

Dark Matter Models (II)

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Outline

- 1 Nonthermal production mechanisms
 - Gravitino
 - Classical scalar field
 - Axion
 - Sterile neutrino
- 2 Guiding principles. . .
- 3 Summary

Outline

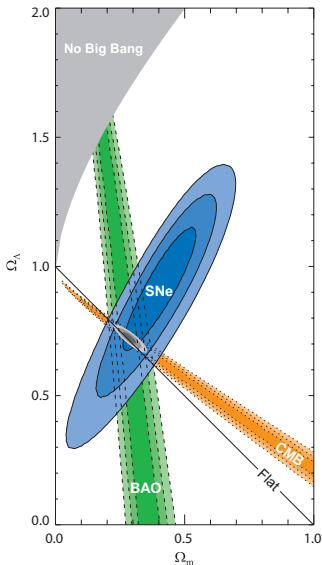
1 Nonthermal production mechanisms

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Astrophysical and cosmological data are in agreement



$$\left(\frac{\dot{a}}{a}\right)^2 = H^2(t) = \frac{8\pi}{3} G \rho_{\text{density}}^{\text{energy}}$$

$$\rho_{\text{density}}^{\text{energy}} = \rho_{\text{radiation}} + \rho_{\text{matter}}^{\text{ordinary}} + \rho_{\text{matter}}^{\text{dark}} + \rho_{\Lambda}$$

$$\rho_{\text{radiation}} \propto 1/a^4(t), \quad \rho_{\text{matter}} \propto 1/a^3(t), \quad \rho_{\Lambda} = \text{const}$$

$$\frac{3H_0^2}{8\pi G} = \rho_{\text{density}}^{\text{energy}}(t_0) \equiv \rho_c \approx 0.53 \times 10^{-5} \frac{\text{GeV}}{\text{cm}^3}$$

radiation:

$$\Omega_{\gamma} \equiv \frac{\rho_{\gamma}}{\rho_c} = 0.5 \times 10^{-4}$$

Baryons (H, He):

$$\Omega_B \equiv \frac{\rho_B}{\rho_c} = 0.05$$

Neutrino:

$$\Omega_{\nu} \equiv \frac{\sum \rho_{\nu_i}}{\rho_c} < 0.01$$

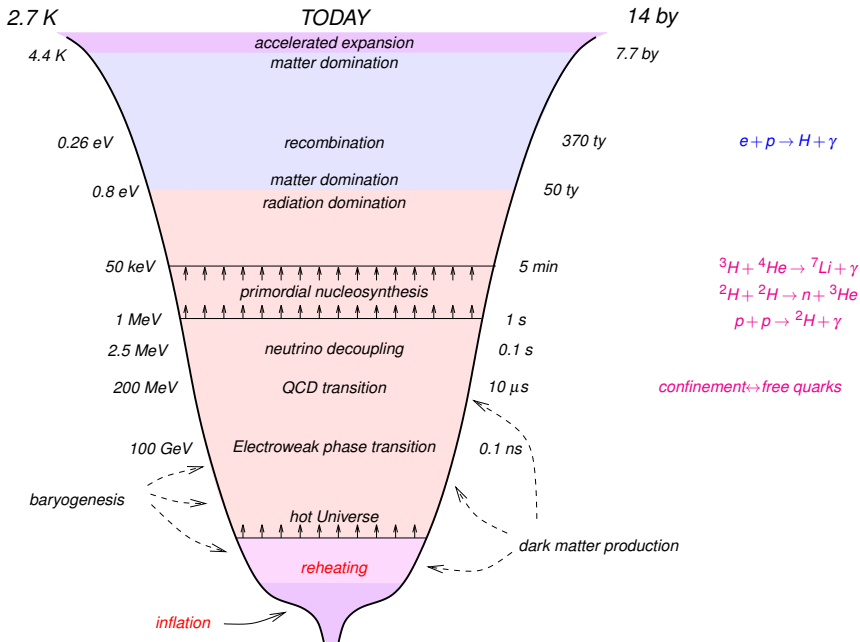
$$N_{\nu} \simeq 3, \quad \sum m_{\nu} \lesssim 0.2 \text{ eV}$$

Dark matter:

$$\Omega_{\text{DM}} \equiv \frac{\rho_{\text{DM}}}{\rho_c} = 0.27$$

Dark energy:

$$\Omega_{\Lambda} \equiv \frac{\rho_{\Lambda}}{\rho_c} = 0.68$$



Dark Matter: non-thermal production

- ① in the primordial plasma of SM particles
(via scatterings, oscillations):

gravitino
sterile neutrino of 1-50 keV

- ② at phase transitions:

axion of $10^{-4} - 10^{-7}$ eV
Q-balls
strangelets (?)

- ③ during reheating (after inflation?):

black holes
any guy coupled (only) to inflaton

- ▶ perturbatively:

inflaton decays
production by external (inflaton) field

- ▶ non-perturbatively:

Bose-enhancement of
coherent production by external field

- ④ while the Universe expands:

gravity produces any particles at $H \sim M_X$

Gravitino production

$$\mathcal{L} = \frac{1}{F} \partial^\mu \psi \cdot J_\mu^{SUSY}, \quad \tilde{G}_\mu \rightarrow \tilde{G}_\mu + i\sqrt{4\pi} \frac{M_{Pl}}{F} \partial_\mu \psi$$

$$m_{3/2} = \sqrt{\frac{8\pi}{3}} \frac{F}{M_{Pl}} \longleftrightarrow \Lambda = 0$$

$$1 \text{ TeV} \lesssim \sqrt{F} \lesssim M_{Pl}, \quad 2 \cdot 10^{-4} \text{ eV} \lesssim m_{3/2} \lesssim M_{Pl}$$

LSP in low scale SUSY breaking models

$$2 \cdot 10^{-4} \text{ eV} \lesssim m_{3/2} \lesssim 100 \text{ GeV} \longrightarrow \sqrt{F} \lesssim 10^{10} \text{ GeV}$$

Thermal equilibrium is forbidden

(fermion; would be hot DM):

$$\Omega_{3/2} = \frac{m_{3/2} \cdot n_{3/2}}{\rho_c} = 0.2 \left(\frac{m_{3/2}}{200 \text{ eV}} \right) \left(\frac{g_{3/2}}{2} \right) \left(\frac{210}{g_*(T_d)} \right) \frac{1}{2h^2}$$

Gravitino production in scatterings and decays

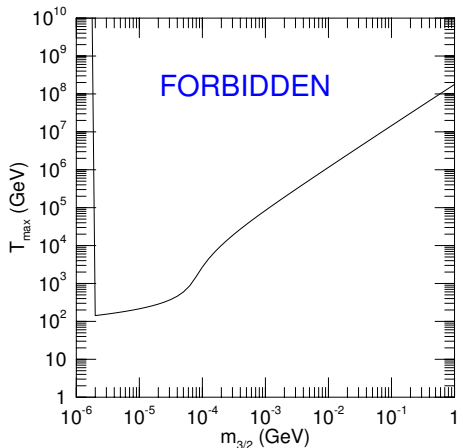
$$\tilde{X}_i \rightarrow \tilde{G} + X_i, \quad X_i + X_j \rightarrow \tilde{X}_k + \tilde{G}$$

$$\Gamma \propto \frac{1}{F^2} \propto \frac{1}{m_{3/2}^2}, \quad \sigma \propto \frac{1}{F^2} \propto \frac{1}{m_{3/2}^2}$$

$$\begin{aligned} \frac{dn_{3/2}}{dt} + 3Hn_{3/2} &= \sum_i \Gamma_{\tilde{X}_i} \cdot \gamma_i^{-1} \cdot n_{\tilde{X}_i} + \sum_{i,j} \langle \sigma_{ij} \rangle \cdot n_{X_i} n_{X_j}, \end{aligned}$$

$$\begin{aligned} \frac{d}{dT} \left(\frac{n_{3/2}}{s} \right) &= - \sum_i \Gamma_{\tilde{X}_i} \cdot \frac{n_{\tilde{X}_i}}{\gamma_i \cdot sHT} - \sum_{i,j} \frac{\langle \sigma_{ij} \rangle \cdot n_{X_i} n_{X_j}}{sHT}, \\ &\propto \frac{1}{T^3} \quad \propto \text{const} \end{aligned}$$

$$\begin{aligned} \Omega_{3/2} &\sim \left(\frac{200 \text{ keV}}{m_{3/2}} \right) \cdot \left(\frac{T_{max}}{10 \text{ TeV}} \right) \\ &\times \left(\frac{M_S}{1 \text{ TeV}} \right)^2 \cdot \left(\frac{15}{\sqrt{g_*(T_{max})}} \right) \cdot \frac{1}{2h^2} \end{aligned}$$



... issues of QFT description
of a gauge theory at finite temperature

Outcome depends on initial conditions !!!

Free massive scalar field

$$\mathcal{L} = \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{2} m_\phi^2 \phi^2$$

For the homogeneous scalar field in FLRW expanding Universe

$$\ddot{\phi} + 3H\dot{\phi} + m_\phi^2 \phi = 0$$

we find two-stage evolution:

$$m_\phi < H(t) \implies \phi = \phi_i = \text{const}$$

$$m_\phi > H(t) \implies \rho = \langle E_k \rangle - \langle E_p \rangle = 0, \quad \rho \sim m_\phi^2 \phi^2 \propto 1/a^3$$

- **dust-like substance** in the late Universe, $\Omega \propto m_\phi^{1/2} \phi_i^2$
depends on initial conditions
- **pressureless** at spatial scales $l > 1/m_\phi$ fuzzy DM
- may help (?) with CDM-problems (core-cusp, lack of dwarfs, etc)

Axion: well-motivated but fine-tuned

$$L_\theta = \frac{\alpha_s}{8\pi} \left(\theta_0 + \text{Arg}(\text{Det}\hat{M}_q) \right) G_{\mu\nu}^a \tilde{G}^{\mu\nu a} \equiv \frac{\alpha_s}{8\pi} \theta G_{\mu\nu}^a \tilde{G}^{\mu\nu a} \equiv \partial_\mu K^\mu$$

P-CP-violation

tree-level and $U(1)_A$ -anomaly contributions, $\bar{q}_L \hat{M}_q q_R + h.c$

strong *CP*-problem

Theory and Nature: neutron EDM $\theta < 10^{-9}$

nonantropic parameter!

Transformation
(PQ-symmetry)

$$q_L^k \rightarrow e^{ie_k^{(PQ)}\beta/2} q_L^k$$

$$q_R^k \rightarrow e^{-ie_k^{(PQ)}\beta/2} q_R^k,$$

with $\sum_k e_k^{(PQ)} \neq 0$

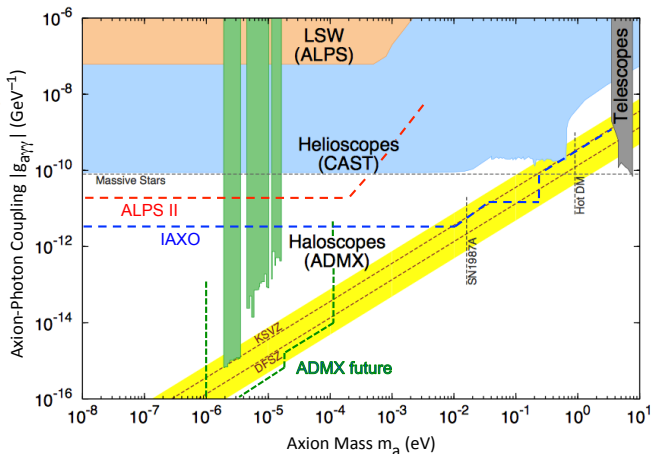
cancels θ with β

$$\theta \rightarrow \theta + a(x)/f_a$$

$$m_{\text{axion}} \simeq f_\pi m_\pi / f_a$$

Dark Matter region

$$\mathcal{L} \sim g_{a\gamma\gamma} \times a(x) F_{\mu\nu} F^{\mu\nu}$$



Axion as Cold Dark Matter

$$\theta \rightarrow \bar{\theta}(x) = \theta + C_g \frac{a(x)}{f_{PQ}}.$$

Free scalar field starts to evolve at $m \simeq H$

$$\mathcal{L} = \frac{f_{PQ}^2}{2} \cdot \left(\frac{d\bar{\theta}}{dt} \right)^2 - \frac{m_a^2(T)}{2} f_{PQ}^2 \bar{\theta}^2,$$

Temperature-dependent mass-term

$$m_a(T) \simeq 0.1 \cdot m_a(0) \cdot \left(\frac{\Lambda_{QCD}}{T} \right)^{3.7}, \quad T > \Lambda_{QCD}$$

$$\Omega_a \simeq 0.2 \cdot \bar{\theta}_i^2 \cdot \left(\frac{4 \cdot 10^{-6} \text{ eV}}{m_a} \right)^{1.2} \cdot \frac{1}{2h^2}$$

- initial conditions
- instantons vs lattice
- domain walls, axion clumps, etc

Sterile neutrino: well-motivated Dark Matter

$$\mathcal{L}_N = \bar{N}_I i \not{\partial} N_I - f_{\alpha I} \bar{L}_\alpha \tilde{H} N_I - \frac{M_{N_I}}{2} \bar{N}_I^c N_I + \text{h.c.}$$

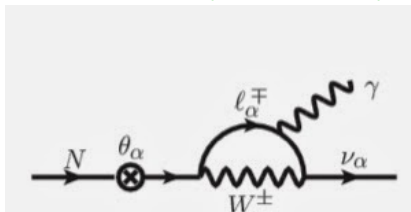
- massive fermions giving mass to active neutrino through mixing

$$m_a \sim \frac{f^2 v^2}{M_N^2} M_N \sim \theta^2 M_N$$

- unstable, but exceeding the age of the Universe at condition

$$\theta^2 < 1.5 \times 10^{-7} \left(\frac{50 \text{ keV}}{M_N} \right)^5$$

- can be searched for because of two-body radiative decay



Production in oscillations

$$\frac{\partial}{\partial t} f_s - H \mathbf{p} \frac{\partial}{\partial \mathbf{p}} f_s = \Gamma_\alpha P(\nu_\alpha \rightarrow \nu_s) f_\alpha(t, \mathbf{p}).$$

where $\Gamma_\alpha \sim G_F^2 T^4 E$ is the **weak interaction** rate in plasma

$$P(\nu_\alpha \rightarrow \nu_s) = \sin^2 2\theta_\alpha^{\text{mat}} \cdot \sin^2 \left(\frac{t}{2t_\alpha^{\text{mat}}} \right),$$

$$t_\alpha^{\text{mat}} = \frac{t_\alpha^{\text{vac}}}{\sqrt{\sin^2 2\theta_\alpha + (\cos 2\theta_\alpha - V_{\alpha\alpha} \cdot t_\alpha^{\text{vac}})^2}},$$

$$\sin 2\theta_\alpha^{\text{mat}} = \frac{t_\alpha^{\text{mat}}}{t_\alpha^{\text{vac}}} \cdot \sin 2\theta_\alpha, \quad t_\alpha^{\text{vac}} = \frac{2E}{M_N^2}$$

and **effective plasma potential** for active neutrinos

$$V_{\alpha\alpha} \sim -\# G_F^2 T^4 E + \# G_F T^2 \mu_{L\alpha}$$

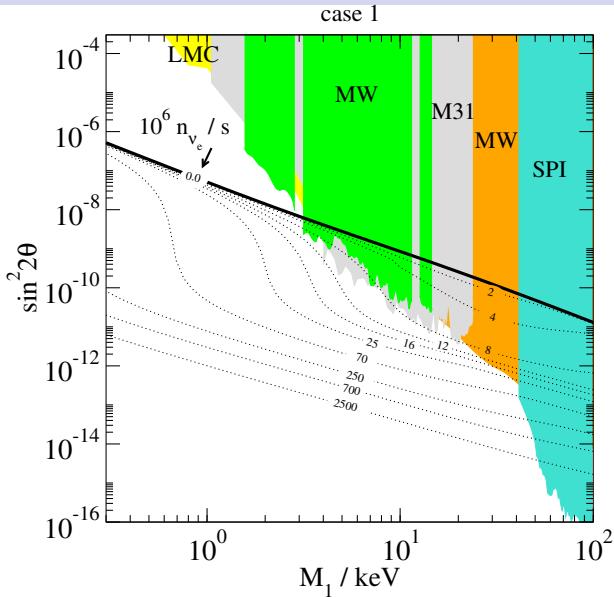
resonant production in the lepton asymmetric plasma

BAU-DM connection?

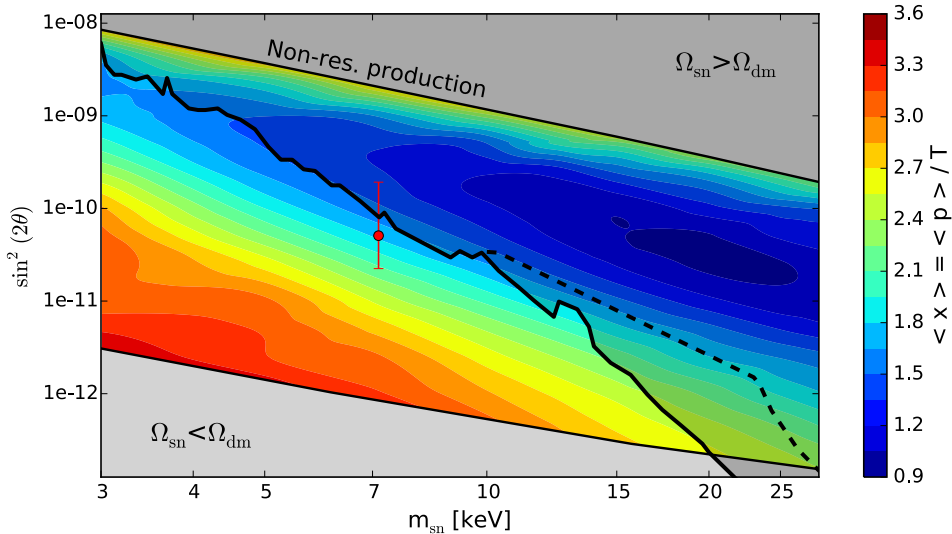
Sterile neutrino Dark Matter

$\sin^2 2\theta$, n_{ν_e} and M_N
to saturate Ω_{DM}

- larger asymmetry $10^6 n_{\nu_e}/s > 2500$ is forbidden by BBN
- above the solid line "0.0" $\Omega_N > \Omega_{DM}$
- selected upper limits from X-ray telescopes
- recall $m > 0.75$ keV for fermionic DM
It can be refined with estimates of neutrino velocities



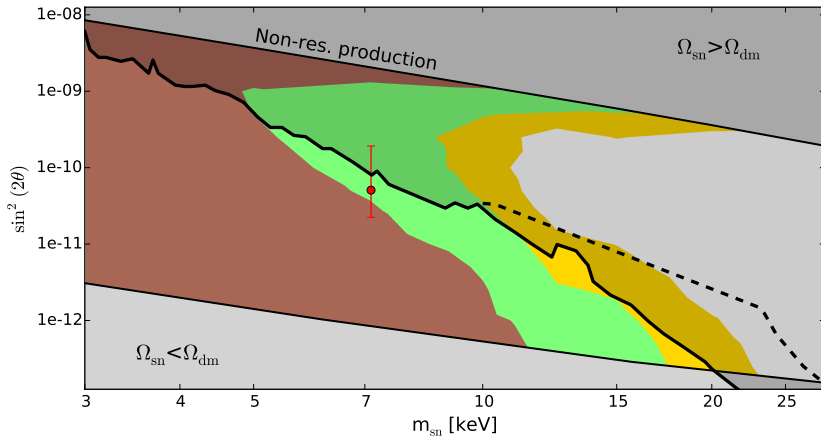
Sterile neutrino Dark Matter



A.Schneider (2016)

Sterile neutrino Dark Matter: ... gone?

A.Schneider (2016)



brown: MW satellite counts

green and yellow: Lyman- α

production by inflaton

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Dark Matter: possible guiding principles

Naturality:

- exploit known interactions

examples: WIMPs,
free particles

- part of a well-motivated model

examples: LSP, axion,
sterile neutrinos

- Why $\Omega_B \sim \Omega_{DM}$?

examples: Mirror World
antibaryonic DM

- Why $\Omega_\Lambda \sim \Omega_{DM}$?

examples:
DM-DE coupling

Minimality:

Use as little new physics as possible

Motivation:

No any hints of
new physics in experiment

Many models are

untestable

example:

gravitationally produced

free massive fermion

Reality:

Deep insight into the gravitational properties of dark matter

what happen

at small scales?

status of:

cusps/core in galactic centers
lack of dwarf galaxies
lack of small galaxies

examples:

cold dark matter
warm dark matter
selfinteracting dark matter

Examples: both Natural and Minimal

Natural source of dark matter production: gravity

Gravity produces any free massive particle when metric changes in the expanding Universe

most efficiently when $H \sim M$

say, at radiation domination stage

$$\Omega_X \sim \left(\frac{M_X}{10^9 \text{ GeV}} \right)^{5/2}$$

S.Mamaev, V.Mostepanenko, A.Starobinsky (1976)

Modified gravity ($R \rightarrow R - R^2/6\mu^2$) may be responsible for inflation and subsequent reheating

A.Starobinsky (1980)

that is (universal) production of all particles, including those of dark matter

$$\Omega_X \simeq 0.15 \times \left(\frac{M_X}{10^7 \text{ GeV}} \right)^3$$

D.Gorbunov, A.Panin (2010)

Untestable

Minimal, but still **testable**: scalar Dark Matter

$$V_S = \frac{1}{2}\mu_S^2 S^2 + \frac{1}{2}\lambda_{hS} S^2 H^\dagger H$$

$$m_S = \sqrt{\mu_S^2 + \frac{1}{2}\lambda_{hS} v^2}$$

$$\Omega_S \propto m_S n_S \propto \frac{1}{\sigma_{ann}} \propto \frac{m_S^2}{\lambda_{hS}^2}$$

indirect:

$$\text{flux}(SS \rightarrow SM) \propto n_S^2 \sigma_{ann} \propto \frac{1}{\lambda_{hS}^2}$$

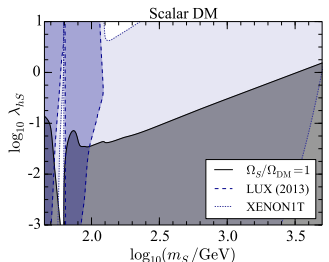
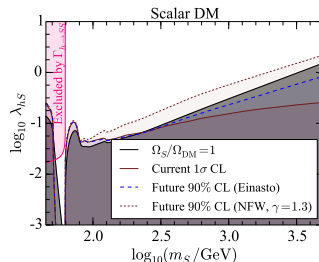
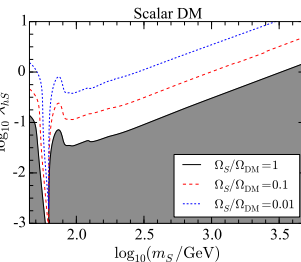
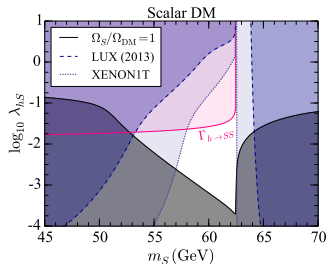
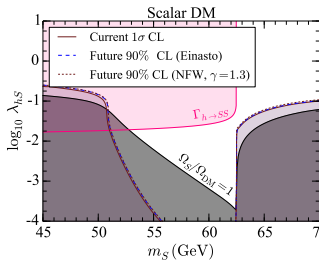
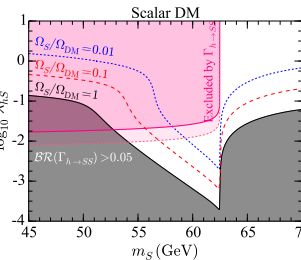
direct:

$$\Gamma(SA \rightarrow SA) \propto n_S \sigma_{ann} \propto \frac{1}{m_S}$$

- EW phase transition of 1 order ?
- EW vacuum stability ?

Constraints on scalar Dark Matter

A.Beniwal et al (2015)



Discussion on WIMPs

Most natural properties:

- to be in equilibrium in primordial plasma up to very freezout (and in kinetic equilibrium even later)
- to form a symmetric component:

$$X = \bar{X} \quad \text{or} \quad n_X = n_{\bar{X}}$$

But what we have in reality?

- We are sure there were
 - ▶ Big Bang Nucleosynthesis (starting from 1 MeV)
 - ▶ Recombination (at about 0.3 eV)
 and both are significantly “out-of-equilibrium” processes
- The visible matter is asymmetric, so that

$$f \neq \bar{f} \quad \text{and} \quad n_f \neq n_{\bar{f}}$$

Asymmetric Dark Matter

Sakharov's conditions are involved !!!

- Two different quantum numbers work in two different (dark and visible) sectors

Then a mechanism similar to baryogenesis is responsible for DM production

- A single quantum number is responsible for both asymmetries

- ▶ Non-zero total asymmetry

- 1 Generation of the asymmetry
- 2 Redistribution of the asymmetry between two Worlds

- ▶ The Universe is neutral

- 1 Simultaneous production of particles in one sector and antiparticles in another
- 2 Annihilation of the symmetric parts of charged particles in the both sectors

Baryon or lepton numbers can be naturally exploited !

Different (as compared to WIMPs) phenomenological signatures

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Summary (II)

- A lot of models with nonthermal DM production
- Outcome generally depends on initial conditions
- Before BBN epoch everything is allowed
- DM may be multicomponent
- DM perturbations are adiabatic