### Search of chiral anomaly in kaon-photon reaction Phys. Rev. **D93**, 094029 (2016); 1512.04438

M. I. Vysotsky, E. V. Zhemchugov

A. I. Alikhanov Institute for Theoretical and Experimental Physics Moscow, Russia

> July 18–30, 2016 Dubna, Russia

Institute for High-Energy Physics Protvino (Serpukhov), Russia OKA Detector

Current experiment  $E_K = 17.7 \text{ GeV}.$ 



Institute for High-Energy Physics Protvino (Serpukhov), Russia OKA Detector

Current experiment  $E_K = 17.7$  GeV.

[Yu. M. Antipov et. al., Phys. Rev. D36, 21 (1987)]  $E_{\pi} = 40$  GeV.



Massless QED Lagrangian:

$$\mathcal{L}_{\mathsf{QED}} = \bar{\psi}(ie\partial_{\mu} - eA)\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$

 $U(1) \times U(1)$  symmetry transformations:

$$\psi \to \mathrm{e}^{i\alpha}\psi, \qquad \psi \to \mathrm{e}^{i\beta\gamma^5}\psi$$

Noether currents:

 $\begin{aligned} j^{\mu} &= \bar{\psi} \gamma^{\mu} \psi, \qquad \partial_{\mu} j^{\mu} = 0, \\ j^{\mu}_{5} &= \bar{\psi} \gamma^{\mu} \gamma^{5} \psi, \quad \partial_{\mu} j^{\mu}_{5} = 0. \end{aligned}$  $\langle \gamma(k), \gamma(k') | \partial_{\mu} j_{5}^{\mu} | 0 \rangle = \partial_{\mu} \langle \gamma(k), \gamma(k') | \bar{\psi} \gamma^{\mu} \gamma^{5} \psi | 0 \rangle$  $=2\partial_{\mu} \ \mu \times$ 

$$\partial_{\mu}j_{5}^{\mu} = \frac{e^{2}}{16\pi^{2}}\varepsilon^{\mu\nu\alpha\beta}F_{\mu\nu}F_{\alpha\beta} \neq 0$$

QCD:

 $\pi^0 \to \gamma \gamma$ 

### $SU(3)_L \times SU(3)_R \to SU(3)_V$

Goldstone bosons  $\phi^a$ ,  $a = 1, \ldots, 8$ :

$$\Sigma = e^{\frac{2i}{F_{\pi}}\phi^{a}T^{a}}, \ F_{\pi} = 92.2 \text{ MeV}, \ \phi^{a}T^{a} = \begin{pmatrix} \frac{\pi^{0}}{\sqrt{2}} + \frac{\eta}{\sqrt{6}} & \pi^{+} & K^{+} \\ \frac{\pi^{+}}{K^{+}} & -\frac{\pi^{0}}{\sqrt{2}} + \frac{\eta}{\sqrt{6}} & K^{0} \\ \frac{\pi^{+}}{K^{+}} & K^{0} & -\frac{2\eta}{\sqrt{6}} \end{pmatrix}$$

Chiral perturbation theory:

$$\mathcal{L}_{\mathsf{ChPT}} = \frac{F_{\pi}^2}{8} \operatorname{tr}(D_{\mu}\Sigma D^{\mu}\Sigma^{\dagger}) + \mathcal{L}_{\mathsf{WZ}}$$
$$D_{\mu}\Sigma = \partial_{\mu}\Sigma + ie[A_{\mu}, Q], \quad Q = \begin{pmatrix} \frac{2}{3} & 0 & 0\\ 0 & -\frac{1}{3} & 0\\ 0 & 0 & -\frac{1}{3} \end{pmatrix}$$

Wess-Zumino term (one of):

$$\mathcal{L}_{\mathsf{WZ}} \supset \frac{1}{2\pi^2 F_{\pi}^5} \varepsilon^{\mu\nu\rho\sigma} \operatorname{tr}(\Sigma \partial_{\mu} \Sigma \partial_{\nu} \Sigma \partial_{\rho} \Sigma \partial_{\sigma} \Sigma) \supset \frac{e^2}{8\pi^2 F_{\pi}^5} \varepsilon^{\mu\nu\rho\sigma} \pi^0 F_{\mu\nu} F_{\rho\sigma}$$

$$\mathcal{L}_{\mathsf{ChPT}} = \frac{F_{\pi}^2}{8} \operatorname{tr}(D_{\mu} \Sigma D^{\mu} \Sigma^{\dagger} + \chi \Sigma^{\dagger} + \Sigma \chi^{\dagger}) + \mathcal{L}_{\mathsf{WZ}}$$



$$\mathcal{L}_{\mathsf{ChPT}} = \frac{F_{\pi}^2}{8} \operatorname{tr}(D_{\mu} \Sigma D^{\mu} \Sigma^{\dagger} + \chi \Sigma^{\dagger} + \Sigma \chi^{\dagger}) + \mathcal{L}_{\mathsf{WZ}}$$



#### Wrong:

- ► No quarks in ChPT.
- ► No mesons in QCD.

$$\begin{split} \Delta \mathcal{L}_{\text{QCD}} &= \bar{Q}i \not \!\!\!D Q + \frac{1}{\Lambda^2} (\bar{q}_R Q_L) (\bar{Q}_R q_L) + \text{h.c.}, \\ q &= \begin{pmatrix} u \\ d \\ s \end{pmatrix}, \ Q = \begin{pmatrix} U \\ D \\ S \end{pmatrix}, \ \Lambda \to \infty \\ \mathcal{L}_{\text{ChPT}} &= \frac{F_\pi^2}{8} \operatorname{tr} (D_\mu \Sigma D^\mu \Sigma^\dagger + \chi \Sigma^\dagger + \Sigma \chi^\dagger) + \bar{Q}i \not \!\!\!D Q \\ &+ M \bar{Q}_R \Sigma Q_L + \text{h.c.} + \mathcal{L}_{\text{WZ}}, \quad M \to \infty \end{split}$$





#### Wrong:

- ▶ No quarks in ChPT.
- No mesons in QCD.

#### [H. Georgi, "Weak Interactions", Ch. 6a.3]

$$\pi^+ \gamma \to \pi^+ \pi^0$$

Another Wess-Zumino term:





$$\begin{aligned} A(\pi^-\gamma \to \pi^-\pi^0) &= h(s,t,u) \cdot \varepsilon^{\mu\alpha\beta\gamma} A_\mu \partial_\alpha \pi^- \partial_\beta \pi^+ \partial_\gamma \pi^0 \\ h(s,t,u) &= h(0) \left\{ 1 + \frac{2f_{\rho\pi\pi} f_{\rho\pi\gamma}}{m_\rho^2 h(0)} \left[ \frac{s}{m_\rho^2 - s} + \frac{t}{m_\rho^2 - t} + \frac{u}{m_\rho^2 - u} \right] + \frac{f_{\omega\gamma} f_{\omega3\pi}}{m_\omega^2 h(0)} \frac{q^2}{m_\omega^2 - q^2} \right\} \end{aligned}$$

h(0) values

$$\begin{array}{l} \mbox{Theory} \ \ \frac{e}{4\pi^2 F_\pi^3} = 9.8 \ \mbox{GeV}^{-3} \\ \mbox{Experiment at LO (1987)} \ \ 12.9 \pm 0.9 \ \ (\mbox{stat.}) \pm 0.5 \ \ (\mbox{syst.}) \pm 1.0 \ \ (\mbox{sign}) \ \ \mbox{GeV}^{-3} \\ \mbox{Experiment at NNLO} + \ \mbox{EMC (2001)} \ \ 10.7 \pm 1.2 \ \ \mbox{GeV}^{-3} \end{array}$$

Update from the COMPASS Collaboration?













 $\pi^+ \gamma \to \pi^+ \pi^0$ 





 $K^+\gamma \to K^+\pi^0$ 



neutral pion production

 $K^+\gamma \to K^0\pi^+$ 



charged pion production





s-channel amplitude:

$$A_{s}^{(0)}(K^{+}\gamma \to K^{+}\pi^{0}) = -\frac{2f_{K^{*}+K^{+}\gamma}f_{K^{*}+K^{+}\pi^{0}}}{s - m_{K^{*}+}^{2} + i\sqrt{s}\Gamma_{K^{*}+}(s)}\varepsilon^{\alpha\beta\gamma\delta}\epsilon_{\alpha}p_{\beta}k_{1\gamma}k_{2\delta}$$
$$A_{s}(K^{+}\gamma \to K^{+}\pi^{0}) = A_{s}^{(0)}(K^{+}\gamma \to K^{+}\pi^{0}) - A_{s}^{(0)}(K^{+}\gamma \to K^{+}\pi^{0})|_{s=0}$$

Cross section:

$$\begin{split} \frac{d\sigma(K^+\gamma \to K^+\pi^0)}{dt} &= \frac{1}{2^7\pi} \left( t + \frac{(st - m_{K^+}^2 m_{\pi^0}^2)(t - m_{\pi^0}^2)}{(s - m_{K^+}^2)^2} \right) \\ &\times \left| \frac{e}{4\pi^2 F_\pi^3} + \frac{2f_{K^*+K^+\gamma} f_{K^*+K^+\pi^0}}{m_{K^{*+}}^2 - s - i\sqrt{s} \Gamma_{K^{*+}}(s)} \cdot \frac{s}{m_{K^{*+}}^2} \right. \\ &+ \frac{2f_{K^*+K^+\gamma} f_{K^{*+}K^+\pi^0}}{m_{K^{*+}}^2 - u} \cdot \frac{u}{m_{K^{*+}}^2} + \frac{2f_{\rho^0\pi^0\gamma} f_{\rho^0K^+K^+}}{m_{\rho^0}^2 - t} \cdot \frac{t}{m_{\rho^0}^2} \\ &+ \frac{2f_{\omega\pi^0\gamma} f_{\omega K^+K^+}}{m_{\omega}^2 - t} \cdot \frac{t}{m_{\omega}^2} + \frac{2f_{\phi\pi^0\gamma} f_{\phi K^+K^+}}{m_{\phi}^2 - t} \cdot \frac{t}{m_{\phi}^2} \right|^2 \end{split}$$

E. V. Zhemchugov

$$\begin{array}{ll} f_{K^{*+}K^{+}\pi^{0}} &=& 3.10 \\ f_{K^{*+}K^{0}\pi^{+}} &=& 4.38 \\ f_{K^{*0}K^{+}\pi^{+}} &=& 4.41 \\ f_{\rho^{0}K^{+}K^{+}} &=& 3.16 \\ f_{\rho^{+}K^{+}K^{0}} &=& -4.47 \\ f_{\omega K^{+}K^{+}} &=& -4.47 \\ f_{K^{*+}K^{+}\gamma} &=& 0.240 \ {\rm GeV^{-1}} \\ f_{K^{*0}K^{0}\gamma} &=& -0.385 \ {\rm GeV^{-1}} \\ f_{\rho^{0}\pi^{0}\gamma} &=& 0.252 \ {\rm GeV^{-1}} \\ f_{\rho^{+}\pi^{+}\gamma} &=& 0.219 \ {\rm GeV^{-1}} \\ f_{\omega\pi^{0}\gamma} &=& 0.696 \ {\rm GeV^{-1}} \\ |f_{\phi\pi^{0}\gamma}| &=& 0.040 \ {\rm GeV^{-1}} \end{array}$$

Decay widths:

$$\begin{split} \Gamma(K^* \to K\pi) &\implies |f_{K^*K\pi}| \\ \Gamma(K^* \to K\gamma) &\implies |f_{K^*K\gamma}| \\ \Gamma(\phi \to K^+K^-) &\implies |f_{\phi K^+K^+}| \\ \Gamma(\rho^+ \to \pi^+\gamma) &\implies |f_{\rho^+\pi^+\gamma}| \\ \Gamma(\rho^0 \to \pi^0\gamma) &\implies |f_{\rho^0\pi^0\gamma}| \\ \Gamma(\omega \to \pi^0\gamma) &\implies |f_{\omega\pi^0\gamma}| \\ \Gamma(\phi \to \pi^0\gamma) &\implies |f_{\phi\pi^0\gamma}| \end{split}$$

SU(3) symmetry:

$$\begin{split} \sqrt{2}f_{K^{*+}K^{+}\pi^{0}} &= f_{K^{*+}K^{0}\pi^{+}} = f_{K^{*0}K^{+}\pi^{+}} = -f_{\rho^{+}K^{+}K^{0}} \\ &= \sqrt{2}f_{\rho^{0}K^{+}K^{+}} = \sqrt{2}f_{\omega K^{+}K^{+}} = -f_{\phi K^{+}K^{+}} \end{split}$$

$$f_{K^{*+}K^+\gamma} = f_{\rho^+\pi^+\gamma} = f_{\rho^0\pi^0\gamma} = \frac{1}{3}f_{\omega\pi^0\gamma} = -\frac{1}{2}f_{K^{*0}K^0\gamma}$$

The sign of the anomaly term is unknown.

E. V. Zhemchugov



Weizsacker-Williams equivalent photons approximation:

$$\frac{d\sigma(K^+N \to K\pi N)}{dt \, ds \, dq_{\perp}^2} = \frac{Z^2 \alpha}{\pi (s - m_{K^+}^2)} \frac{q_{\perp}^2}{\left(q_{\perp}^2 + \left(\frac{s - m_{K^+}^2}{2E_K}\right)^2\right)^2} \frac{d\sigma(K^+\gamma \to K\pi)}{dt} |F(\vec{q}\,^2)|^2}{F(\vec{q}\,^2)} \\ F(\vec{q}\,^2) = \exp\left(-\frac{\langle r^2 \rangle \vec{q}\,^2}{6}\right) \\ \frac{d\sigma(K^+N \to K\pi N)}{dt \, ds} = \frac{Z^2 \alpha}{\pi} \frac{E_1(a) - 1}{s - m_{K^+}^2} \frac{d\sigma(K^+\gamma \to K\pi)}{dt} \\ E_1(a) = \int_a^\infty \frac{e^{-z}}{z} dz, \ a = \frac{1}{3}r_0^2 A^{2/3} \left(\frac{s - m_{K^+}^2}{2E_K}\right)^2$$

[Berestetskiy, Lifshitz, Pitaevsky, "Quantum Electrodynamics", §99]

E. V. Zhemchugov





## Conclusions

- ► A theoretical prediction has been made for the cross sections of  $K^+\gamma \rightarrow K^+\pi^0$ and  $K^+\gamma \rightarrow K^0\pi^+$  reactions at low energies. For the anomalous reaction, we predict two possible values depending on the a priori unknown sign of the interference term, which should be resolved by the experiment.
- ▶ It is possible to observe the chiral anomaly through comparison of cross section of  $K^+$  Cu  $\rightarrow K^+\pi^0$  Cu reaction with that of  $K^+$  Cu  $\rightarrow K^0\pi^+$  Cu reaction at  $s \leq 0.6$  GeV<sup>2</sup>. The point is that only the first one has the anomaly which manifests itself as an increase in the cross section at low s.
- Luminosity of 60  $\mu$ b<sup>-1</sup> at 0.4 <  $s < 0.6 \text{ GeV}^2$  is planned to be collected in the Protvino experiment. In this case expected observations are  $\approx 10$  events of  $K^0\pi^+$  production and either  $\approx 20 \text{ or } \approx 70$  events of  $K^+\pi^0$  production, depending on the sign of the interference term.

# Thank you for your attention!















Total: 1

 $K^+\gamma \to K^+\pi^0$ 









Total: 0