d/2-2} (1+z)^{d/2-3} \cdot \int_0^1 dx [x(1-x)]^{1-d/2} \\
 & 1+z = 1/u \\
 &= \int_0^1 du u^{3-d} (1-u)^{d/2-2} \cdot \int_0^1 dx [x(1-x)]^{1-d/2} \\
 & \frac{\Gamma(d/2 - q)}{\Gamma(3 - d/2)} \Gamma^2(2 - d/2)
 \end{aligned}$$

$$\begin{aligned} I_s &= \Gamma(d/2 - 1)\Gamma^2(2 - d/2)\frac{(p_+p'_-)^{d/2-2}}{p-p'_+} \\ &= \Gamma(1 - \varepsilon)\Gamma^2(\varepsilon)\frac{1}{-q^2}\left(\frac{(-p^2)(-p'^2)}{-q^2}\right)^{-\varepsilon} \end{aligned}$$

Result

$$\begin{aligned} I &= \frac{\Gamma(1 + \varepsilon)}{\Gamma(1 - 2\varepsilon)} \Gamma^2(-\varepsilon) (-q^2)^{-1-\varepsilon} \\ &+ \frac{\Gamma(1 - \varepsilon)}{\Gamma(1 - 2\varepsilon)} \Gamma(\varepsilon) \Gamma(-\varepsilon) (-q^2)^{-1} [(-p^2)^{-\varepsilon} + (-p'^2)^{-\varepsilon}] \\ &+ \Gamma(1 - \varepsilon) \Gamma^2(\varepsilon) (-q^2)^{-1} \left(\frac{(-p^2)(-p'^2)}{-q^2} \right)^{-\varepsilon} \\ &= \frac{\Gamma(1 + \varepsilon)}{\Gamma(1 - 2\varepsilon)} \Gamma^2(-\varepsilon) (-q^2)^{-1-\varepsilon} \left[1 - \left(\frac{-p^2}{-q^2} \right)^{-\varepsilon} - \left(\frac{-p'^2}{-q^2} \right)^{-\varepsilon} \right. \\ &\quad \left. + \frac{\Gamma(1 + \varepsilon) \Gamma(1 - 2\varepsilon)}{\Gamma(1 - \varepsilon)} \left(\frac{-p^2}{-q^2} \right)^{-\varepsilon} \left(\frac{-p'^2}{-q^2} \right)^{-\varepsilon} \right] \\ I &= \frac{1}{-q^2} \left(\log \frac{-p^2}{-q^2} \log \frac{-p'^2}{-q^2} + \frac{\pi^2}{3} \right) \end{aligned}$$

Divergences cancel

Other regions?

Let's try

$$k \sim (\lambda, \lambda, \lambda)Q \quad p \sim (\lambda^2, 1, 0)Q \quad p' \sim (1, \lambda^2, 0)Q$$

$$-k^2 = -k_+k_- + \vec{k}_\perp^2 = \mathcal{O}(\lambda^2)$$

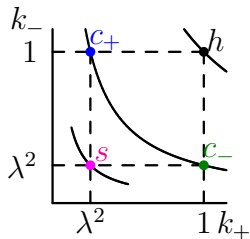
$$\begin{aligned} -(k+p)^2 &= -(k_+ + p_+)(k_- + p_-) + \vec{k}_\perp^2 \\ &= -\underbrace{p_-k_+}_{\mathcal{O}(\lambda)} + \mathcal{O}(\lambda^2) \end{aligned}$$

$$\begin{aligned} -(k+p')^2 &= -(k_+ + p'_+)(k_- + p'_-) + \vec{k}_\perp^2 \\ &= -\underbrace{p'_+k_-}_{\mathcal{O}(\lambda)} + \mathcal{O}(\lambda^2) \end{aligned}$$

$$I_{sh} = \frac{1}{p_-p'_+} \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{(-k^2 - i0)(-k_+ - i0)(-k_- - i0)} = 0$$

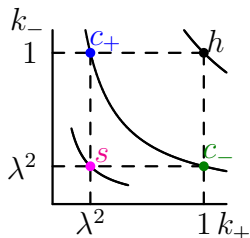
Effective theory

	k	k^2
Hard	$(1, 1, 1)Q$	Q^2
Collinear c_+	$(\lambda^2, 1, \lambda)Q$	$Q^2\lambda^2$
Collinear c_-	$(1, \lambda^2, \lambda)Q$	$Q^2\lambda^2$
Soft	$(\lambda^2, \lambda^2, \lambda^2)Q$	$Q^2\lambda^4$



Effective theory

	k	k^2
Hard	$(1, 1, 1)Q$	Q^2
Collinear c_+	$(\lambda^2, 1, \lambda)Q$	$Q^2\lambda^2$
Collinear c_-	$(1, \lambda^2, \lambda)Q$	$Q^2\lambda^2$
Soft	$(\lambda^2, \lambda^2, \lambda^2)Q$	$Q^2\lambda^4$



$$\varphi \rightarrow \varphi_{c_+} + \varphi_{c_-} + \varphi_s$$

Effective theory

	k	k^2	
Hard	$(1, 1, 1)Q$	Q^2	
Collinear c_+	$(\lambda^2, 1, \lambda)Q$	$Q^2\lambda^2$	
Collinear c_-	$(1, \lambda^2, \lambda)Q$	$Q^2\lambda^2$	
Soft	$(\lambda^2, \lambda^2, \lambda^2)Q$	$Q^2\lambda^4$	

$$\varphi \rightarrow \varphi_{c_+} + \varphi_{c_-} + \varphi_s$$

$$L = L_{c_+} + L_{c_-} + L_s + L_{cs}$$

$$L_{c_+} = \frac{1}{2}(\partial_\mu \varphi_{c_+})(\partial^\mu \varphi_{c_+}) - \frac{g}{3!}\varphi_{c_+}^3$$

$$L_{c_-} = \frac{1}{2}(\partial_\mu \varphi_{c_-})(\partial^\mu \varphi_{c_-}) - \frac{g}{3!}\varphi_{c_-}^3$$

$$L_s = \frac{1}{2}(\partial_\mu \varphi_s)(\partial^\mu \varphi_s) - \frac{g}{3!}\varphi_s^3$$

Power counting

Soft

$$k_s \sim \lambda^2 \quad \partial_s \sim \lambda^2 \quad x_s \sim \lambda^{-2}$$

Power counting

Soft

$$k_s \sim \lambda^2 \quad \partial_s \sim \lambda^2 \quad x_s \sim \lambda^{-2}$$

$$\langle T \varphi_s(x) \varphi_s(0) \rangle = \int \frac{d^d k}{(2\pi)^d} e^{-ik \cdot x} \frac{i}{k^2 + i0}$$

$$\varphi_s \sim \lambda^{d-2}$$

Power counting

Soft

$$k_s \sim \lambda^2 \quad \partial_s \sim \lambda^2 \quad x_s \sim \lambda^{-2}$$

$$\langle T\varphi_s(x)\varphi_s(0) \rangle = \int \frac{d^d k}{(2\pi)^d} e^{-ik \cdot x} \frac{i}{k^2 + i0}$$

$$\varphi_s \sim \lambda^{d-2}$$

$$L_s \sim (\partial_\mu \varphi_s)(\partial^\mu \varphi_s) \sim \lambda^{2d} \quad S_s = \int d^d x L_s \sim 1$$

Power counting

Collinear c_+ ($k_+x_- \sim 1$, $k_-x_+ \sim 1$, $k_\perp x_\perp \sim 1$)

$$k_{c_+} \sim (\lambda^2, 1, \lambda) \quad \partial_{c_+} \sim (\lambda^2, 1, \lambda) \quad x_{c_+} \sim (1, \lambda^{-2}, \lambda^{-1})$$

Power counting

Collinear c_+ ($k_+x_- \sim 1$, $k_-x_+ \sim 1$, $k_\perp x_\perp \sim 1$)

$$k_{c_+} \sim (\lambda^2, 1, \lambda) \quad \partial_{c_+} \sim (\lambda^2, 1, \lambda) \quad x_{c_+} \sim (1, \lambda^{-2}, \lambda^{-1})$$

$$\langle T \varphi_{c_+}(x) \varphi_{c_+}(0) \rangle = \int \frac{d^d k}{(2\pi)^d} e^{-ik \cdot x} \frac{i}{k^2 + i0}$$

$$\varphi_{c_+} \sim \lambda^{d/2-1}$$

Power counting

Collinear c_+ ($k_+x_- \sim 1$, $k_-x_+ \sim 1$, $k_\perp x_\perp \sim 1$)

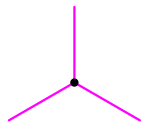
$$k_{c_+} \sim (\lambda^2, 1, \lambda) \quad \partial_{c_+} \sim (\lambda^2, 1, \lambda) \quad x_{c_+} \sim (1, \lambda^{-2}, \lambda^{-1})$$

$$\langle T \varphi_{c_+}(x) \varphi_{c_+}(0) \rangle = \int \frac{d^d k}{(2\pi)^d} e^{-ik \cdot x} \frac{i}{k^2 + i0}$$

$$\varphi_{c_+} \sim \lambda^{d/2-1}$$

$$L_{c_+} \sim (\partial_+ \varphi_{c_+})(\partial_- \varphi_{c_+}) \sim \lambda^d \quad S = \int d^d x L \sim 1$$

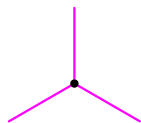
Interactions



$$-\frac{g}{3!}\varphi_s^3$$

$$(\lambda^2, \lambda^2, \lambda^2) + (\lambda^2, \lambda^2, \lambda^2) = (\lambda^2, \lambda^2, \lambda^2)$$

Interactions



$$-\frac{g}{3!}\varphi_s^3$$

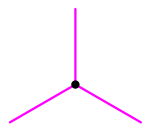
$$(\lambda^2, \lambda^2, \lambda^2) + (\lambda^2, \lambda^2, \lambda^2) = (\lambda^2, \lambda^2, \lambda^2)$$



$$-\frac{g}{3!}\varphi_{c+}^3$$

$$(\lambda^2, 1, \lambda) + (\lambda^2, 1, \lambda) = (\lambda^2, 1, \lambda)$$

Interactions



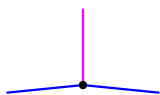
$$-\frac{g}{3!}\varphi_s^3$$

$$(\lambda^2, \lambda^2, \lambda^2) + (\lambda^2, \lambda^2, \lambda^2) = (\lambda^2, \lambda^2, \lambda^2)$$



$$-\frac{g}{3!}\varphi_{c_+}^3$$

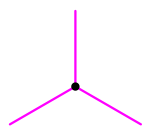
$$(\lambda^2, 1, \lambda) + (\lambda^2, 1, \lambda) = (\lambda^2, 1, \lambda)$$



$$L_{cs} = -\frac{g}{2}\varphi_{c_+}^2\varphi_s - \frac{g}{2}\varphi_{c_-}^2\varphi_s$$

$$(\lambda^2, 1, \lambda) + (\lambda^2, \lambda^2, \lambda^2) = (\lambda^2, 1, \lambda)$$

Interactions



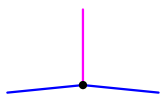
$$-\frac{g}{3!}\varphi_s^3$$

$$(\lambda^2, \lambda^2, \lambda^2) + (\lambda^2, \lambda^2, \lambda^2) = (\lambda^2, \lambda^2, \lambda^2)$$



$$-\frac{g}{3!}\varphi_{c_+}^3$$

$$(\lambda^2, 1, \lambda) + (\lambda^2, 1, \lambda) = (\lambda^2, 1, \lambda)$$



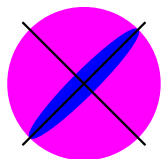
$$L_{cs} = -\frac{g}{2}\varphi_{c_+}^2\varphi_s - \frac{g}{2}\varphi_{c_-}^2\varphi_s$$

$$(\lambda^2, 1, \lambda) + (\lambda^2, \lambda^2, \lambda^2) = (\lambda^2, 1, \lambda)$$

Other interactions ($\varphi_{c_+}^2\varphi_{c_-}$, $\varphi_{c_+}\varphi_{c_-}\varphi_s$, $\varphi_{c_+}^2\varphi_s$, ...) are not allowed by momentum conservation

Multipole expansion

$$S_{cs} = \frac{g}{2} \int d^d x \varphi_{c_+}^2(x) \varphi_s(x)$$



$$x \sim (1, \lambda^{-2}, \lambda^{-1}) Q^{-1}$$

$$x \sim (\lambda^{-2}, \lambda^{-2}, \lambda^{-2}) Q^{-1}$$

$$\partial \sim (\lambda^2, \lambda^2, \lambda^2) Q$$

$$\begin{aligned} \varphi_s(x) = & \varphi_s(\bar{x}_-) - \underbrace{\vec{x}_\perp \cdot \vec{\partial}_\perp}_{\mathcal{O}(\lambda)} \varphi_s(\bar{x}_-) + \frac{1}{2} \underbrace{x_+ \partial_-}_{\mathcal{O}(\lambda^2)} \varphi_s(\bar{x}_-) \\ & + \frac{1}{2} \underbrace{x_{\perp i} x_{\perp j} \partial_{\perp i} \partial_{\perp j}}_{\mathcal{O}(\lambda^2)} \varphi_s(\bar{x}_-) + \mathcal{O}(\lambda^3) \quad \bar{x}_- \equiv \frac{1}{2} x_- e_+^\mu \end{aligned}$$

$$S_{cs} = \frac{g}{2} \int d^d x \varphi_{c_+}^2(x) \varphi_s(\bar{x}_-) + \mathcal{O}(\lambda)$$

Multipole expansion

$$\begin{aligned} S_{cs} &= \frac{g}{2} \int d^d x \varphi_{c_+}^2(x) \varphi_s(x) \\ &= \frac{g}{2} \int d^d x \frac{d^d k_1}{(2\pi)^d} \frac{d^d k_2}{(2\pi)^d} \frac{d^d k_s}{(2\pi)^d} e^{-i(k_1+k_2+k_s)\cdot x} \tilde{\varphi}_{c_+}(k_1) \tilde{\varphi}_{c_+}(k_2) \tilde{\varphi}_s(k_s) \end{aligned}$$

$$k_1 + k_2 + k_s = (\lambda^2, 1, \lambda)Q \quad x \sim (1, \lambda^{-2}, \lambda^{-1})Q^{-1}$$

$$k_s \sim (\lambda^2, \lambda^2, \lambda^2)Q$$

$$k_s \cdot x = \frac{1}{2} \underbrace{k_{s+} x_-}_{\mathcal{O}(1)} + \frac{1}{2} \underbrace{k_{s-} x_+}_{\mathcal{O}(\lambda^2)} - \underbrace{\vec{k}_{s\perp} \cdot \vec{x}_\perp}_{\mathcal{O}(\lambda)}$$

φ_s when interacting with φ_{c_+} carries $k_s = (k_{s+}, 0, \vec{0})$

$$\int \frac{d^d k_s}{(2\pi)^d} e^{i k_s \cdot x} \tilde{\varphi}_s(k_s) = \varphi_s(\bar{x}_-)$$

SCET Lagrangian

$$\begin{aligned} L = & \frac{1}{2}(\partial_\mu\varphi_{c_+}(x))(\partial^\mu\varphi_{c_+}(x)) - \frac{g}{3!}\varphi_{c_+}^3(x) \\ & + \frac{1}{2}(\partial_\mu\varphi_{c_-}(x))(\partial^\mu\varphi_{c_-}(x)) - \frac{g}{3!}\varphi_{c_-}^3(x) \\ & + \frac{1}{2}(\partial_\mu\varphi_s(x))(\partial^\mu\varphi_s(x)) - \frac{g}{3!}\varphi_s^3(x) \\ & - \frac{g}{2}\varphi_{c_+}^2(x)\varphi_s(\bar{x}_-) - \frac{g}{2}\varphi_{c_-}^2(x)\varphi_s(\bar{x}_+) + \mathcal{O}(\lambda) \end{aligned}$$

Translation invariant up to $\mathcal{O}(\lambda)$

Matching

Effective theory

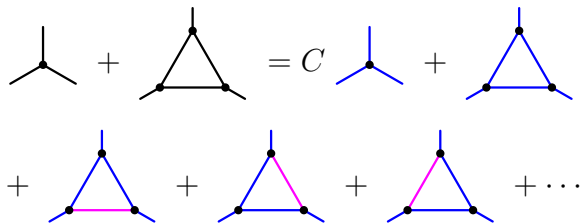
- ▶ Write the most general Lagrangian up to some order with unknown coefficients
- ▶ Calculate some scattering amplitudes in the full theory; if necessary, expand in small parameters up to some order
- ▶ Calculate the same amplitudes in the effective theory
- ▶ Equate, find the coefficients

Matching

For example, the c_+ collinear interaction

$$L_{c_+} = -\frac{g}{3!} C \varphi_{c_+}^3 \quad C = 1 + C_1 g^2 + \dots$$

All 3 external momenta c_+



All external momenta $\parallel e_+ \Rightarrow$ all loop diagrams $= 0$

$$C = 1$$

Current

$$J(x) = \frac{1}{2}\varphi^2(x) \rightarrow J_2(x) + J_3(x) + \dots$$

Kinematics: $c_+ \rightarrow c_-$

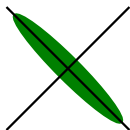
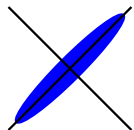
$$J_2 = C_2\varphi_{c_+}\varphi_{c_-} \quad J_3 = \frac{1}{2}C_3 [\varphi_{c_+}^2\varphi_{c_-} + \varphi_{c_+}\varphi_{c_-}^2]$$

Current

$$J(x) = \frac{1}{2}\varphi^2(x) \rightarrow J_2(x) + J_3(x) + \dots$$

Kinematics: $c_+ \rightarrow c_-$

$$J_2 = C_2\varphi_{c_+}\varphi_{c_-} \quad J_3 = \frac{1}{2}C_3 [\varphi_{c_+}^2\varphi_{c_-} + \varphi_{c_+}\varphi_{c_-}^2]$$



$$t \sim 1/Q$$

$$t' \sim 1/Q$$

$$\varphi_{c_+}(x + te_-) = \varphi_{c_+}(x) + \underbrace{t\partial_-}_{\mathcal{O}(1)}\varphi_{c_+}(x) + \frac{1}{2}\underbrace{t^2\partial_-^2}_{\mathcal{O}(1)}\varphi_{c_+}(x) + \dots$$

$$\varphi_{c_-}(x + t'e_+) = \varphi_{c_-}(x) + \underbrace{t'\partial_+}_{\mathcal{O}(1)}\varphi_{c_-}(x) + \frac{1}{2}\underbrace{t'^2\partial_+^2}_{\mathcal{O}(1)}\varphi_{c_-}(x) + \dots$$

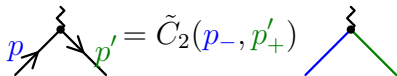
J_2 : tree level

$$J_2(x) = \int dt dt' C_2(t, t') \varphi_{c_+}(x + te_-) \varphi_{c_-}(x + t'e_+)$$

J_2 : tree level

$$J_2(x) = \int dt dt' C_2(t, t') \varphi_{c_+}(x + te_-) \varphi_{c_-}(x + t'e_+)$$

On-shell $p = (0, p_-, \vec{0})$, $p' = (p'_+, 0, \vec{0})$

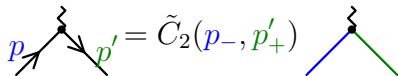

$$p \text{ --- } \text{---} \text{---} p' = \tilde{C}_2(p_-, p'_+) \text{ ---} \text{---}$$

$$\tilde{C}_2(p_-, p'_+) = \int dt dt' e^{ip_-t - ip'_+t'} C_2(t, t')$$

J_2 : tree level

$$J_2(x) = \int dt dt' C_2(t, t') \varphi_{c_+}(x + te_-) \varphi_{c_-}(x + t'e_+)$$

On-shell $p = (0, p_-, \vec{0})$, $p' = (p'_+, 0, \vec{0})$

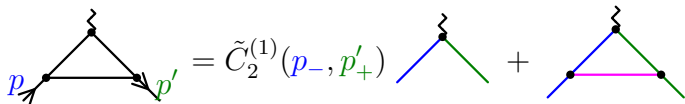

$$p \text{ --- } \text{---} p' = \tilde{C}_2(p_-, p'_+) \text{ ---} \text{---}$$

$$\tilde{C}_2(p_-, p'_+) = \int dt dt' e^{ip_-t - ip'_+t'} C_2(t, t')$$

$$\tilde{C}_2^{(0)}(p_-, p'_+) = 1 \quad C_2^{(0)}(t, t') = \delta(t)\delta(t')$$

J_2 : 1 loop

On-shell $p = (0, p_-, \vec{0})$, $p' = (p'_+, 0, \vec{0})$



The diagram shows an equation for a one-loop triangle diagram. On the left, a triangle with three black lines has an incoming blue line from the bottom-left and an outgoing green line to the bottom-right. A curly line is attached to the top vertex. This is equal to the function $\tilde{C}_2^{(1)}(p_-, p'_+)$ multiplied by a tree-level diagram consisting of a blue line and a green line meeting at a vertex with a curly line. This is then added to another tree-level diagram consisting of a blue line, a green line, and a magenta line meeting at two vertices, with a curly line attached to the top vertex.

SCET loop vanishes

$$\tilde{C}_2^{(1)}(p_-, p'_+) = \frac{g^2}{(4\pi)^{d/2}} I_h(p_-, p'_+)$$

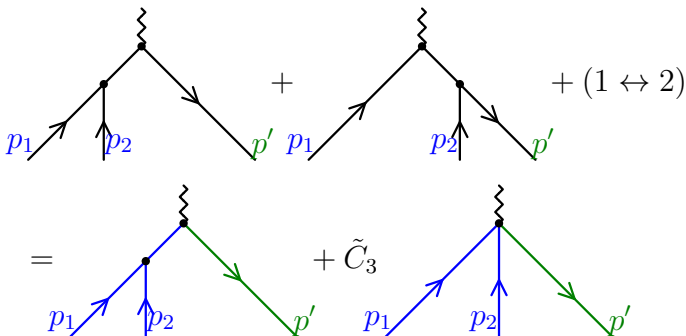
J_3 : tree level

$$J_3(x) = \frac{1}{2} \int dt_1 dt_2 dt' C_3(t_1, t_2, t') \\ \varphi_{c_+}(x + t_1 e_-) \varphi_{c_+}(x + t_2 e_-) \varphi_{c_-}(x + t' e_+) + (+ \leftrightarrow -)$$

J_3 : tree level

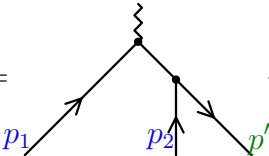
$$J_3(x) = \frac{1}{2} \int dt_1 dt_2 dt' C_3(t_1, t_2, t')$$

$$\varphi_{c_+}(x + t_1 e_-) \varphi_{c_+}(x + t_2 e_-) \varphi_{c_-}(x + t' e_+) + (+ \leftrightarrow -)$$



J_3 : tree level

On-shell $p_1 = (0, p_{1-}, \vec{0})$, $p_2 = (0, p_{2-}, \vec{0})$, $p' = (p'_+, 0, \vec{0})$

$$\tilde{C}_3(p_{1-}, p_{2-}, p'_+) = \text{Diagram} + (1 \leftrightarrow 2)$$


J_3 : tree level

On-shell $p_1 = (0, p_{1-}, \vec{0})$, $p_2 = (0, p_{2-}, \vec{0})$, $p' = (p'_+, 0, \vec{0})$

$$\begin{aligned} \tilde{C}_3(p_{1-}, p_{2-}, p'_+) &= \text{Diagram} + (1 \leftrightarrow 2) \\ &= \frac{g}{p_{2-} p'_+ - i0} + (1 \leftrightarrow 2) \end{aligned}$$

J_3 : tree level

On-shell $p_1 = (0, p_{1-}, \vec{0})$, $p_2 = (0, p_{2-}, \vec{0})$, $p' = (p'_+, 0, \vec{0})$

$$\tilde{C}_3(p_{1-}, p_{2-}, p'_+) = \text{Diagram} + (1 \leftrightarrow 2)$$

$$= \frac{g}{p_{2-} p'_+ - i0} + (1 \leftrightarrow 2)$$

$$C_3(t_1, t_2, t') = \int \frac{dp_{1-}}{2\pi} \frac{dp_{2-}}{2\pi} \frac{dp'_+}{2\pi} e^{-ip_{1-}t_1 - ip_{2-}t_2 + ip'_+t'}$$
$$\times \tilde{C}_3(p_{1-}, p_{2-}, p'_+) = g\delta(t_1)\theta(-t_2)\theta(t') + (1 \leftrightarrow 2)$$

J_3 : coordinate space

$$\begin{aligned} J_3(x) &= g \varphi_{c_+}(x) \int_{-\infty}^0 dt \varphi_{c_+}(x + te_-) \int_0^{\infty} dt' \varphi_{c_-}(x + t'e_+) \\ &= g \varphi_{c_+}(x) \frac{1}{i\partial_- - i0} \varphi_{c_+}(x) \frac{1}{-i\partial_+ - i0} \varphi_{c_-}(x) \end{aligned}$$

Incoming $p_- \rightarrow i\partial_-$, outgoing $p'_+ \rightarrow -i\partial_+$

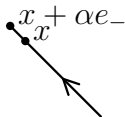
J_3 : coordinate space

$$\begin{aligned} J_3(x) &= g \varphi_{c_+}(x) \int_{-\infty}^0 dt \varphi_{c_+}(x + te_-) \int_0^{\infty} dt' \varphi_{c_-}(x + t'e_+) \\ &= g \varphi_{c_+}(x) \frac{1}{i\partial_- - i0} \varphi_{c_+}(x) \frac{1}{-i\partial_+ - i0} \varphi_{c_-}(x) \end{aligned}$$

Incoming $p_- \rightarrow i\partial_-$, outgoing $p'_+ \rightarrow -i\partial_+$

$$\begin{aligned} \int_{-\infty}^0 dt \varphi(x + te_-) &= \int \frac{d^d k}{(2\pi)^d} e^{-ik \cdot x} \varphi(\tilde{k}) \int_{-\infty}^0 dt e^{-ik_- t + 0t} \\ &= \int \frac{d^d k}{(2\pi)^d} \frac{i}{k_- + i0} e^{-ik \cdot x} \varphi(\tilde{k}) = \frac{i}{i\partial_- + i0} \int \frac{d^d k}{(2\pi)^d} e^{-ik \cdot x} \varphi(\tilde{k}) \\ &= \frac{i}{i\partial_- + i0} \varphi(x) \\ \int_0^{\infty} dt \varphi(x + te_+) &= \frac{-i}{i\partial_+ - i0} \varphi(x) \end{aligned}$$

J_3 : coordinate space



$$F(x) = \int_{-\infty}^0 dt \varphi(x + te_-)$$

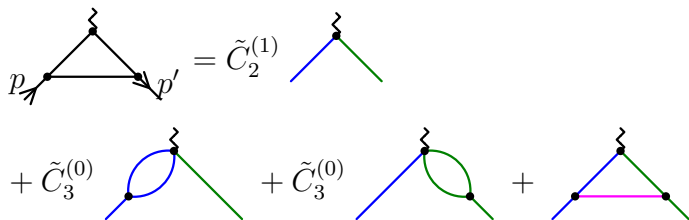
$$F(x + \alpha e_-) = F(x) + \alpha \partial_- F(x)$$

$$= \int_{-\infty}^{\alpha} dt \varphi(x + te_-) = F(x) + \alpha \varphi(x)$$

$$\partial_- F(x) = \varphi(x)$$

Form factor

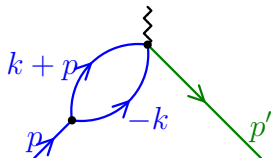
Off-shell $p = (p_+, p_-, \vec{0})$, $p' = (p'_+, p'_-, \vec{0})$,



The diagram shows the expansion of a form factor. On the left, a triangle diagram with external momenta p and p' is equated to $\tilde{C}_2^{(1)}$ times a vertex diagram. This vertex diagram is then expanded into three terms: $\tilde{C}_3^{(0)}$ times a self-energy diagram on the left leg, $\tilde{C}_3^{(0)}$ times a self-energy diagram on the right leg, and a triangle diagram with a pink horizontal internal line.

$$\tilde{C}_2^{(1)} = \frac{g^2}{(4\pi)^{d/2}} I_h$$

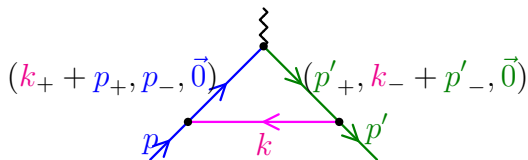
Form factor



$$\begin{aligned}
 & \frac{1}{2} \frac{g^2}{(4\pi)^{d/2}} \int \frac{d^d k}{i\pi^{d/2}} \frac{\tilde{C}_3(k_- + p_-, -k_-, p'_+)}{(-k^2 - i0)(-(k+p)^2 - i0)} \\
 &= \frac{g^2}{(4\pi)^{d/2}} \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{(-k^2 - i0)(-(k+p)^2 - i0)(-p'_+ k_- - i0)} \\
 &= \frac{g^2}{(4\pi)^{d/2}} I_{c_+}
 \end{aligned}$$

Form factor

$$p = (p_+, p_-, \vec{0}), p' = (p'_+, p'_-, \vec{0}), k = (k_+, k_-, \vec{k}_\perp)$$



Momentum conservation up to $\mathcal{O}(\lambda)$

$$\frac{g^2}{(4\pi)^{d/2}} \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{(-k^2 - i0)(-p_-(k_+ + p_+) - i0)(-p'_+(k_- + p'_-) - i0)}$$
$$= \frac{g^2}{(4\pi)^{d/2}} I_s$$

QCD

$$k_c \sim (\lambda^2, 1, \lambda)Q, \quad k_s \sim (\lambda^2, \lambda^2, \lambda^2)Q$$

$$\psi \rightarrow \psi_c + \psi_s \quad A^\mu \rightarrow A_c^\mu + A_s^\mu$$

QCD

$$k_c \sim (\lambda^2, 1, \lambda)Q, \quad k_s \sim (\lambda^2, \lambda^2, \lambda^2)Q$$

$$\psi \rightarrow \psi_c + \psi_s \quad A^\mu \rightarrow A_c^\mu + A_s^\mu$$

Soft

$$\langle T \psi_s(x) \bar{\psi}_s(0) \rangle = \int \frac{d^4 k}{(2\pi)^4} e^{-ik \cdot x} \frac{i \not{k}}{k^2 + i0} \sim \lambda^6$$

$$\langle T A_s^\mu(x) A_s^\nu(0) \rangle = \int \frac{d^4 k}{(2\pi)^4} e^{-ik \cdot x} \frac{i}{k^2 + i0} \left[-g^{\mu\nu} + \xi \frac{k^\mu k^\nu}{k^2} \right] \sim \lambda^4$$

$$\psi_s \sim \lambda^3 \quad A_s^\mu \sim \lambda^2 \quad iD_s^\mu = i\partial_s^\mu + gA_s^\mu \sim \lambda^2$$

Collinear fields

$$\psi_c = \xi + \eta \quad \xi = \frac{\gamma_+ \gamma_-}{4} \psi_c \quad \eta = \frac{\gamma_- \gamma_+}{4} \psi_c$$

$$\langle T\xi(x)\xi(0) \rangle = \int \frac{d^4k}{(2\pi)^4} e^{-ik \cdot x} \frac{i}{k^2 + i0} \frac{\gamma_+ \gamma_-}{4} \underbrace{k}_{\frac{1}{2}k - \gamma_+} \frac{\gamma_- \gamma_+}{4} \sim \lambda^2$$

$$\langle T\eta(x)\eta(0) \rangle = \int \frac{d^4k}{(2\pi)^4} e^{-ik \cdot x} \frac{i}{k^2 + i0} \frac{\gamma_- \gamma_+}{4} \underbrace{k}_{\frac{1}{2}k + \gamma_-} \frac{\gamma_+ \gamma_-}{4} \sim \lambda^4$$

$$\xi \sim \lambda \quad \eta \sim \lambda^2$$

Collinear fields

$$\psi_c = \xi + \eta \quad \xi = \frac{\gamma_+ \gamma_-}{4} \psi_c \quad \eta = \frac{\gamma_- \gamma_+}{4} \psi_c$$

$$\langle T \xi(x) \xi(0) \rangle = \int \frac{d^4 k}{(2\pi)^4} e^{-ik \cdot x} \frac{i}{k^2 + i0} \frac{\gamma_+ \gamma_-}{4} \underbrace{k}_{\frac{1}{2}k_- \gamma_+} \frac{\gamma_- \gamma_+}{4} \sim \lambda^2$$

$$\langle T \eta(x) \eta(0) \rangle = \int \frac{d^4 k}{(2\pi)^4} e^{-ik \cdot x} \frac{i}{k^2 + i0} \frac{\gamma_- \gamma_+}{4} \underbrace{k}_{\frac{1}{2}k_+ \gamma_-} \frac{\gamma_+ \gamma_-}{4} \sim \lambda^4$$

$$\xi \sim \lambda \quad \eta \sim \lambda^2$$

$$\langle T A_c^\mu(x) A_c^\nu(0) \rangle = \int \frac{d^4 k}{(2\pi)^4} e^{-ik \cdot x} \frac{i}{k^2 + i0} \left[-g^{\mu\nu} + \xi \frac{k^\mu k^\nu}{k^2} \right]$$

$$\langle A_{c+} A_{c+} \rangle \sim \lambda^4 \quad \langle A_{c-} A_{c-} \rangle \sim 1 \quad \langle A_{c\perp} A_{c\perp} \rangle \sim \lambda^2$$

$$A_c \sim (\lambda^2, 1, \lambda) \quad iD_c^\mu = i\partial_c^\mu + gA_c^\mu \sim (\lambda^2, 1, \lambda)$$

Soft Lagrangian

$$L_s = \bar{\psi}_s i \not{D}_s \psi_s - \frac{1}{4} F_s^{a\mu\nu} F_{s\mu\nu}^a \quad ig F_s^{\mu\nu} = [iD_s^\mu, iD_s^\nu]$$

Collinear quark Lagrangian

$$\begin{aligned} L &= i(\bar{\xi} + \bar{\eta}) \left(\frac{1}{2} D_+ \gamma_- + \frac{1}{2} D_- \gamma_+ + \not{D}_\perp \right) (\xi + \eta) \\ &= \frac{i}{2} \bar{\xi} D_+ \gamma_- \xi + \frac{i}{2} \bar{\eta} D_- \gamma_+ \eta + i \bar{\xi} \not{D}_\perp \eta + i \bar{\eta} \not{D}_\perp \xi \end{aligned}$$

Collinear quark Lagrangian

$$\begin{aligned} L &= i(\bar{\xi} + \bar{\eta}) \left(\frac{1}{2} D_+ \gamma_- + \frac{1}{2} D_- \gamma_+ + \not{D}_\perp \right) (\xi + \eta) \\ &= \frac{i}{2} \bar{\xi} D_+ \gamma_- \xi + \frac{i}{2} \bar{\eta} D_- \gamma_+ \eta + i \bar{\xi} \not{D}_\perp \eta + i \bar{\eta} \not{D}_\perp \xi \end{aligned}$$

Equations of motion

$$\not{D} \psi_c = \left(\frac{1}{2} D_+ \gamma_- + \frac{1}{2} D_- \gamma_+ + \not{D}_\perp \right) (\xi + \eta) = 0$$

$$\gamma_+ \not{D} \psi_c = 2 D_+ \xi + \not{D}_\perp \eta = 0$$

$$\gamma_- \not{D} \psi_c = 2 D_- \eta + \not{D}_\perp \xi = 0$$

Collinear quark Lagrangian

$$\begin{aligned} L &= i(\bar{\xi} + \bar{\eta}) \left(\frac{1}{2} D_+ \gamma_- + \frac{1}{2} D_- \gamma_+ + \not{D}_\perp \right) (\xi + \eta) \\ &= \frac{i}{2} \bar{\xi} D_+ \gamma_- \xi + \frac{i}{2} \bar{\eta} D_- \gamma_+ \eta + i \bar{\xi} \not{D}_\perp \eta + i \bar{\eta} \not{D}_\perp \xi \end{aligned}$$

Equations of motion

$$\not{D} \psi_c = \left(\frac{1}{2} D_+ \gamma_- + \frac{1}{2} D_- \gamma_+ + \not{D}_\perp \right) (\xi + \eta) = 0$$

$$\gamma_+ \not{D} \psi_c = 2 D_+ \xi + \not{D}_\perp \eta = 0$$

$$\gamma_- \not{D} \psi_c = 2 D_- \eta + \not{D}_\perp \xi = 0$$

$$\eta = -\frac{1}{2} \gamma_- \frac{1}{i D_- + i0} i \not{D}_\perp \xi \quad \bar{\eta} = -\frac{1}{2} \bar{\xi} i \overleftarrow{\not{D}}_\perp \frac{1}{i \overleftarrow{D}_- + i0} \gamma_-$$

Collinear quark Lagrangian

Substituting η via $\xi =$ integrating out η

$$\begin{aligned} L &= \frac{i}{2} \bar{\xi} D_+ \gamma_- \xi - \frac{1}{2} \bar{\xi} i \not{D}_\perp \frac{\gamma_-}{i D_- + i0} i \not{D}_\perp \xi - \frac{1}{2} \bar{\xi} i \overleftarrow{\not{D}}_\perp \frac{\gamma_-}{i \overleftarrow{D}_- + i0} i \not{D}_\perp \xi \\ &\quad + \frac{1}{8} \bar{\xi} i \overleftarrow{\not{D}}_\perp \frac{\gamma_-}{i \overleftarrow{D}_- + i0} i D_- \gamma_+ \frac{\gamma_-}{D_-} \not{D}_\perp \xi \\ &= \frac{1}{2} \bar{\xi} i D_+ \gamma_- \xi + \frac{1}{2} \bar{\xi} i \not{D}_\perp \frac{1}{i D_- + i0} i \not{D}_\perp \gamma_- \xi \end{aligned}$$

Leading-order Lagrangian

$$iD_- \rightarrow iD_{c-} = i\partial_- + gA_{c-}, \quad iD_\perp \rightarrow iD_{c\perp} = i\partial_\perp + gA_{c\perp}$$

$$\begin{aligned}
 L = & \bar{\psi}_s i \not{D}_s \psi_s - \frac{1}{4} F_s^{a\mu\nu} F_{s\mu\nu}^a \\
 & + \frac{1}{2} \bar{\xi} \left[i \underbrace{D_+}_{i\partial_+ + gA_{c+} + gA_{s+}(\bar{x}_-)} + i \not{D}_{c\perp} \frac{1}{iD_{c-} + i0} i \not{D}_{c\perp} \right] \gamma_- \xi \\
 & - \frac{1}{4} F_c^{a\mu\nu} F_{c\mu\nu}^a \qquad \bar{x}_- \equiv \frac{1}{2} x_- e_+^\mu
 \end{aligned}$$

Leading-order Lagrangian

$$iD_- \rightarrow iD_{c-} = i\partial_- + gA_{c-}, \quad iD_\perp \rightarrow iD_{c\perp} = i\partial_\perp + gA_{c\perp}$$

$$L = \bar{\psi}_s i\not{D}_s \psi_s - \frac{1}{4} F_s^{a\mu\nu} F_{s\mu\nu}^a$$

$$+ \frac{1}{2} \bar{\xi} \left[i \underbrace{D_+}_{i\partial_+ + gA_{c+} + gA_{s+}(\bar{x}_-)} + i\not{D}_{c\perp} \frac{1}{iD_{c-} + i0} i\not{D}_{c\perp} \right] \gamma_- \xi$$

$$- \frac{1}{4} F_c^{a\mu\nu} F_{c\mu\nu}^a \quad \bar{x}_- \equiv \frac{1}{2} x_- e_+^\mu$$

$$igF_c^{\mu\nu} = [iD^\mu, iD^\nu]$$

$$iD^\mu = \frac{1}{2} iD_+ e_-^\mu + \frac{1}{2} iD_{c-} e_+^\mu + iD_{c\perp}^\mu$$

Soft gauge transformations

$$U_s(x) = e^{if_s^a(x)t^a} \quad x \sim (\lambda^{-2}, \lambda^{-2}, \lambda^{-2}) \quad \partial \sim (\lambda^2, \lambda^2, \lambda^2)$$

$$\psi_s(x) \rightarrow U_s(x)\psi_s(x)$$

$$A_s^\mu(x) \rightarrow U_s(x)A_s^\mu(x)U_s^{-1}(x) - \frac{i}{g}(\partial^\mu U_s(x))U_s^{-1}(x)$$

Soft gauge transformations

$$U_s(x) = e^{if_s^a(x)t^a} \quad x \sim (\lambda^{-2}, \lambda^{-2}, \lambda^{-2}) \quad \partial \sim (\lambda^2, \lambda^2, \lambda^2)$$

$$\psi_s(x) \rightarrow U_s(x)\psi_s(x)$$

$$A_s^\mu(x) \rightarrow U_s(x)A_s^\mu(x)U_s^\dagger(x) - \frac{i}{g}(\partial^\mu U_s(x))U_s^{-1}(x)$$

$$\xi(x) \rightarrow U_s(\bar{x}_-)\xi(x)$$

$$A_c^\mu(x) \rightarrow U_s(\bar{x}_-)A_c^\mu(x)U_s^{-1}(\bar{x}_-)$$

Soft gauge transformations

$$U_s(x) = U_s(\bar{x}_-) - \underbrace{\vec{x}_\perp \cdot \vec{\partial}_\perp U_s(\bar{x}_-)}_{\mathcal{O}(\lambda)} + \mathcal{O}(\lambda^2)$$

Soft gauge transformations

$$U_s(x) = U_s(\bar{x}_-) - \underbrace{\vec{x}_\perp \cdot \vec{\partial}_\perp U_s(\bar{x}_-)}_{\mathcal{O}(\lambda)} + \mathcal{O}(\lambda^2)$$

$(\partial^\mu U_s) U_s^{-1} \sim \lambda^2$: not needed for A_{c-} and $A_{c\perp}$
 A_{c+} only appears in D_+

$$\begin{aligned} & A_{c+}(x) + A_{s+}(\bar{x}_-) \rightarrow \\ & U_s(\bar{x}_-) [A_{c+}(x) + A_{s+}(\bar{x}_-)] U_s^{-1}(\bar{x}_-) \\ & + \frac{i}{g} U_s(\bar{x}_-) (\partial_+ U_s^{-1}(\bar{x}_-)) \end{aligned}$$

Soft gauge transformations

$$U_s(x) = U_s(\bar{x}_-) - \underbrace{\vec{x}_\perp \cdot \vec{\partial}_\perp U_s(\bar{x}_-)}_{\mathcal{O}(\lambda)} + \mathcal{O}(\lambda^2)$$

$(\partial^\mu U_s) U_s^{-1} \sim \lambda^2$: not needed for A_{c-} and $A_{c\perp}$
 A_{c+} only appears in D_+

$$\begin{aligned} & A_{c+}(x) + A_{s+}(\bar{x}_-) \rightarrow \\ & U_s(\bar{x}_-) [A_{c+}(x) + A_{s+}(\bar{x}_-)] U_s^{-1}(\bar{x}_-) \\ & + \frac{i}{g} U_s(\bar{x}_-) (\partial_+ U_s^{-1}(\bar{x}_-)) \\ & iD_+ \rightarrow U_s(\bar{x}_-) iD_+ U_s^{-1}(\bar{x}_-) \end{aligned}$$

Collinear gauge transformations

$$U_c(x) = e^{if_c^a(x)t^a} \quad x \sim (1, \lambda^{-2}, \lambda^{-1}) \quad \partial \sim (\lambda^2, 1, \lambda)$$
$$\psi_s(x) \rightarrow \psi_s(x) \quad A_s^\mu(x) \rightarrow A_s^\mu(x)$$

Collinear gauge transformations

$$U_c(x) = e^{if_c^a(x)t^a} \quad x \sim (1, \lambda^{-2}, \lambda^{-1}) \quad \partial \sim (\lambda^2, 1, \lambda)$$

$$\psi_s(x) \rightarrow \psi_s(x) \quad A_s^\mu(x) \rightarrow A_s^\mu(x)$$

$$\xi(x) \rightarrow U_c(x)\xi(x)$$

$$A_c^\mu(x) \rightarrow U_c(x)A_c^\mu(x)U_c^{-1}(x)$$

$$+ \frac{1}{g}U_c(x) \left[i\partial^\mu + \frac{1}{2}gA_{s+}(\bar{x}_-)e^\mu, U_c^{-1}(x) \right]$$

Collinear gauge transformations

$$A_{c\perp}(x) \rightarrow U_c(x)A_{c\perp}(x)U_c^{-1}(x) + \frac{i}{g}U_c(x) (\partial_{\perp}U_c^{-1}(x))$$

$$A_{c-}(x) \rightarrow U_c(x)A_{c-}(x)U_c^{-1}(x) + \frac{i}{g}U_c(x) (\partial_{-}U_c^{-1}(x))$$

$$A_{c+}(x) \rightarrow U_c(x)A_{c+}(x)U_c^{-1}(x) + \frac{i}{g}U_c(x) [D_{s+}(\bar{x}_{-}), U_c^{-1}(x)]$$

Collinear gauge transformations

$$A_{c\perp}(x) \rightarrow U_c(x)A_{c\perp}(x)U_c^{-1}(x) + \frac{i}{g}U_c(x) (\partial_{\perp}U_c^{-1}(x))$$

$$A_{c-}(x) \rightarrow U_c(x)A_{c-}(x)U_c^{-1}(x) + \frac{i}{g}U_c(x) (\partial_{-}U_c^{-1}(x))$$

$$A_{c+}(x) \rightarrow U_c(x)A_{c+}(x)U_c^{-1}(x) + \frac{i}{g}U_c(x) [D_{s+}(\bar{x}_{-}), U_c^{-1}(x)]$$

$$iD_{+} \rightarrow U_c(x)iD_{+}U_c^{-1}(x)$$

Reparametrization invariance

$$\begin{aligned} e_+ &\rightarrow (1 + \varepsilon)e_+ & e_- &\rightarrow (1 - \varepsilon)e_- \\ k_+ &\rightarrow (1 + \varepsilon)k_+ & k_- &\rightarrow (1 - \varepsilon)k_- \end{aligned}$$

Lorentz boost

Reparametrization invariance

$$\begin{aligned}e_+ &\rightarrow (1 + \varepsilon)e_+ & e_- &\rightarrow (1 - \varepsilon)e_- \\k_+ &\rightarrow (1 + \varepsilon)k_+ & k_- &\rightarrow (1 - \varepsilon)k_-\end{aligned}$$

Lorentz boost

$$\begin{aligned}e_+ &\rightarrow e_+ & e_- &\rightarrow e_- + \varepsilon_\perp \\k_- &\rightarrow \underbrace{k_-}_{\mathcal{O}(1)} + \underbrace{k \cdot \varepsilon_\perp}_{\mathcal{O}(\lambda)} \\k_\perp &= k - \frac{1}{2}(k \cdot e_+)e_- - \frac{1}{2}(k \cdot e_-)e_+ \\&\rightarrow \underbrace{k_\perp}_{\mathcal{O}(\lambda)} - \underbrace{\frac{1}{2}k_+\varepsilon_\perp}_{\mathcal{O}(\lambda^2)} - \underbrace{\frac{1}{2}(k \cdot \varepsilon_\perp)e_+}_{\mathcal{O}(\lambda)}\end{aligned}$$

Reparametrization invariance

$$e_+ \rightarrow e_+ + \lambda \varepsilon_\perp \quad e_- \rightarrow e_-$$

$$k_+ \rightarrow \underbrace{k_+}_{\mathcal{O}(\lambda^2)} + \underbrace{\lambda k \cdot \varepsilon_\perp}_{\mathcal{O}(\lambda^2)}$$

$$k_\perp \rightarrow \underbrace{k_\perp}_{\mathcal{O}(\lambda)} - \underbrace{\frac{\lambda}{2}(k \cdot \varepsilon_\perp)e_-}_{\mathcal{O}(\lambda^2)} - \underbrace{\frac{\lambda}{2}k_- \varepsilon_\perp}_{\mathcal{O}(\lambda)}$$

Mixes different orders in λ

Wilson lines

$$[x', x] = P \exp ig \int_x^{x'} dy_\mu A^\mu(y)$$

Wilson lines

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Gauge transformation

$$1 + igA^\mu(x) dx_\mu \rightarrow$$

$$1 + ig \left[U(x)A^\mu(x)U^{-1}(x) - \frac{i}{g} (\partial^\mu U(x)) U^{-1}(x) \right] dx_\mu$$

$$= U(x + dx)U^{-1}(x) + igU(x)A^\mu(x)u^{-1}(x) dx_\mu$$

$$= U(x + dx) [1 + igA^\mu(x) dx_\mu] U^{-1}(x)$$

$$[x', x] \rightarrow U(x')[x', x]U^{-1}(x) \quad \bar{\psi}(x')[x', x]\psi(x) = \text{gauge invariant}$$



Wilson lines

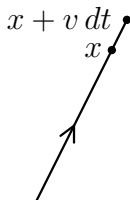


$$W(x) = [x, x - \infty v] = P \exp ig \int_{-\infty}^0 dt A^\mu(x + vt)$$

$$U(x - \infty v) = 1 \quad W(x) \rightarrow U(x)W(x)$$

$$W^{-1}(x)\psi(x), \bar{\psi}(x)W(x) = \text{gauge invariant}$$

Wilson lines



$$W(x) = [x, x - \infty v] = P \exp ig \int_{-\infty}^0 dt A^\mu(x + vt)$$

$$U(x - \infty v) = 1 \quad W(x) \rightarrow U(x)W(x)$$

$$W^{-1}(x)\psi(x), \bar{\psi}(x)W(x) = \text{gauge invariant}$$

$$W(x + v dt) = [1 + igv \cdot A(x) dt] W(x)$$

$$v \cdot \partial W(x) = igv \cdot A(x) W(x) \quad v \cdot DW(x) = 0$$

Collinear-soft interaction: quarks

$$L_q = \frac{1}{2} \bar{\xi}(x) i D_+ \gamma_- \xi(x)$$

$$i D_+ = i \partial_+ + g A_{c+}(x) + g A_{s+}(\bar{x}_-) = i D_{c+} + g A_{s+}(\bar{x}_-)$$

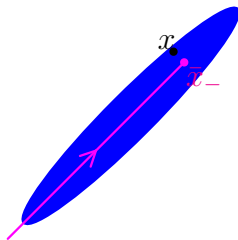
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$$S(x) = P \exp ig \int_{-\infty}^0 dt A_{s+}(x + e_+ t) \quad D_{s+} S(x) = 0$$

$$\xi(x) = S(\bar{x}_-) \xi_0(x) \quad A_c^\mu(x) = S(\bar{x}_-) A_{c0}^\mu(x) S^{-1}(\bar{x}_-)$$



Collinear-soft interaction: quarks

$$\begin{aligned} & iD_+\xi(x) \\ &= [S(\bar{x}_-)i\partial_+ + ((i\partial_+ + gA_s(\bar{x}_-))S(\bar{x}_-)) + gS(\bar{x}_-)A_{c0+}(x)]\xi_0(x) \\ &= S(\bar{x}_-)iD_{c0+}\xi_0(x) \end{aligned}$$

$$L_q = \frac{1}{2}\bar{\xi}_0(x)iD_{c0+}\xi_0(x)$$

Collinear-soft interaction: gluons

$$L_g = -\frac{1}{2} \text{Tr} F_{c\mu\nu} F_c^{\mu\nu} \quad igF_c^{\mu\nu} = [iD^\mu, iD^\nu]$$

$$iD^\mu = \frac{1}{2} iD_+ e_-^\mu + \frac{1}{2} iD_{c-} e_+^\mu + iD_{c\perp}^\mu$$

$$iD_+ = i\partial_+ + gA_{c+}(x) + gA_{s+}(\bar{x}_-)$$

$$A_c(x) = S(\bar{x}_-) A_{c0}(x) S^{-1}(\bar{x}_-)$$

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$$A_c(x) = S(\bar{x}_-) A_{c0}(x) S^{-1}(\bar{x}_-)$$

$$iD_{c-} = i\partial_- + gS(\bar{x}_-) A_{c0-}(x) S^{-1}(\bar{x}_-) = S(\bar{x}_-) iD_{c0-} S^{-1}(\bar{x}_-)$$

$$\text{because } \partial_- S(\bar{x}_-) = 2 \frac{\partial S(\bar{x}_-)}{\partial x_+} = 0$$

$$iD_{c\perp} = i\partial_\perp + gS(\bar{x}_-) A_{c0\perp}(x) S^{-1}(\bar{x}_-)$$

Collinear-soft interaction: gluons

$$\begin{aligned}iD_+ &= S(\bar{x}_-)i\partial_+S^{-1}(\bar{x}_-) - S(\bar{x}_-)(i\partial_+S^{-1}(\bar{x}_-)) \\ &+ gA_{s+}(\bar{x}_-) + gS(\bar{x}_-)A_{c0+}(x)S^{-1}(\bar{x}_-) \\ &= S(\bar{x}_-)iD_{c0+}S^{-1}(\bar{x}_-)\end{aligned}$$

because

$$\begin{aligned}S(\bar{x}_-)(i\partial_+S^{-1}(\bar{x}_-)) &= -(i\partial_+S(\bar{x}_-))S^{-1}(\bar{x}_-) = gA_{s+}(\bar{x}_-), \\ iD_{s+}S &= (i\partial_+ + gA_{s+})S = 0\end{aligned}$$

Collinear-soft interaction: gluons

$$\begin{aligned}iD_+ &= S(\bar{x}_-)i\partial_+S^{-1}(\bar{x}_-) - S(\bar{x}_-)(i\partial_+S^{-1}(\bar{x}_-)) \\ &+ gA_{s+}(\bar{x}_-) + gS(\bar{x}_-)A_{c0+}(x)S^{-1}(\bar{x}_-) \\ &= S(\bar{x}_-)iD_{c0+}S^{-1}(\bar{x}_-)\end{aligned}$$

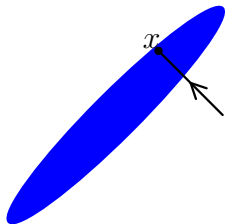
because

$$\begin{aligned}S(\bar{x}_-)(i\partial_+S^{-1}(\bar{x}_-)) &= -(i\partial_+S(\bar{x}_-))S^{-1}(\bar{x}_-) = gA_{s+}(\bar{x}_-), \\ iD_{s+}S &= (i\partial_+ + gA_{s+})S = 0\end{aligned}$$

$$F_c^{\mu\nu} = S(\bar{x}_-)F_{c0}^{\mu\nu}S^{-1}(\bar{x}_-) \quad L_g = -\frac{1}{2}\text{Tr} F_{c0\mu\nu}F_{c0}^{\mu\nu}$$

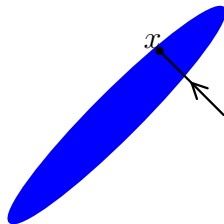
Collinear fields

$$W(x) = P \exp ig \int_{-\infty}^0 dt A_{c-}(x + te_-)$$



Collinear fields

$$W(x) = P \exp ig \int_{-\infty}^0 dt A_{c-}(x + te_-)$$



$$\xi(x) = W(x)\chi(x) \quad \mathcal{A}^\mu(x) = W^{-1}(x)(iD_c^\mu W(x))$$

$$\mathcal{A}_- = 0 \quad \mathcal{A}_+ \sim \lambda^2 \quad \mathcal{A}_\perp \sim \lambda$$

$$i\mathcal{D}^\mu = W^{-1}(x)iD_c^\mu W(x) = i\partial^\mu + \mathcal{A}^\mu$$

Collinear quark Lagrangian

$$\begin{aligned}L_q &= \frac{1}{2}\bar{\xi}iD_{c+}\gamma-\xi + \frac{1}{2}\bar{\xi}i\mathcal{D}_{c\perp}\frac{1}{iD_{c-}+i0}i\mathcal{D}_{c\perp}\gamma-\xi \\&= \frac{1}{2}\bar{\chi}i\mathcal{D}_{c+}\gamma-\chi + \frac{1}{2}\bar{\chi}i\mathcal{D}_{c\perp}\frac{1}{i\partial_-+i0}i\mathcal{D}_{c\perp}\gamma-\chi \\&= \frac{1}{2}\bar{\chi}i\mathcal{D}_{c+}\gamma-\chi - \frac{i}{2}\bar{\chi}\gamma-i\mathcal{D}_{c\perp}\int_{-\infty}^0 dt (i\mathcal{D}_{c\perp}\chi)_{x+te_-}\end{aligned}$$

Collinear quark Lagrangian

$$\begin{aligned}L_q &= \frac{1}{2}\bar{\xi}iD_{c+}\gamma-\xi + \frac{1}{2}\bar{\xi}i\mathcal{D}_{c\perp}\frac{1}{iD_{c-}+i0}i\mathcal{D}_{c\perp}\gamma-\xi \\&= \frac{1}{2}\bar{\chi}i\mathcal{D}_{c+}\gamma-\chi + \frac{1}{2}\bar{\chi}i\mathcal{D}_{c\perp}\frac{1}{i\partial_-+i0}i\mathcal{D}_{c\perp}\gamma-\chi \\&= \frac{1}{2}\bar{\chi}i\mathcal{D}_{c+}\gamma-\chi - \frac{i}{2}\bar{\chi}\gamma-i\mathcal{D}_{c\perp}\int_{-\infty}^0 dt (i\mathcal{D}_{c\perp}\chi)_{x+te_-}\end{aligned}$$

Nonlocality in e_- direction $\sim Q^{-1}$

$$\begin{aligned}W^{-1}(x)iD_{c-}W(x) &= i\partial_- \\ \Rightarrow W^{-1}(x)\frac{1}{iD_{c-}+i0}W(x) &= \frac{1}{i\partial_-+i0}\end{aligned}$$

Collinear gluon Lagrangian

$$W^{-1}F_{c\mu\nu}W = \mathcal{F}_{\mu\nu} = \frac{1}{g} (\partial_\mu \mathcal{A}_\nu - \partial_\nu \mathcal{A}_\mu - i[\mathcal{A}_\mu, \mathcal{A}_\nu])$$

$$L_g = -\frac{1}{2} \text{Tr} F_{c\mu\nu} F_c^{\mu\nu} = -\frac{1}{2g^2} \text{Tr} \mathcal{F}_{\mu\nu} \mathcal{F}^{\mu\nu}$$

Quark vector current

$$J^\mu(x) = \bar{\psi}(x)\gamma^\mu\psi(x)$$

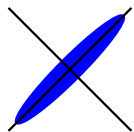
Kinematics: $c_- \rightarrow c_+$

$$\psi(x) \rightarrow S_-(\bar{x}_+)\chi_-(x) \quad (\bar{x}_+ = \frac{1}{2}x_+e_-, \gamma_- \chi_- = 0)$$

$$\bar{\psi}(x) \rightarrow S_+(\bar{x}_-)\bar{\chi}_+(x) \quad (\bar{x}_- = \frac{1}{2}x_-e_+, \gamma_+ \chi_+ = 0)$$

$$J^\mu(x) = \int dt dt' C_V(t, t') \bar{\chi}_+(x + te_-) S_+(\bar{x}_-) S_-(\bar{x}_+) \gamma_\perp^\mu \chi_-(x + t'e_+) + \dots$$

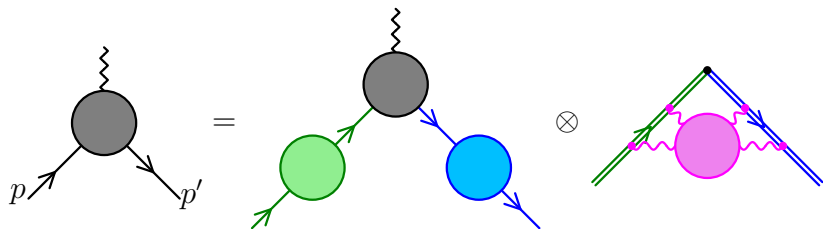
Quark vector current



$$x^\mu \sim (1, 1, \lambda^{-1})Q^{-1}$$

$$J^\mu(x) = \int dt dt' C_V(t, t') \bar{\chi}_+(\bar{x}_+ + x_\perp + te_-) S_+^+(0) S_-(0) \\ \gamma_\perp^\mu \chi_-(\bar{x}_- + x_\perp + t'e_+) + \dots$$

Factorization

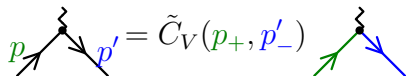


$$F(-q^2, -p^2, -p'^2) = \tilde{C}_V(-q^2) J_-(-p^2) J_+(-p'^2) S \left(\frac{(-p^2)(-p'^2)}{-q^2} \right)$$

Matching coefficient

On-shell $p = (p_+, 0, \vec{0})$, $p' = (0, p'_-, \vec{0})$

Tree level


$$p \text{ --- } \text{---} \text{---} p' = \tilde{C}_V(p_+, p'_-)$$

$$\tilde{C}_V^{(0)}(p_+, p'_-) = 1 \quad C_V^{(0)}(t, t') = \delta(t)\delta(t')$$

Matching coefficient

On-shell $p = (p_+, 0, \vec{0})$, $p' = (0, p'_-, \vec{0})$

Tree level

$$\begin{aligned} \text{Diagram} &= \tilde{C}_V(p_+, p'_-) \text{Diagram} \\ \tilde{C}_V^{(0)}(p_+, p'_-) &= 1 \quad C_V^{(0)}(t, t') = \delta(t)\delta(t') \end{aligned}$$

1 loop

$$\text{Diagram} = \tilde{C}_V^{(1)}(p_+, p'_-) \text{Diagram}$$

SCET loops vanish

Matching coefficient: 1 loop

$$-C_F \frac{g_0^2}{(4\pi)^{d/2}} \int \frac{d^d k}{i\pi^{d/2}} \frac{\gamma^\nu (\not{k} + \not{p}') \gamma_\perp^\mu (\not{k} + \not{p}) \gamma_\nu}{(-k^2 - i0)[-(k+p)^2 - i0][-(k+p')^2 - i0]}$$

Matching coefficient: 1 loop

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Numerator

$$\begin{aligned} & \gamma^\nu \not{k} \gamma_\perp^\mu \not{k} \gamma_\nu + 2\gamma_\perp^\mu \not{k} \not{p}' + 2\not{p} \not{k} \gamma_\perp^\mu + 2p_+ p'_- \gamma_\perp^\mu \\ &= -(d-2) \not{k} \gamma_\perp^\mu \not{k} + 2p'_- k_+ \gamma_\perp^\mu + 2p_+ k_- \gamma_\perp^\mu + 2p_+ p'_- \gamma_\perp^\mu \\ &= [(d-2)k^2 - 2k_\perp^2 + 2p_+ k_- + 2p'_- k_+ + 2p_+ p'_-] \gamma_\perp^\mu \end{aligned}$$

Matching coefficient: 1 loop

$$-C_F \frac{g_0^2}{(4\pi)^{d/2}} \int \frac{d^d k}{i\pi^{d/2}} \frac{\gamma^\nu (\not{k} + \not{p}') \gamma_\perp^\mu (\not{k} + \not{p}) \gamma_\nu}{(-k^2 - i0)[-(k+p)^2 - i0][-(k+p')^2 - i0]}$$

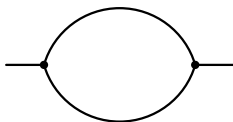
Numerator

$$\begin{aligned} & \gamma^\nu \not{k} \gamma_\perp^\mu \not{k} \gamma_\nu + 2\gamma_\perp^\mu \not{k} \not{p}' + 2\not{p} \not{k} \gamma_\perp^\mu + 2p_+ p'_- \gamma_\perp^\mu \\ &= -(d-2) \not{k} \gamma_\perp^\mu \not{k} + 2p'_- k_+ \gamma_\perp^\mu + 2p_+ k_- \gamma_\perp^\mu + 2p_+ p'_- \gamma_\perp^\mu \\ &= [(d-2)k^2 - 2k_\perp^2 + 2p_+ k_- + 2p'_- k_+ + 2p_+ p'_-] \gamma_\perp^\mu \end{aligned}$$

$$C_V(-q^2) = -C_F \frac{g_0^2}{(4\pi)^{d/2}} \int \frac{d^d k}{i\pi^{d/2}} \frac{2(-q^2) - (d-8)(-k^2) + 2(-k^2)^2/(-q^2)}{(-k^2 - i0)(-k^2 - p_+ k_- - i0)(-k^2 - p'_- k_+ - i0)}$$

$$\begin{aligned}
C_V(-q^2) = & -C_F \frac{g_0^2}{(4\pi)^{d/2}} \left[2(-q^2) I_h \right. \\
& - (d-8) \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{[-(k+p)^2 - i0][-(k+p')^2 - i0]} \\
& \left. + \frac{2}{-q^2} \int \frac{d^d k}{i\pi^{d/2}} \frac{-k^2}{[-(k+p)^2 - i0][-(k+p')^2 - i0]} \right]
\end{aligned}$$

$$\begin{aligned}
& \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{[-(k+p)^2 - i0][-(k+p')^2 - i0]} \\
&= \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{(-k^2 - i0)[-(k+q)^2 - i0]} \\
&= \frac{\Gamma(\varepsilon)\Gamma^2(1-\varepsilon)}{\Gamma(2-2\varepsilon)} (-q^2)^{-\varepsilon}
\end{aligned}$$



$$\begin{aligned}
& \int \frac{d^d k}{i\pi^{d/2}} \frac{-k^2}{[-(k+p)^2 - i0][-(k+p')^2 - i0]} \\
&= \int \frac{d^d k}{i\pi^{d/2}} \frac{-(k-p)^2}{(-k^2 - i0)[-(k+q)^2 - i0]} \\
&= 2p \cdot \int \frac{d^d k}{i\pi^{d/2}} \frac{k}{(-k^2 - i0)[-(k+q)^2 - i0]} \\
&= \frac{2p \cdot q}{q^2} \int \frac{d^d k}{i\pi^{d/2}} \frac{k \cdot q}{(-k^2 - i0)[-(k+q)^2 - i0]} \\
&= -\frac{q^2}{2} \int \frac{d^d k}{i\pi^{d/2}} \frac{1}{(-k^2 - i0)[-(k+q)^2 - i0]}
\end{aligned}$$

Result

$$\begin{aligned} C_V(-q^2) &= 1 - C_F \frac{g_0^2(-q^2)^{-\varepsilon}}{(4\pi)^{d/2}} \frac{\Gamma(1+\varepsilon)\Gamma^2(1-\varepsilon)}{\Gamma(1-2\varepsilon)} \left[\frac{2}{\varepsilon^2} + \frac{3+2\varepsilon}{\varepsilon(1-2\varepsilon)} \right] \\ &= 1 - C_F \frac{g_0^2(-q^2)^{-\varepsilon}}{(4\pi)^{d/2}} e^{-\gamma\varepsilon} \left(\frac{2}{\varepsilon^2} + \frac{3}{\varepsilon} - \frac{\pi^2}{6} + 8 \right) \end{aligned}$$

Renormalized matching coefficient

$$\overline{\text{MS}} \quad \frac{g_0^2}{(4\pi)^{d/2}} = \frac{\alpha_s(\mu)}{4\pi} \mu^{-2\epsilon} e^{\gamma\epsilon}$$

$$C_V(-q^2) = 1 - C_F \frac{\alpha_s(\mu)}{4\pi} \left(\frac{-q^2}{\mu^2} \right)^{-\epsilon} \left(\frac{2}{\epsilon^2} + \frac{3}{\epsilon} - \frac{\pi^2}{6} + 8 \right)$$

Renormalized matching coefficient

$$\overline{\text{MS}} \quad \frac{g_0^2}{(4\pi)^{d/2}} = \frac{\alpha_s(\mu)}{4\pi} \mu^{-2\epsilon} e^{\gamma\epsilon}$$

$$C_V(-q^2) = 1 - C_F \frac{\alpha_s(\mu)}{4\pi} \left(\frac{-q^2}{\mu^2} \right)^{-\epsilon} \left(\frac{2}{\epsilon^2} + \frac{3}{\epsilon} - \frac{\pi^2}{6} + 8 \right)$$

Renormalization

$$C_V(-q^2) = Z^{-1} C_V(-q^2, \mu)$$

$$Z = 1 + C_F \left(\frac{2}{\epsilon^2} - \frac{2}{\epsilon} \log \frac{-q^2}{\mu^2} + \frac{3}{\epsilon} \right) \frac{\alpha_s(\mu)}{4\pi}$$

$$C_V(-q^2, \mu) = 1 + C_F \left(-\log^2 \frac{-q^2}{\mu^2} + 3 \log \frac{-q^2}{\mu^2} + \frac{\pi^2}{6} - 8 \right) \frac{\alpha_s(\mu)}{4\pi}$$

RG equation

$$\frac{d \log C_V(-q^2)}{d \log \mu} = 0 = -\frac{d \log Z}{d \log \mu} + \frac{d \log C_V(-q^2, \mu)}{d \log \mu}$$

$$\frac{d \log \alpha_s(\mu)}{d \log \mu} = -2\varepsilon + \mathcal{O}(\alpha_s)$$

$$\frac{d \log Z}{d \log \mu} = \Gamma(\alpha_s(\mu)) \log \frac{-q^2}{\mu^2} + \gamma_V(\alpha_s(\mu))$$

$$\frac{d \log C_V(-q^2, \mu)}{d \log \mu} = \Gamma(\alpha_s(\mu)) \log \frac{-q^2}{\mu^2} + \gamma_V(\alpha_s(\mu))$$

$$\Gamma(\alpha_s) = 4C_F \frac{\alpha_s}{4\pi} \quad \gamma_V(\alpha_s) = -6C_F \frac{\alpha_s}{4\pi}$$

Solution

$$C_V(-q^2, \mu) = U(\mu_0, \mu) C_V(-q^2, \mu_0)$$

$$\frac{d \log U(\mu_0, \mu)}{d \log \mu} = \Gamma(\alpha_s(\mu)) \log \frac{-q^2}{\mu^2} + \gamma_V(\alpha_s(\mu))$$

$$\Gamma(\alpha_s) = \Gamma_0 \frac{\alpha_s}{4\pi} + \Gamma_1 \left(\frac{\alpha_s}{4\pi} \right)^2 + \dots$$

$$\gamma_V(\alpha_s) = \gamma_{V0} \frac{\alpha_s}{4\pi} + \gamma_{V1} \left(\frac{\alpha_s}{4\pi} \right)^2 + \dots$$

Solution

$$C_V(-q^2, \mu) = U(\mu_0, \mu) C_V(-q^2, \mu_0)$$

$$\frac{d \log U(\mu_0, \mu)}{d \log \mu} = \Gamma(\alpha_s(\mu)) \log \frac{-q^2}{\mu^2} + \gamma_V(\alpha_s(\mu))$$

$$\Gamma(\alpha_s) = \Gamma_0 \frac{\alpha_s}{4\pi} + \Gamma_1 \left(\frac{\alpha_s}{4\pi} \right)^2 + \dots$$

$$\gamma_V(\alpha_s) = \gamma_{V0} \frac{\alpha_s}{4\pi} + \gamma_{V1} \left(\frac{\alpha_s}{4\pi} \right)^2 + \dots$$

$$\frac{d \log \alpha_s(\mu)}{d \log \mu} = -2\beta(\alpha_s(\mu)) \quad \beta(\alpha_s) = \beta_0 \frac{\alpha_s}{4\pi} + \beta_1 \left(\frac{\alpha_s}{4\pi} \right)^2 + \dots$$

Solution

$$C_V(-q^2, \mu) = U(\mu_0, \mu) C_V(-q^2, \mu_0)$$

$$\frac{d \log U(\mu_0, \mu)}{d \log \mu} = \Gamma(\alpha_s(\mu)) \log \frac{-q^2}{\mu^2} + \gamma_V(\alpha_s(\mu))$$

$$\Gamma(\alpha_s) = \Gamma_0 \frac{\alpha_s}{4\pi} + \Gamma_1 \left(\frac{\alpha_s}{4\pi} \right)^2 + \dots$$

$$\gamma_V(\alpha_s) = \gamma_{V0} \frac{\alpha_s}{4\pi} + \gamma_{V1} \left(\frac{\alpha_s}{4\pi} \right)^2 + \dots$$

$$\frac{d \log \alpha_s(\mu)}{d \log \mu} = -2\beta(\alpha_s(\mu)) \quad \beta(\alpha_s) = \beta_0 \frac{\alpha_s}{4\pi} + \beta_1 \left(\frac{\alpha_s}{4\pi} \right)^2 + \dots$$

$$\frac{d \log U(\mu_0, \mu)}{d \log \alpha_s} = -\frac{1}{2\beta(\alpha_s)} \left[\Gamma(\alpha_s) \left(\log \frac{-q^2}{\mu_0^2} - 2 \log \frac{\mu}{\mu_0} \right) + \gamma_V(\alpha_s) \right]$$

$$\log \frac{\mu}{\mu_0} = - \int_{\alpha_s(\mu_0)}^{\alpha_s(\mu)} \frac{d\alpha_s}{\alpha_s} \frac{1}{2\beta(\alpha_s)}$$

Solution

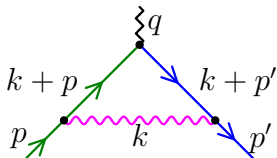
$$U(\mu_0, \mu) = \exp [S(\mu_0, \mu) - A_{\gamma_V}(\mu_0, \mu)] \left(\frac{-q^2}{\mu_0^2} \right)^{-A_{\Gamma}(\mu_0, \mu)}$$

$$A_{\gamma}(\mu_0, \mu) = \int_{\alpha_s(\mu_0)}^{\alpha_s(\mu)} \frac{d\alpha_s}{\alpha_s} \frac{\gamma(\alpha_s)}{2\beta(\alpha_s)} = \frac{\gamma_0}{2\beta_0} \log r + \mathcal{O}(\alpha_s)$$

$$\begin{aligned} S(\mu_0, \mu) &= - \int_{\alpha_s(\mu_0)}^{\alpha_s(\mu)} \frac{d\alpha_s}{\alpha_s} \frac{\Gamma(\alpha_s)}{2\beta(\alpha_s)} \int_{\alpha_s(\mu_0)}^{\alpha_s} \frac{d\alpha'_s}{\alpha'_s} \frac{1}{\beta(\alpha'_s)} \\ &= \frac{\Gamma_0}{2\beta_0^2} \left[\frac{4\pi}{\alpha_s(\mu_0)} \left(\frac{r-1}{r} - \log r \right) + \left(\frac{\Gamma_1}{\Gamma_0} - \frac{\beta_1}{\beta_0} \right) (1-r + \log r) \right. \\ &\quad \left. + \frac{\beta_1}{2\beta_0} \log^2 r \right] + \mathcal{O}(\alpha_s) \end{aligned}$$

$$r = \frac{\alpha_s(\mu)}{\alpha_s(\mu_0)}$$

Soft function



$$\begin{aligned} S_1 &= C_F \int \frac{d^d k}{(2\pi)^d} i g_0 e_+^\mu \frac{i}{e_+ \cdot (k + p') + i0} \frac{i}{e_- \cdot (k + p) + i0} i g_0 e_{-\mu} \\ &\quad \frac{-i}{k^2 + i0} \\ &= -2C_F(-q^2)I_s \end{aligned}$$

Soft function

$$S(\Lambda_s^2) = 1 - 2C_F \frac{g_0^2(\Lambda_s^2)^{-\varepsilon}}{(4\pi)^{d/2}} \Gamma(1 - \varepsilon) \Gamma^2(\varepsilon)$$

$$= 1 - 2C_F \frac{\alpha_s(\mu)}{4\pi} \left(\frac{\Lambda_s^2}{\mu^2} \right)^{-\varepsilon} \left(\frac{1}{\varepsilon^2} + \frac{\pi^2}{4} \right)$$

$$= Z_S^{-1} S(\Lambda_s^2, \mu)$$

$$Z_S = 1 + 2C_F \left(\frac{1}{\varepsilon^2} - \frac{1}{\varepsilon} \log \frac{\Lambda_s^2}{\mu^2} \right) \frac{\alpha_s(\mu)}{4\pi}$$

$$S(\Lambda_s^2, \mu) = 1 - C_F \left(\log^2 \frac{\Lambda_s^2}{\mu^2} + \frac{\pi^2}{2} \right) \frac{\alpha_s(\mu)}{4\pi}$$

$$\frac{d \log S(\Lambda_s^2, \mu)}{d \log \mu} = \Gamma(\alpha_s(\mu)) \log \frac{\Lambda_s^2}{\mu^2} + \gamma_S(\alpha_s(\mu))$$

$$\gamma_S(\alpha_s) = 0 + \mathcal{O}(\alpha_s^2)$$

Resummed form factor

$$F(-q^2, -p^2, -p'^2) = C_V(-q^2, \mu) J(-p^2, \mu) J(-p'^2, \mu) S(\Lambda_s^2, \mu)$$

$$\Lambda_s^2 = \frac{(-p^2)(-p'^2)}{-q^2}$$

$$\begin{aligned} \frac{d \log F(-q^2, -p^2, -p'^2)}{d \log \mu} &= 0 = \frac{d \log C_V(-q^2, \mu)}{d \log \mu} \\ &+ \frac{d \log J(-p^2, \mu)}{d \log \mu} + \frac{d \log J(-p'^2, \mu)}{d \log \mu} + \frac{d \log S(\Lambda_s^2, \mu)}{d \log \mu} \end{aligned}$$

Resummed form factor

$$F(-q^2, -p^2, -p'^2) = C_V(-q^2, \mu) J(-p^2, \mu) J(-p'^2, \mu) S(\Lambda_s^2, \mu)$$

$$\Lambda_s^2 = \frac{(-p^2)(-p'^2)}{-q^2}$$

$$\frac{d \log F(-q^2, -p^2, -p'^2)}{d \log \mu} = 0 = \frac{d \log C_V(-q^2, \mu)}{d \log \mu} + \frac{d \log J(-p^2, \mu)}{d \log \mu} + \frac{d \log J(-p'^2, \mu)}{d \log \mu} + \frac{d \log S(\Lambda_s^2, \mu)}{d \log \mu}$$

$$\frac{d \log C_V(-q^2, \mu)}{d \log \mu} = \Gamma(\alpha_s(\mu)) \log \frac{-q^2}{\mu^2} + \gamma_V(\alpha_s(\mu))$$

$$\frac{d \log J(-p^2, \mu)}{d \log \mu} = -\Gamma(\alpha_s(\mu)) \log \frac{-p^2}{\mu^2} - \gamma_J(\alpha_s(\mu))$$

$$\frac{d \log S(\Lambda_s^2, \mu)}{d \log \mu} = \Gamma(\alpha_s(\mu)) \log \frac{\Lambda_s^2}{\mu^2} + \gamma_S(\alpha_s(\mu))$$

Resummed form factor

$$\Gamma(\alpha_s(\mu)) \left[\log \frac{-q^2}{\mu^2} - \log \frac{-p^2}{\mu^2} - \log \frac{-p'^2}{\mu^2} + \log \frac{(-p^2)(-p'^2)}{(-q^2)\mu^2} \right] \\ + \gamma_V(\alpha_s(\mu)) - 2\gamma_J(\alpha_s(\mu)) + \gamma_S(\alpha_s(\mu)) = 0$$

Resummed form factor

$$\Gamma(\alpha_s(\mu)) \left[\log \frac{-q^2}{\mu^2} - \log \frac{-p^2}{\mu^2} - \log \frac{-p'^2}{\mu^2} + \log \frac{(-p^2)(-p'^2)}{(-q^2)\mu^2} \right] \\ + \gamma_V(\alpha_s(\mu)) - 2\gamma_J(\alpha_s(\mu)) + \gamma_S(\alpha_s(\mu)) = 0$$

