



TWO PARTICLES BOSE-EINSTEIN CORRELATIONS AT 0.9 AND 7 TEV WITH THE ATLAS DETECTOR

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ON BEHALF OF ATLAS COLLABORATION



**“NEW TRENDS IN HIGH-ENERGY
PHYSICS”**

**JOINT INSTITUTE FOR NUCLEAR
RESEARCH CONFERENCE,
BUDVA, BECICI, MONTENEGRO,
2 – 8 OCTOBER, 2016**

MOTIVATION

- Bose-Einstein correlations (BEC) represent a unique probe of the *space-time geometry* of the *hadronization region* and allow the *determination the size and shape of the source* from which particles are *emitted*.
- Studies of the dependence of BEC on *particle multiplicity* and *transverse momentum* are of special interest. They help in the understanding of multiparticle production mechanisms.
- High-multiplicity data in proton interactions can serve as a reference for studies in nucleus-nucleus collisions. The effect is reproduced in both hydrodynamical/hydrokinetic and Pomeron-based approaches for hadronic interactions where high multiplicities play a crucial role.

PARAMETRIZATION MODELS

- **GSSg model.** *The Goldhaber spherical source model.*

$$\Omega^{(G)} = C_0 \left(1 + \lambda e^{-R^2 Q^2} \right) \cdot (1 + Q\varepsilon)$$

- **GSSe model.** *Empirical model. Used since it represents well the shape of the correlation.*

$$\Omega^{(G)} = C_0 \left(1 + \lambda e^{-RQ} \right) \cdot (1 + Q\varepsilon)$$

R is the source radius

λ is the *incoherence factor* (0, 1) introduced empirically.

- **QOg model.** *Quantum Optics model.*

$$\Omega^{(GO)} = C_0 \left(1 + 2p(1-p)e^{-R^2 Q^2} + p^2 e^{-2R^2 Q^2} \right) \cdot (1 + Q\varepsilon)$$

- **QOe model.** *Empirical but inspired to the Quantum Optics model.*

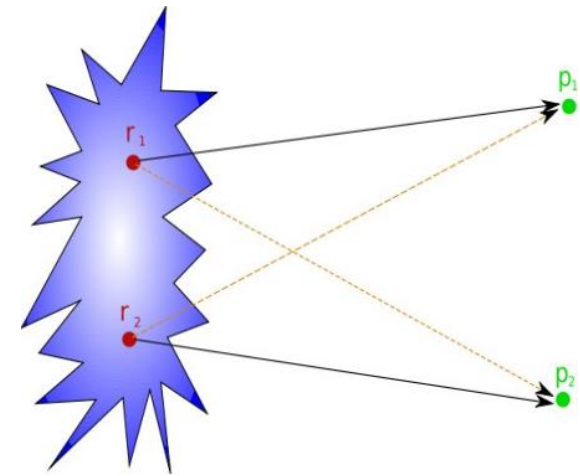
$$\Omega^{(EO)} = C_0 \left(1 + 2p(1-p)e^{-RQ} + p^2 e^{-2RQ} \right) \cdot (1 + Q\varepsilon)$$

p is the chaoticity: =0 (=1) for purely coherent (**chaotic**) sources.

BOSE-EINSTEIN CORRELATIONS

Correlations in phase space between two identical bosons from symmetry of wave functions.

- ▶ Enhances likelihood of two particles close in phase space
- ▶ Allows one to ‘probe’ the source of the bosons in size and shape
- ▶ Dependence on particle multiplicity and transverse momentum probes the production mechanism



Correlation function $C_2(Q)$ a ratio of probabilities:

$$C_2(Q) = \frac{\rho(p_1, p_2)}{\rho_0(p_1, p_2)} = C_0(1 + \Omega(\lambda, RQ)) \cdot (1 + Q\varepsilon), \quad Q^2 = -(p_1 - p_2)^2$$

$\Omega^G(\lambda, RQ) = \lambda e^{-R^2 Q^2}$

$\Omega^E(\lambda, RQ) = \lambda e^{-RQ}$

C_0 is a normalisation, ε accounts for long range effects, \mathbf{R} is the effective radius parameter of the source, λ is the strength of the effect parameter, 0/1 for coherent/chaotic source. Two possible parameterisation: Gaussian and Exponential.

$$C_2(Q) = \frac{N(p_1^\pm, p_2^\pm)(Q)}{N^{ref}(p_1, p_2)(Q)}$$

N^{ref} without BEC effect from: unlike-charge particles (UCP), opposite hemispheres, event mixing. **Basic Reference:** distribution of UCP pairs of non-identical particle taken from the same event.

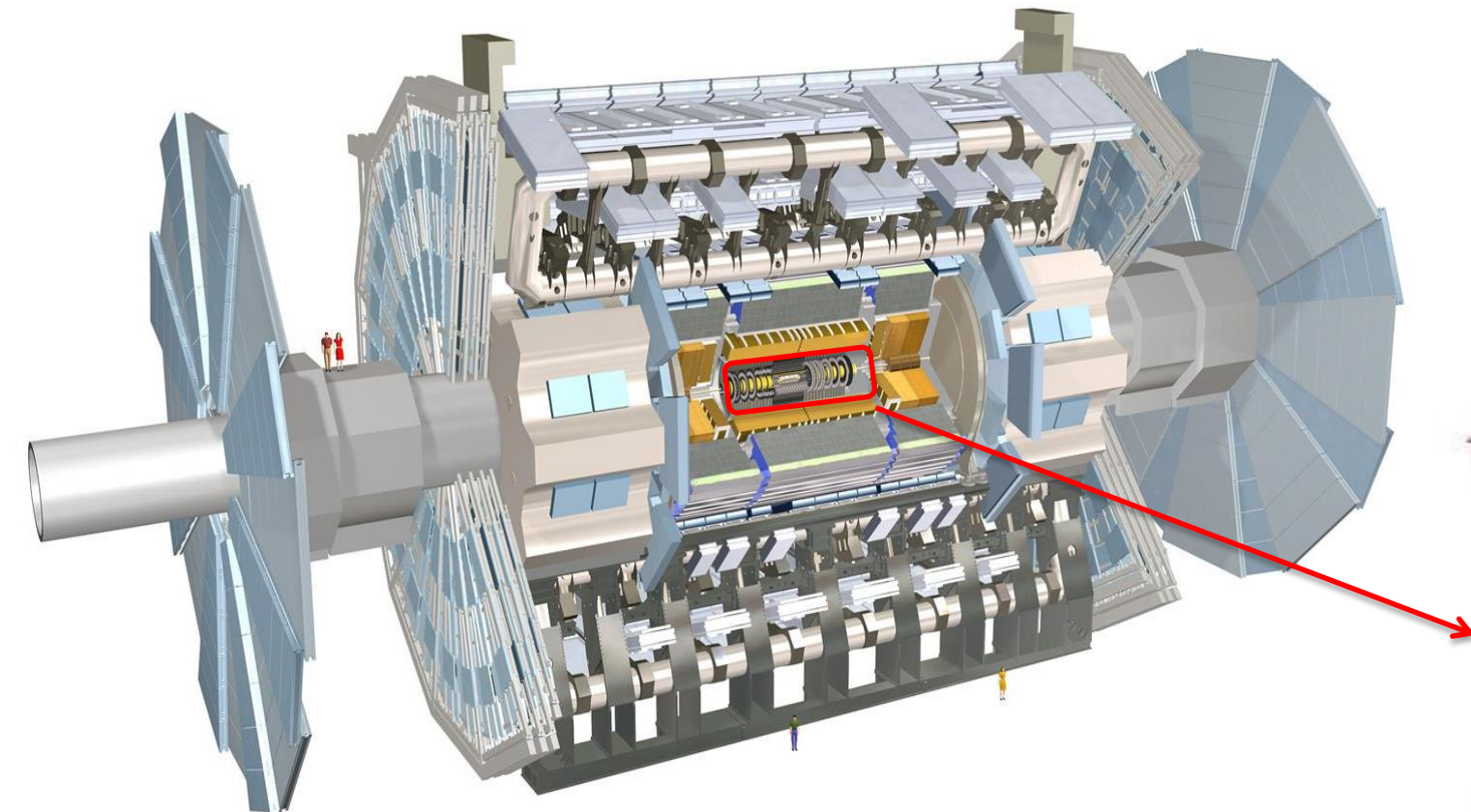
$$R_2(Q) = \frac{C_2^{Data}(Q)}{C_2^{MC}(Q)} = \frac{\rho^{Data}(p_1^\pm, p_2^\pm) / \rho_0^{Data}(p_1^\pm, p_2^\mp)}{\rho^{MC}(p_1^\pm, p_2^\pm) / \rho_0^{MC}(p_1^\pm, p_2^\mp)}$$

The studies are carried out using the **double ratio correlation function**. The $R_2(Q)$ eliminates problems with energy-momentum conservation, topology, resonances etc. **MC without BEC.**

A TOROIDAL LHC APPARATUS

The focus of ATLAS is high- p_T physics, and also provides a window onto important softer QCD processes. ► *Two-particle Bose-Einstein correlations in pp collisions at $\sqrt{s}=0.9$ and 7 TeV measured with the ATLAS detector.*

Published in EPJ C75 (2015) 10, 466

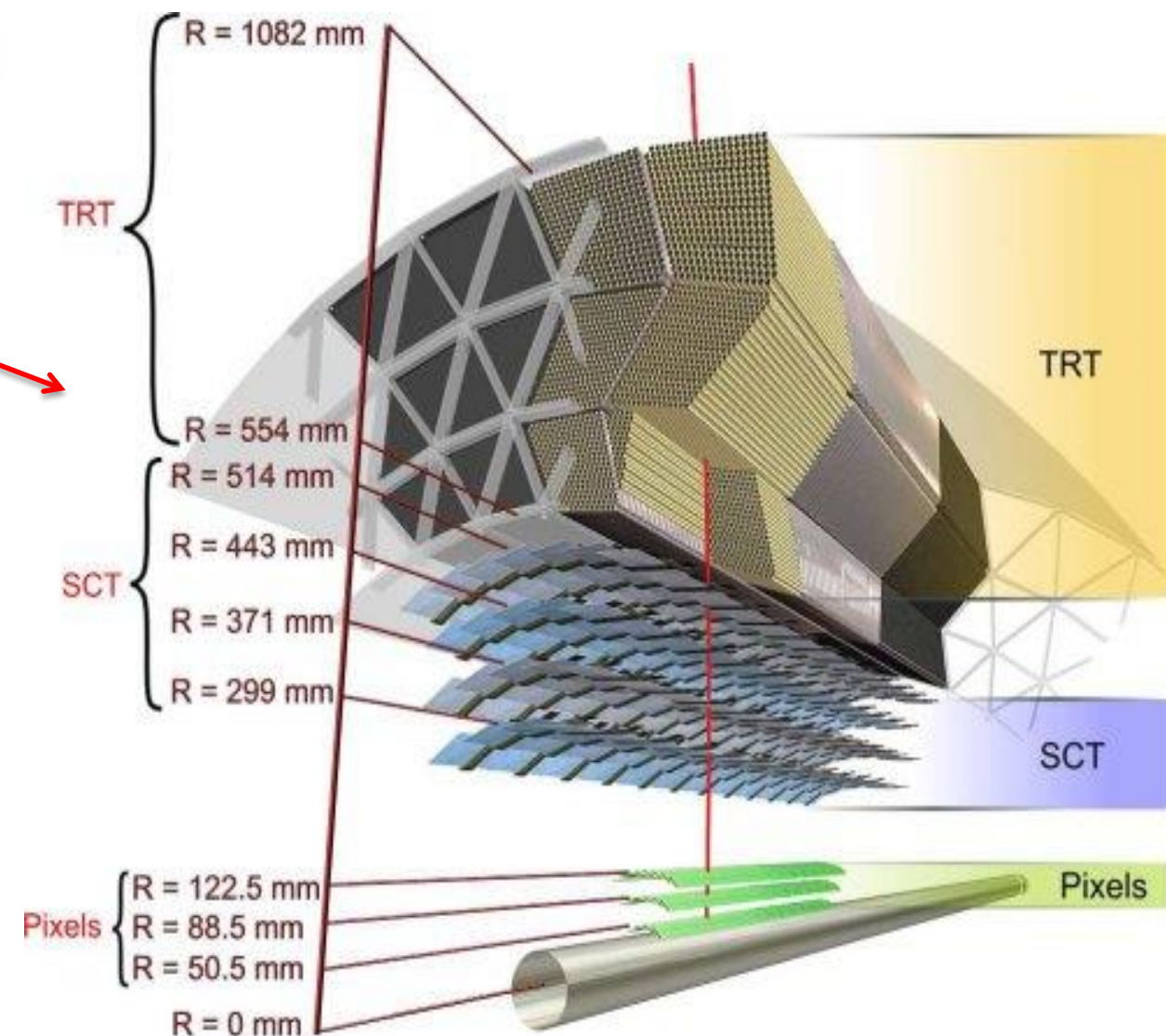


**ATLAS Inner Detector:
main tracking device (Run 1)**

Inner Detector ($|\eta| < 2.5$, $p_T > 100$ MeV):

Tracking; **2T Solenoid Magnet**

- Silicon Pixels $50 \times 400 \mu\text{m}^2$
- Silicon Strips (SCT) $40 \mu\text{m}$ rad stereo strips
- Transition Radiation Tracker (TRT) up to 36 points/track

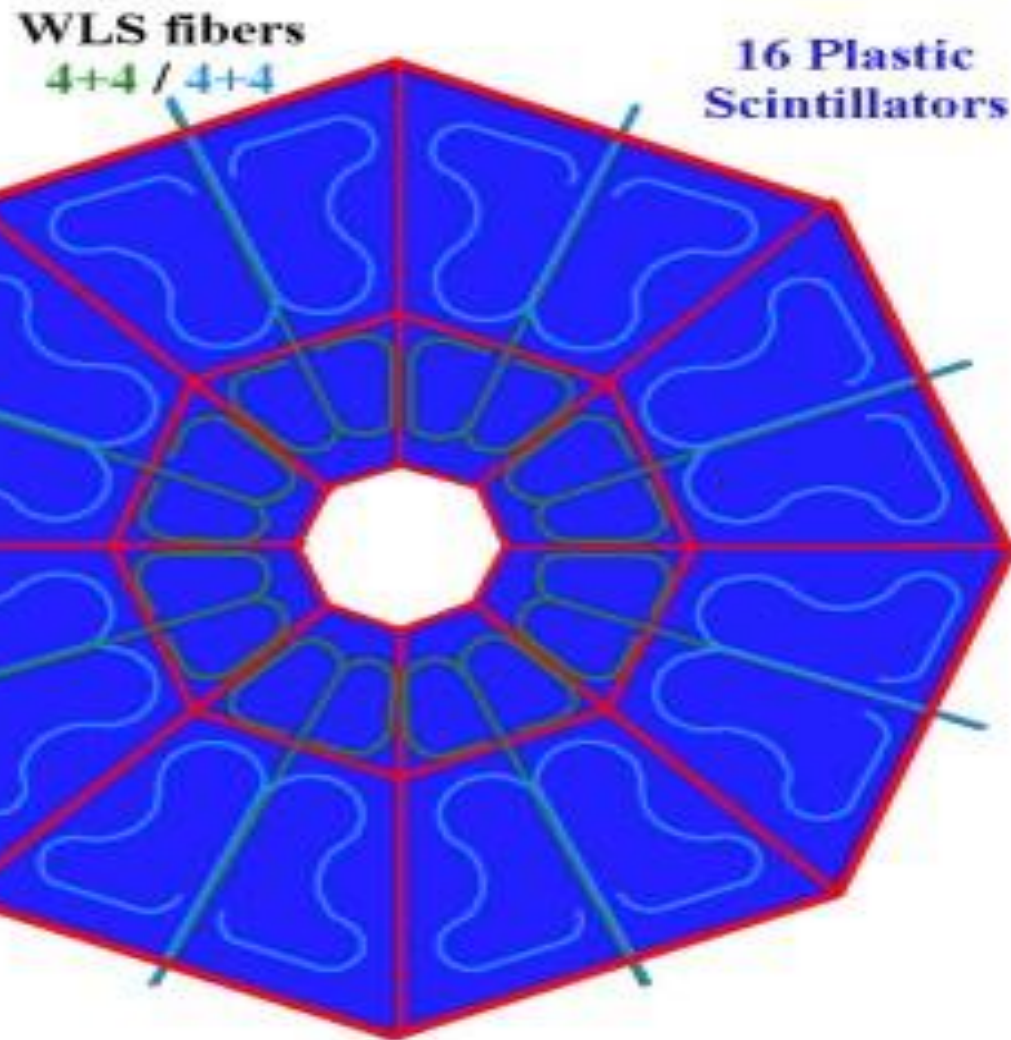


MINIMUM BIAS TRIGGER SCINTILLATOR (MBTS)

32 independent wedge-shaped plastic scintillators (16 per side) read out by PMTs, $2.09 < |\eta| < 3.84^*$

* Pseudorapidity is defined as $\eta = -\frac{1}{2} \ln(\tan(\theta/2))$, θ is the polar angle with respect to the beam.

Run 1



- Designed for triggering on min bias events, $>99\%$ efficiency
- **MBTS** timing used to veto halo and beam gas events
- Also being used as gap trigger for various diffractive subjects

MINIMUM-BIAS EVENT SELECTION CRITERIA

Events pass the data quality criteria (“**Good events**”: 1) *all ID sub-systems nominal cond.*, 2) *stable beam*, 3) *defined beam spot*)

- Accept on signal-arm Minimum Bias Trigger Scintillator
- **Primary vertex** (*2 tracks with $p_T > 100 \text{ MeV}$*),
- **Veto** to any additional vertices with ≥ 4 tracks.
- **At least 2 tracks with $p_T > 100 \text{ MeV}$, $|\eta| < 2.5$**
- **At least 1 first Pixel layer hit & 2, 4, or 6 SCT hits** for *$p_T > 100, 200, 300 \text{ MeV}$* respectively.
- Cuts on the **transverse impact parameter**: *$|d_0^{BL}| < 1.5 \text{ mm}$*
- Cuts on the **longitudinal impact parameter**:
 $|\Delta z_0 \sin\Theta| < 1.5 \text{ mm}$ (Δz_0 is difference between tracks z_0 and vertex z position)
- Track fit **χ^2 probability > 0.01** for tracks with *$p_T > 10 \text{ GeV}$*

DATA AND MC SAMPLES

- ❑ Statistics *for 0.9 TeV 357,523 events with 4,532,663 tracks* and for *7 TeV 10,066,072 events with 209,809,430 tracks* passed selection criteria.
- ❑ Integrated luminosities: *$9 \mu\text{b}^{-1}$ at 0.9 TeV and $190 \mu\text{b}^{-1}$ at 7 TeV.*
- ❑ Track and event selection criteria as in the Min Bias 2.0 analysis.
- ❑ The tracking and event efficiencies, unfolding – follow this study.
- ❑ In addition the High Multiplicity (HM) dataset at 7 TeV is studied, for the first time in BEC analyses. Statistic *for 7 TeV (HM) 17784 events with 2,719,536 tracks* were selected with
 - Integrated luminosities: *12.4nb^{-1} at 7 TeV HM.*
- ❑ Closure tests show good agreement between the reconstructed unfolded and truth MC spectra.
- ❑ Study based on the Min Bias 2.0 datasets and MC samples: Pythia MC09 (main), Perugia0, Phojet, and EPOS.

MONTE-CARLO MODELS

Four recent versions of MC event generators were used to provide calculation of R_2 correlation functions and for systematic studies.

➤ **The MC models do not contain BEC effects.**

Generator	Version	Tune	PDF	Focus of Tune	Statistic
PYTHIA 6	6.421	MC09	MRST LO	MB/UE	1.1×10 ⁷ at 0.9 TeV 2.7×10 ⁷ at 7 TeV MBT 1.8×10 ⁶ at 7 TeV HMT
PYTHIA 6	6.421	Perugia0	CTEQ 5L	MB	
PHOJET	1.12.1.35	No tune	MRST LO	MB/UE	
EPOS	1.99 v2965	LHC	CTEQ6.6 LO	MB	for HMT only

➤ Large MC samples of minimum-bias and high-multiplicity events were generated with **PYTHIA 6.421 ATLAS MC09** set of optimised parameters with non-diffractive, single-diffractive and double-diffractive processes included in proportion to the cross sections predicted by the model.

For the study of systematic effects, additional MC samples were produced using the **PHOJET 1.12.1.35**, **PYTHIA with the Perugia0** tune, and the **EPOS 1.99 v2965** for the HM analysis.

- ❑ The **PHOJET** program uses the Dual Parton Model for low- p_T physics and is interfaced to PYTHIA for the fragmentation of partons.
- ❑ The **EPOS** generator is based on an implementation of the QCD inspired Gribov-Regge field theory describing soft and hard scattering simultaneously, and relies on the same parton distribution functions as used in PYTHIA.

TRACK RECONSTRUCTION CORRECTIONS

Performed corrections on:

- ❑ *The reconstruction track efficiency – $\varepsilon(p_T, \eta)$,*
- ❑ *The fraction of secondaries particles – $f_{sec}(p_T, \eta)$,*
- ❑ *The fraction of selected tracks for which the corresponding primary particles are outside the kinematic range – $f_{okr}(p_T, \eta)$,*
- ❑ *The fake tracks – $f_{fake}(p_T, \eta)$,*

We use the formula:

$$w_i(p_T, \eta) = \frac{(1 - f_{sec}(p_T, \eta)) \cdot (1 - f_{okr}(p_T, \eta)) \cdot (1 - f_{fake}(p_T, \eta))}{\varepsilon(p_T, \eta)}$$

- *The effect of events lost due to the trigger and vertex reconstruction corrected using even-by-event weights applied to a pair of particles*
- *The resolution of Q obtained to be better than 5 MeV so to exclude track fake reconstruction the Q-threshold taken 20 MeV*

COULOMB CORRECTION

The measured $N(Q)$ *distribution* for the like or unlike signed particle (track) pairs in presence of the Coulomb interaction is given by:

$$N_{meas}(Q) = G(Q)N(Q)$$

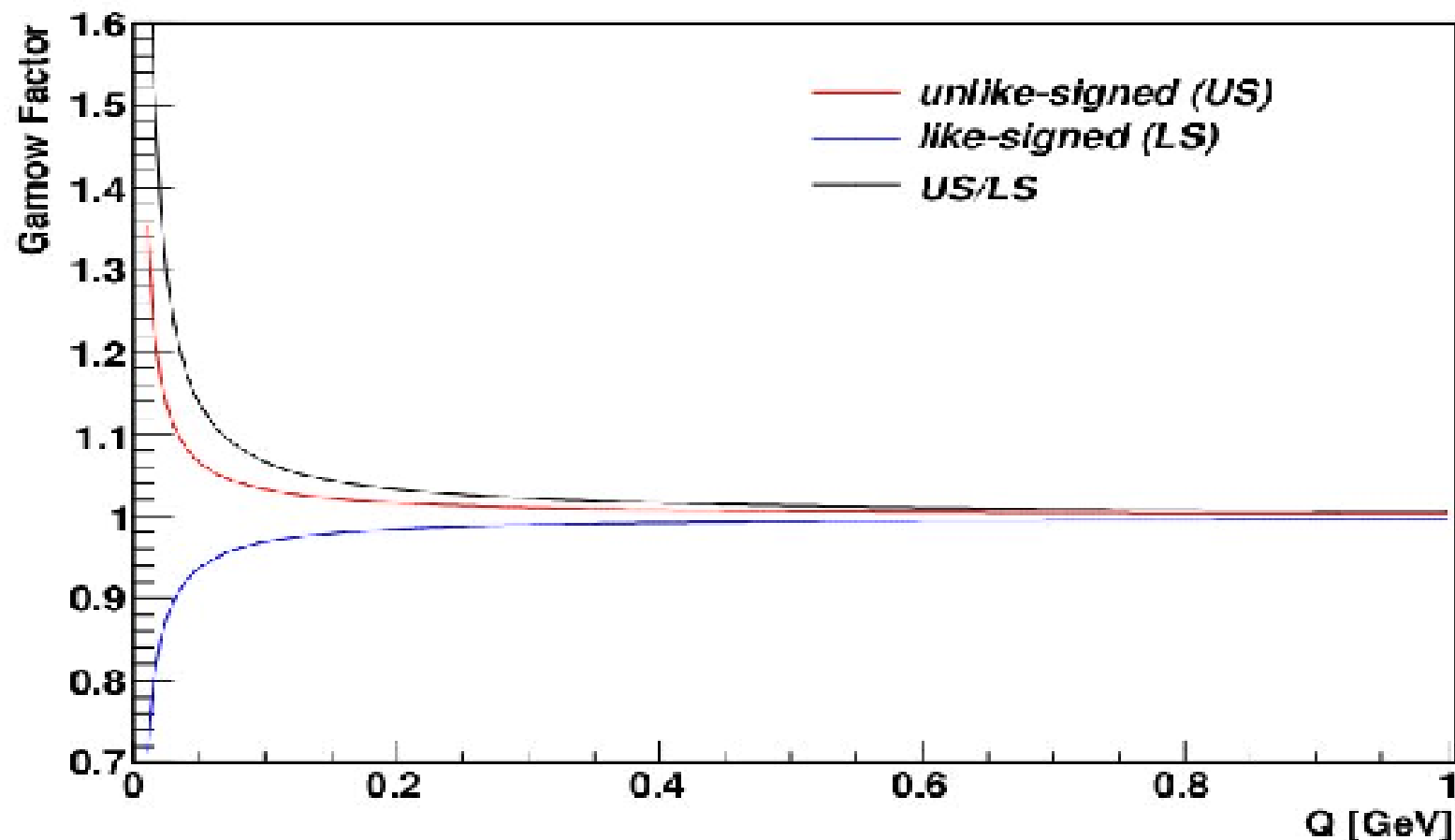
where $N_{meas}(Q)$ is the measured distribution, $N(Q)$ is the distribution free of Coulomb correlations.

Gamov penetration factor

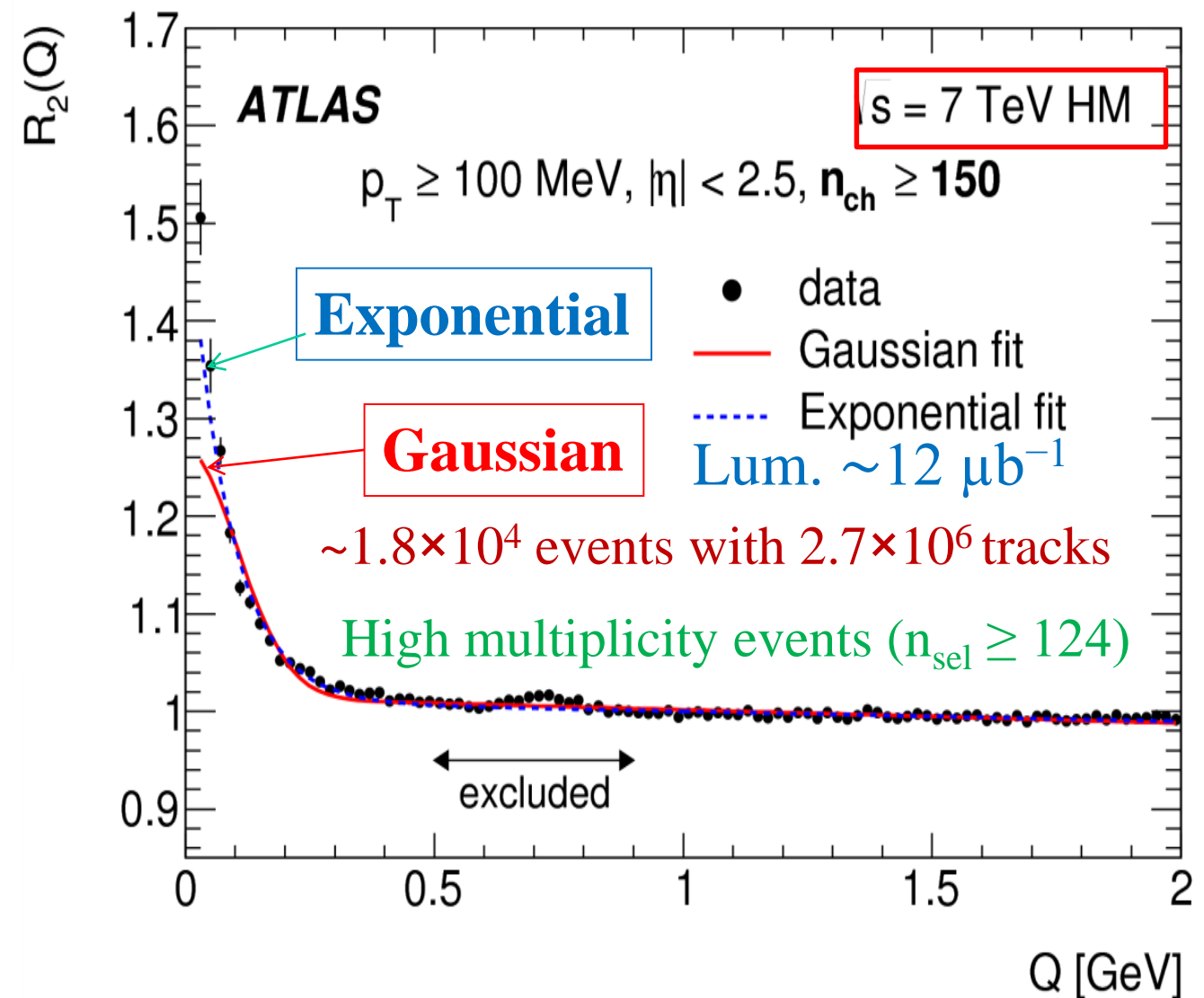
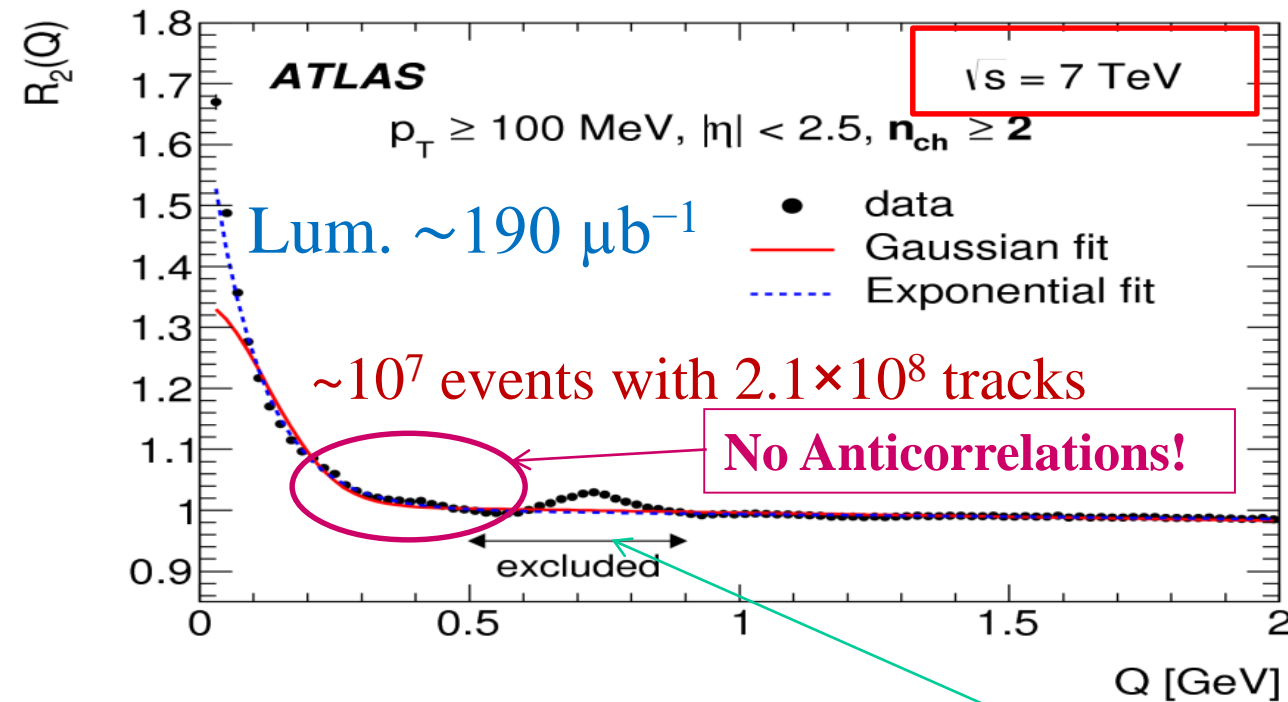
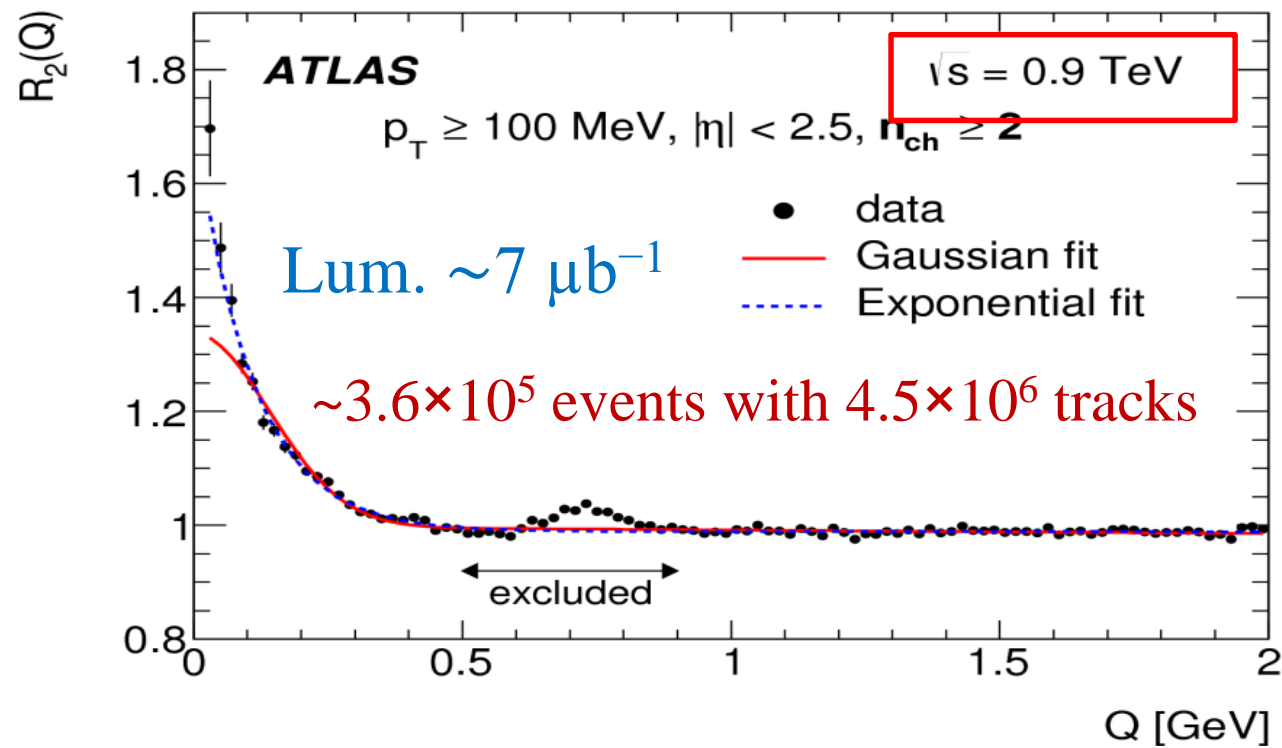
$$G(Q) = \frac{2\pi\eta}{e^{2\pi\eta} - 1}$$

Sommerfeld parameter

$$\eta = \frac{\pm\alpha_m}{|Q|}$$



Distribution of Gamma factor for unlike-signed particle pairs, like-signed particle pairs and ratio of unlike-signed to like-signed.



Studies of one-dimensional BEC effects in pp collisions for $p_T > 100 \text{ MeV}$ and $|\eta| < 2.5$ at centre-of-mass energies of 0.9 and 7 TeV.

Fit to extract strength and source size. **Goldhaber** spherical shape with a **Gaussian** distribution of the source. **Exponential**, radial Lorentzian distribution of the source -> much better at low Q. Bump in ρ -meson region because **MC overestimates $\rho \rightarrow \pi\pi$** , therefore **region 0.5 – 0.9 GeV excluded** from the fit. **Q region is from 0.02 to 2 GeV.**

- The systematic uncertainties of the inclusive Bose-Einstein correlation parameters, \mathbf{R} (the effective radius parameter of the source) and λ (the strength of the effect parameter), of the *fit of $R_2(Q)$ correlation functions with exponential model* are summarized in the Table.
- The systematic uncertainties are combined by adding them in quadrature and the resulting values are given in the bottom row.
- The same sources of uncertainty are considered for the differential measurements in n_{ch} and the average transverse momentum k_T of a pair, and their impact on the fit parameters is found to be similar in size.

Source	0.9 TeV		7 TeV		7 TeV (HM)	
	λ	R	λ	R	λ	R
Track reconstruction efficiency	0.6%	0.7%	0.3%	0.2%	1.3%	0.3%
Track splitting and merging	negligible		negligible		negligible	
Monte Carlo samples	14.5%	12.9%	7.6%	10.4%	5.1%	8.4%
Coulomb correction	2.6%	0.1%	5.5%	0.1%	3.7%	0.5%
Fitted range of Q	1.0%	1.6%	1.6%	2.2%	5.5%	6.0%
Starting value of Q	0.4%	0.3%	0.9%	0.6%	0.5%	0.3%
Bin size	0.2%	0.2%	0.9%	0.5%	4.1%	3.4%
Exclusion interval	0.2%	0.2%	1%	0.6%	0.7%	1.1%
Total	14.8%	13.0%	9.6%	10.7%	9.4%	10.9%

COMPARISON WITH OTHER EXPERIMENTS

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The results of BEC parameters for Exponential fits of R_2 used total uncertainties

Statistical uncertainties are below 2–4 %

Energy [GeV]	n_{ch}	λ	R [fm]
0.9	≥ 2	0.74 ± 0.10	1.83 ± 0.25
7	≥ 2	0.71 ± 0.07	2.06 ± 0.22
7 (HM)	≥ 150	0.74 ± 0.06	2.36 ± 0.30

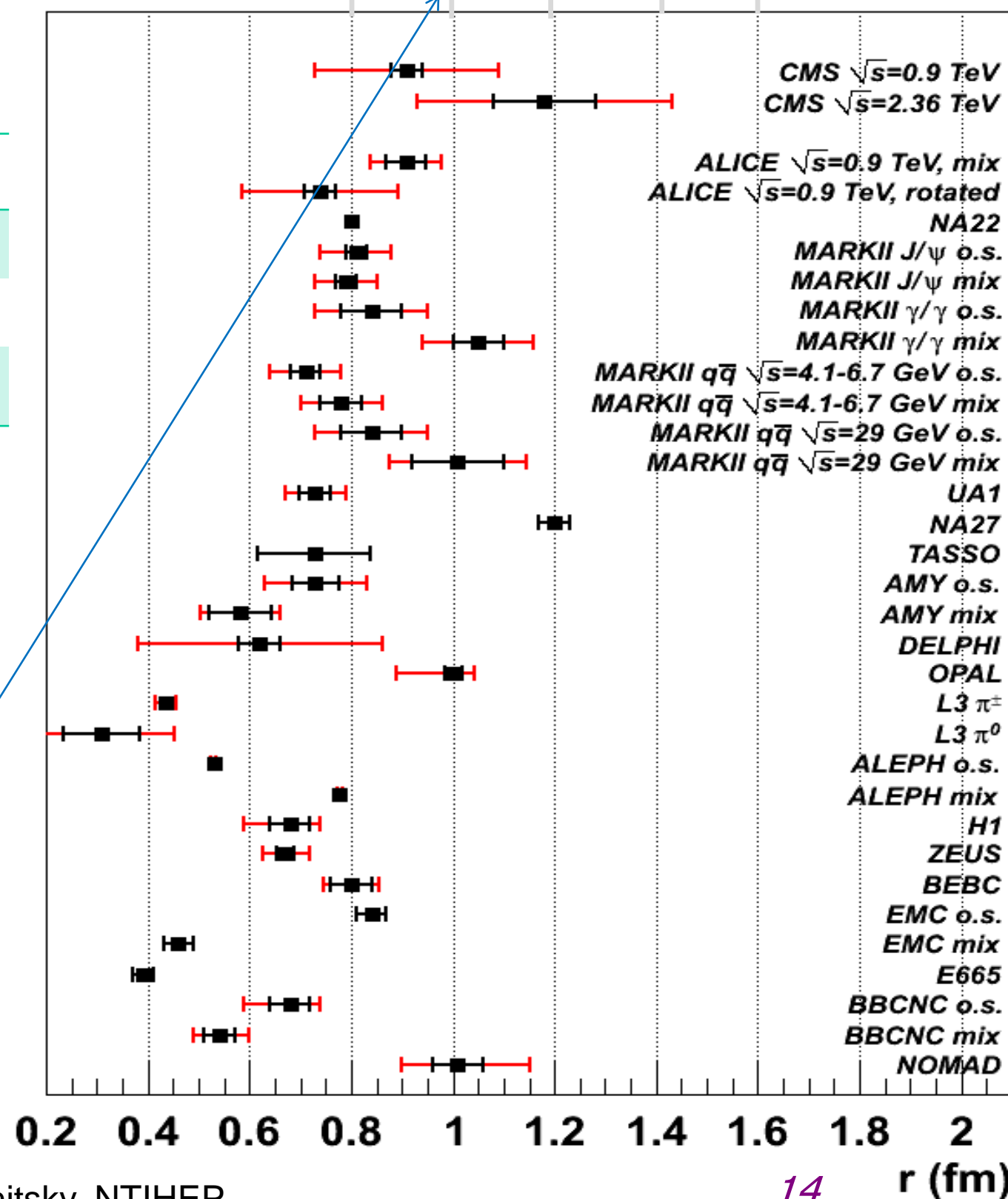
Comparison with results of previous experiments

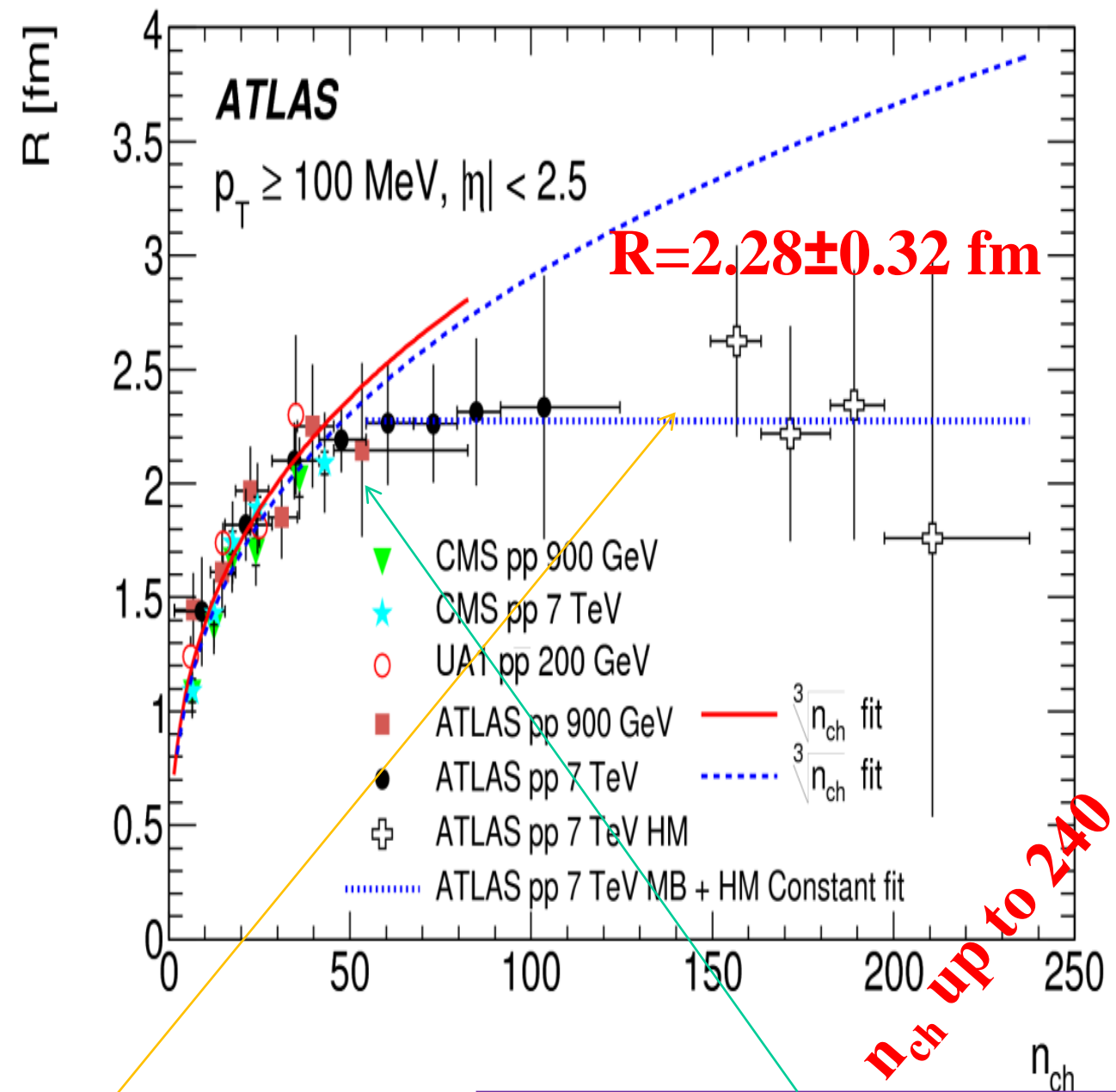
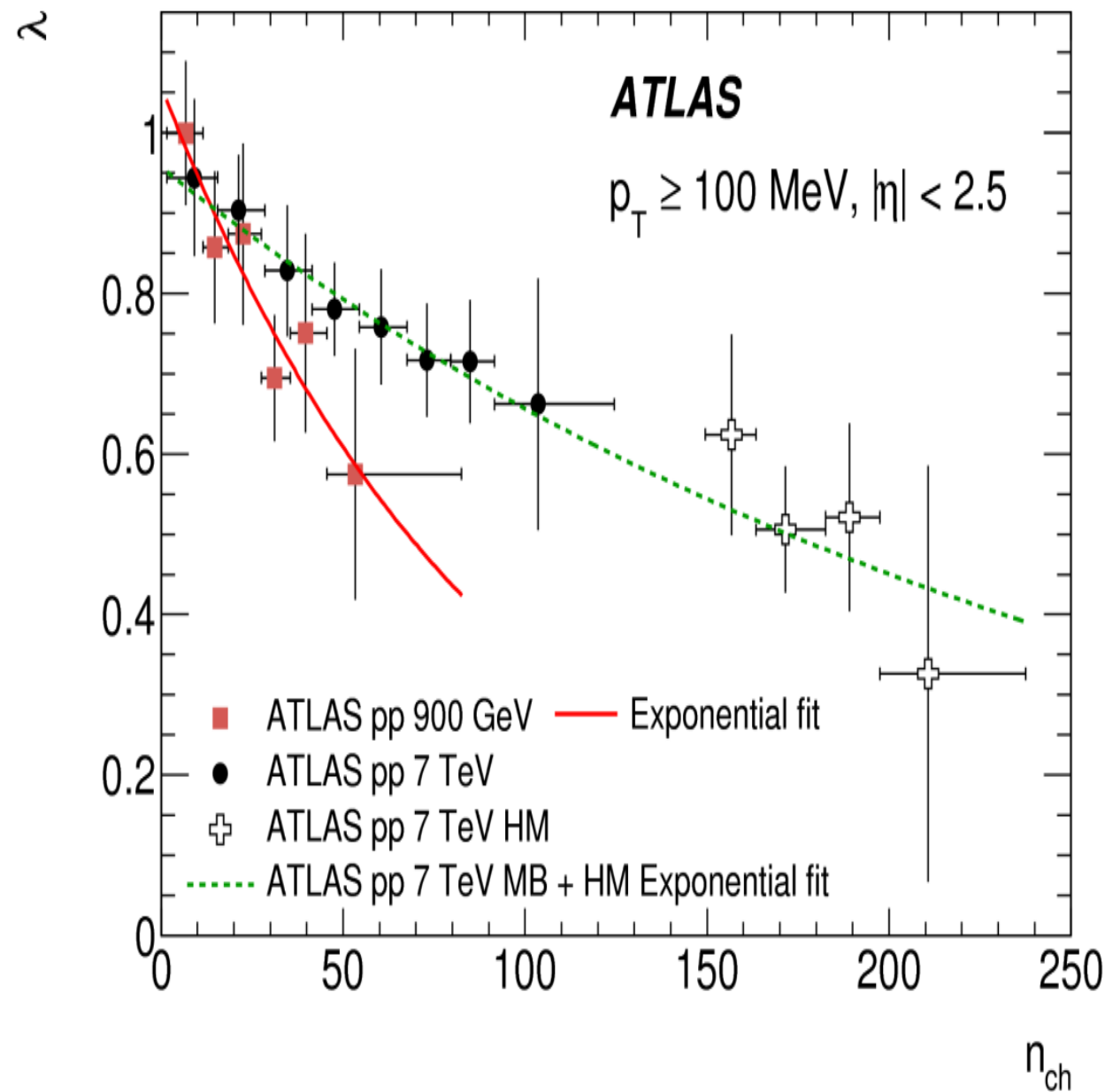
Most of the previous experiments provided R measurement using the Gaussian fit. The calculation of Gaussian result from the Exponential fit can be done using the scale factor $\sqrt{\pi}$:

$$R^{(G)} = R^{(E)} \sqrt{\pi}$$

Energy [GeV]	n_{ch}	R [fm]
0.9	≥ 2	1.03 ± 0.14
7	≥ 2	1.16 ± 0.12
7 (HM)	≥ 150	1.33 ± 0.17

ATLAS $\sqrt{s} = 7$ TeV HMT
 ATLAS $\sqrt{s} = 7$ TeV
 ATLAS $\sqrt{s} = 0.9$ TeV





- ▶ λ and R are energy independent within the uncertainties
- ▶ λ exponentially decrease with multiplicity

Good Agreement with
CMS & UA1

- ▶ R of the $\alpha \cdot n_{ch}^{1/3}$ fit for $n_{ch} \leq 55$: 0.9 TeV is $\alpha = 0.64 \pm 0.07$ fm, 7 TeV is $\alpha = 0.63 \pm 0.05$ fm
- ▶ R is a **Constant** for $n_{ch} > 55$ at 7 TeV $R = 2.28 \pm 0.32$ fm observed for the first time

THEORY PREDICTION

V.A. Shegelsky, et al, Phys Letter B703 (2011) 288; Nucl Phys B219 (2011) 10

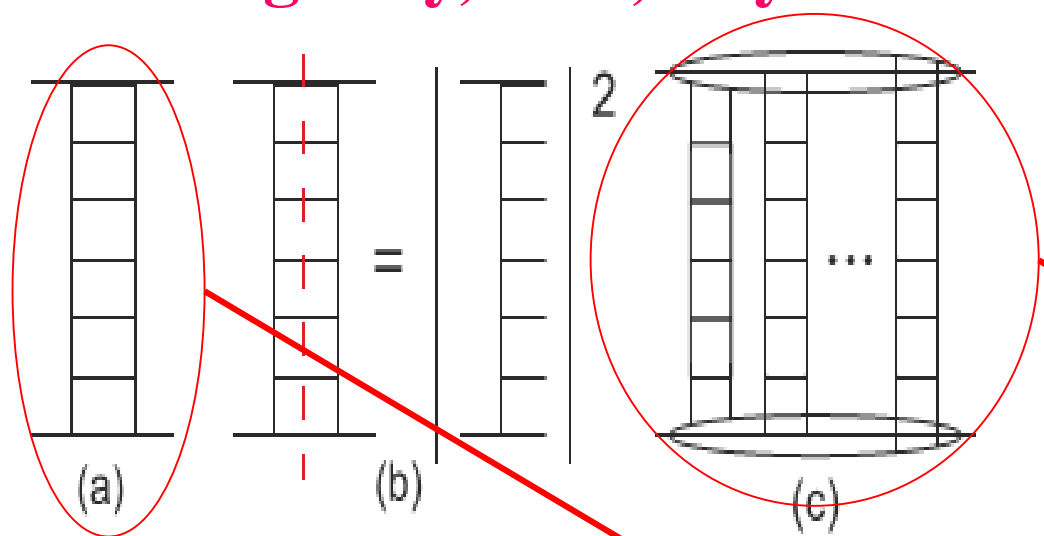
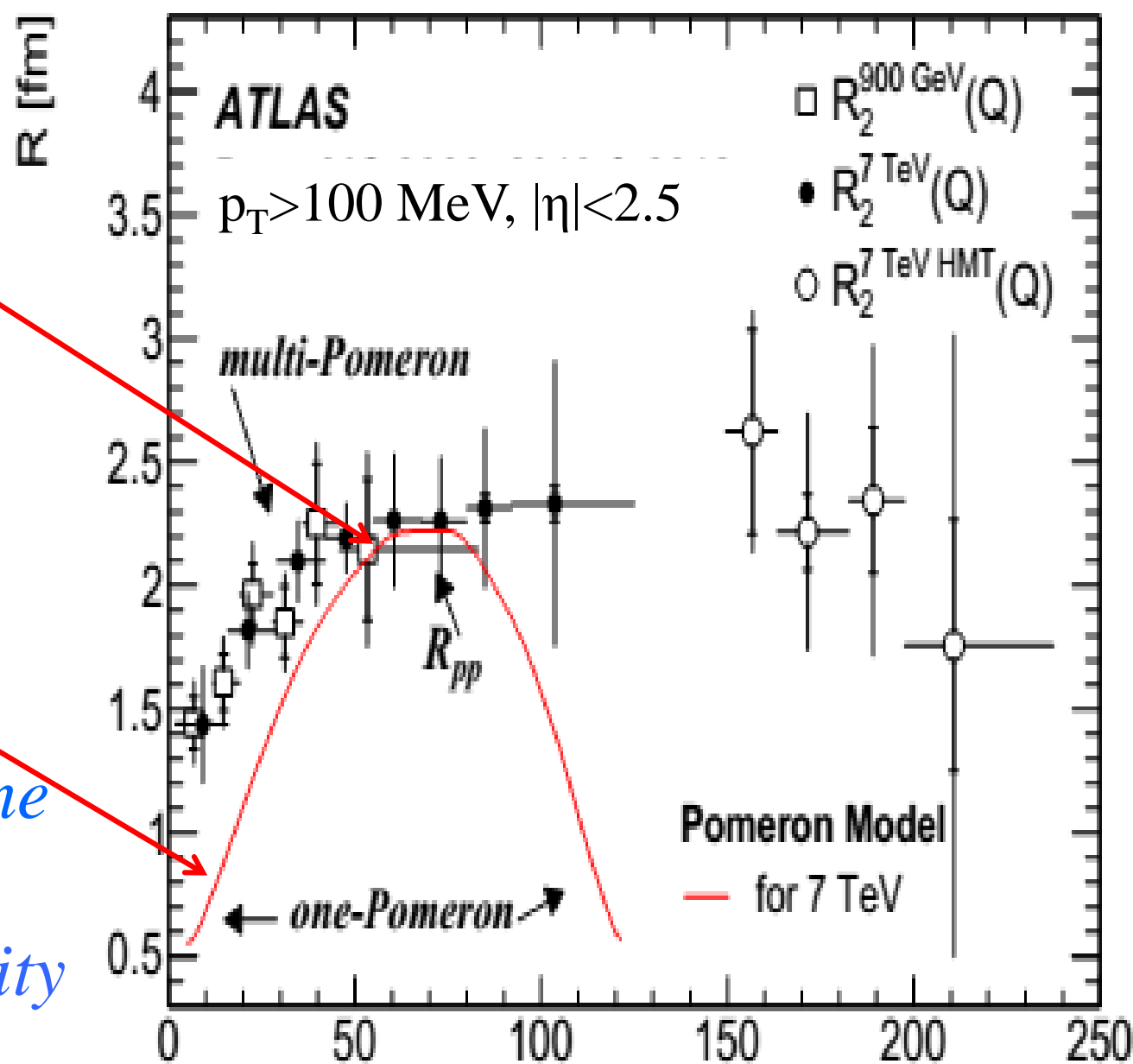


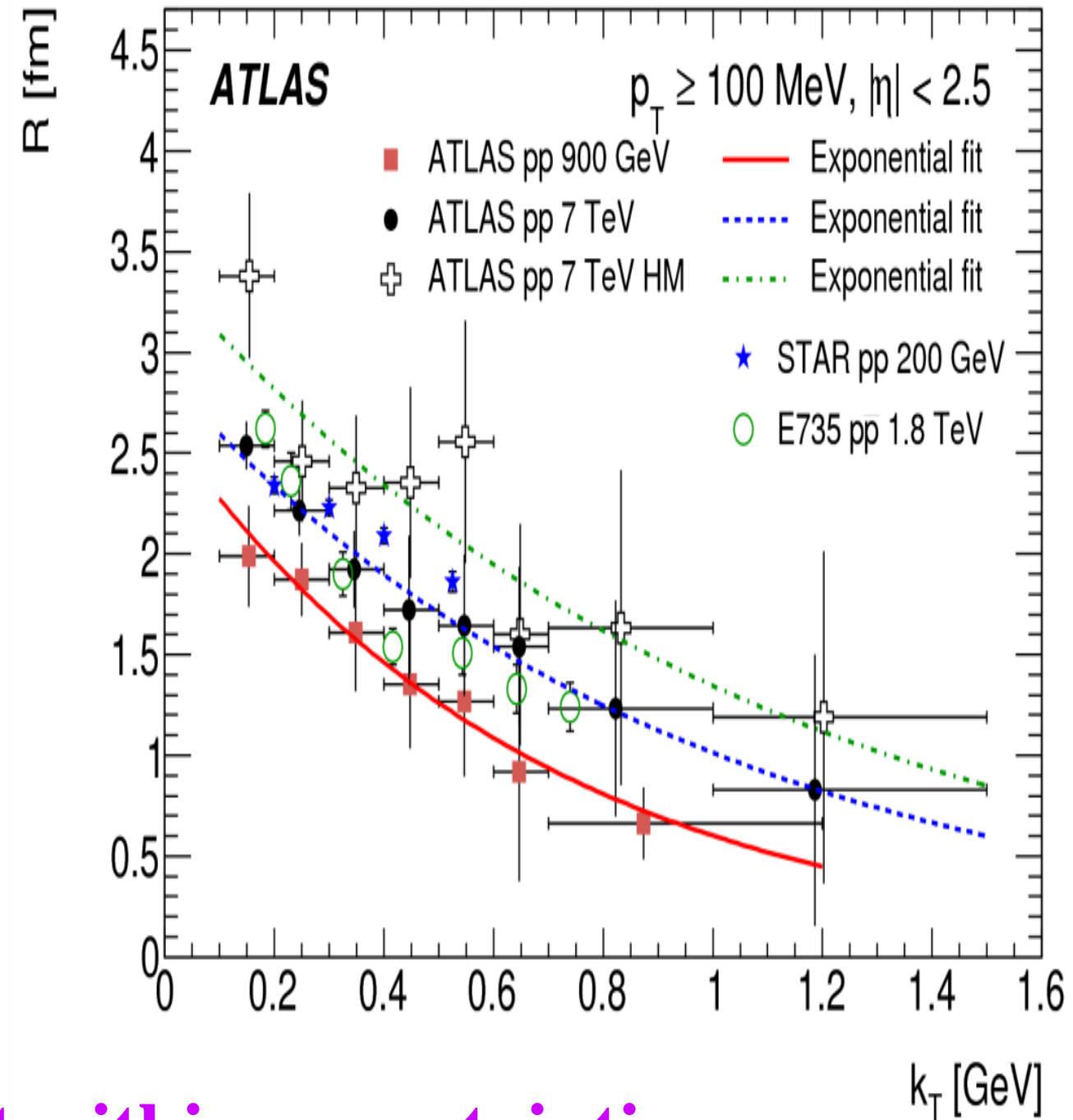
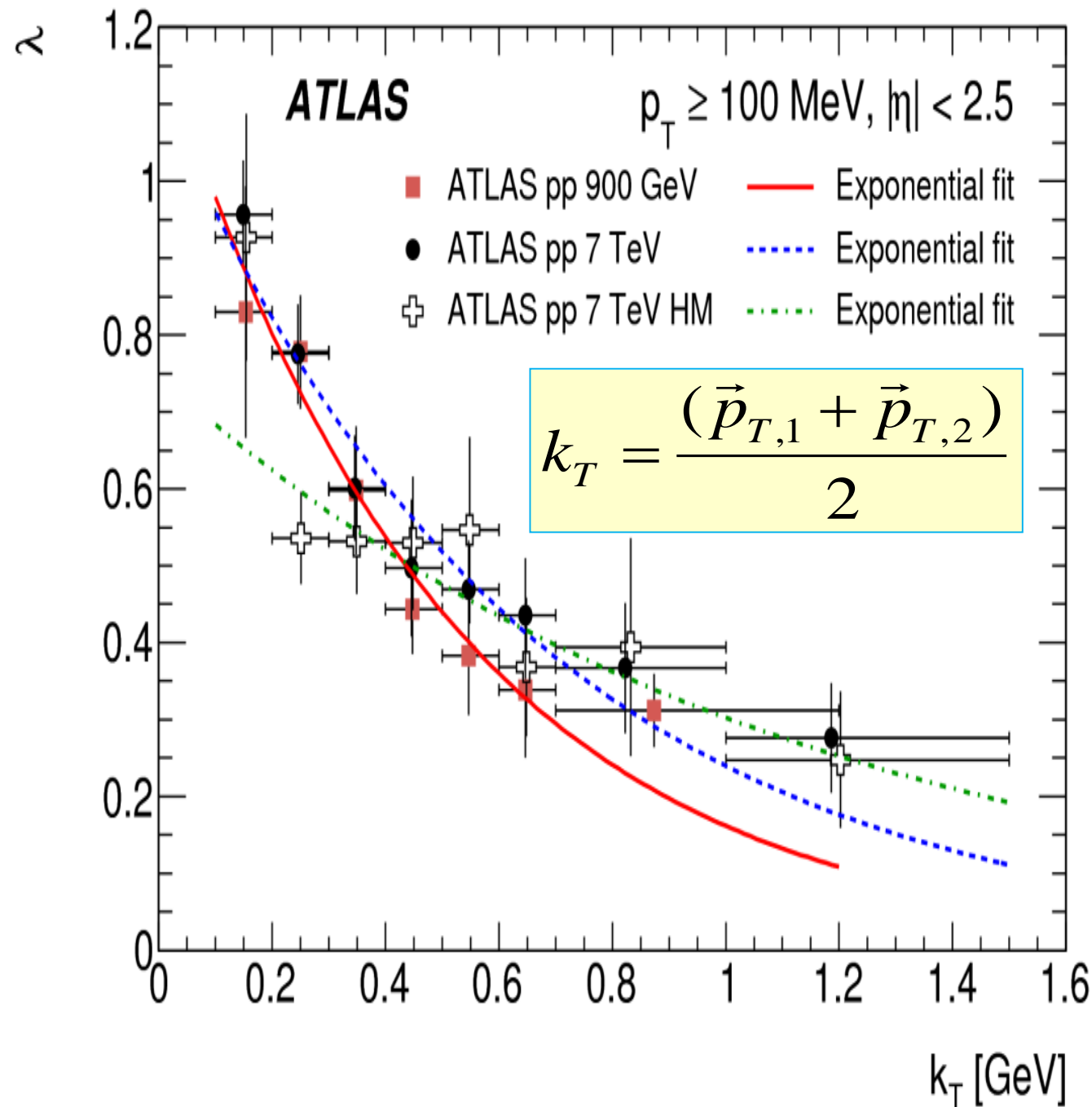
Fig. 1. (a) The ladder diagram for one-Pomeron exchange; (b) cutting one-Pomeron exchange leads to the multiperipheral chain of final state particles; (c) a multi-Pomeron exchange diagram.



The prediction of **Pomeron model** for n_{ch} $R_{\max}=2.2$ fm is in agreement with our saturated radius $R=2.3$ fm. There is not principal agreement with data for $n_{ch} \geq 80$

k_T DEPENDENCE OF λ AND R PARAMETERS

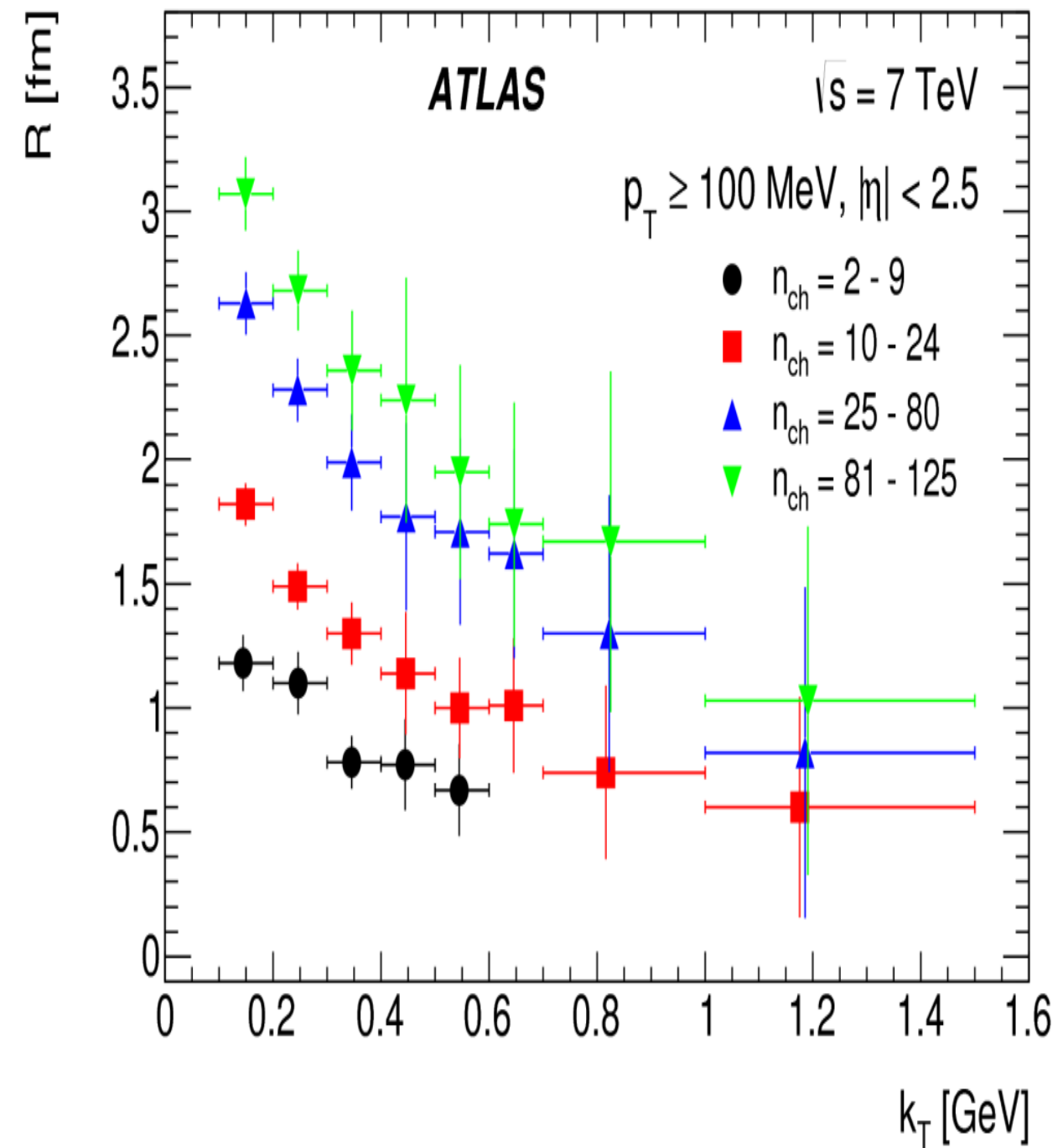
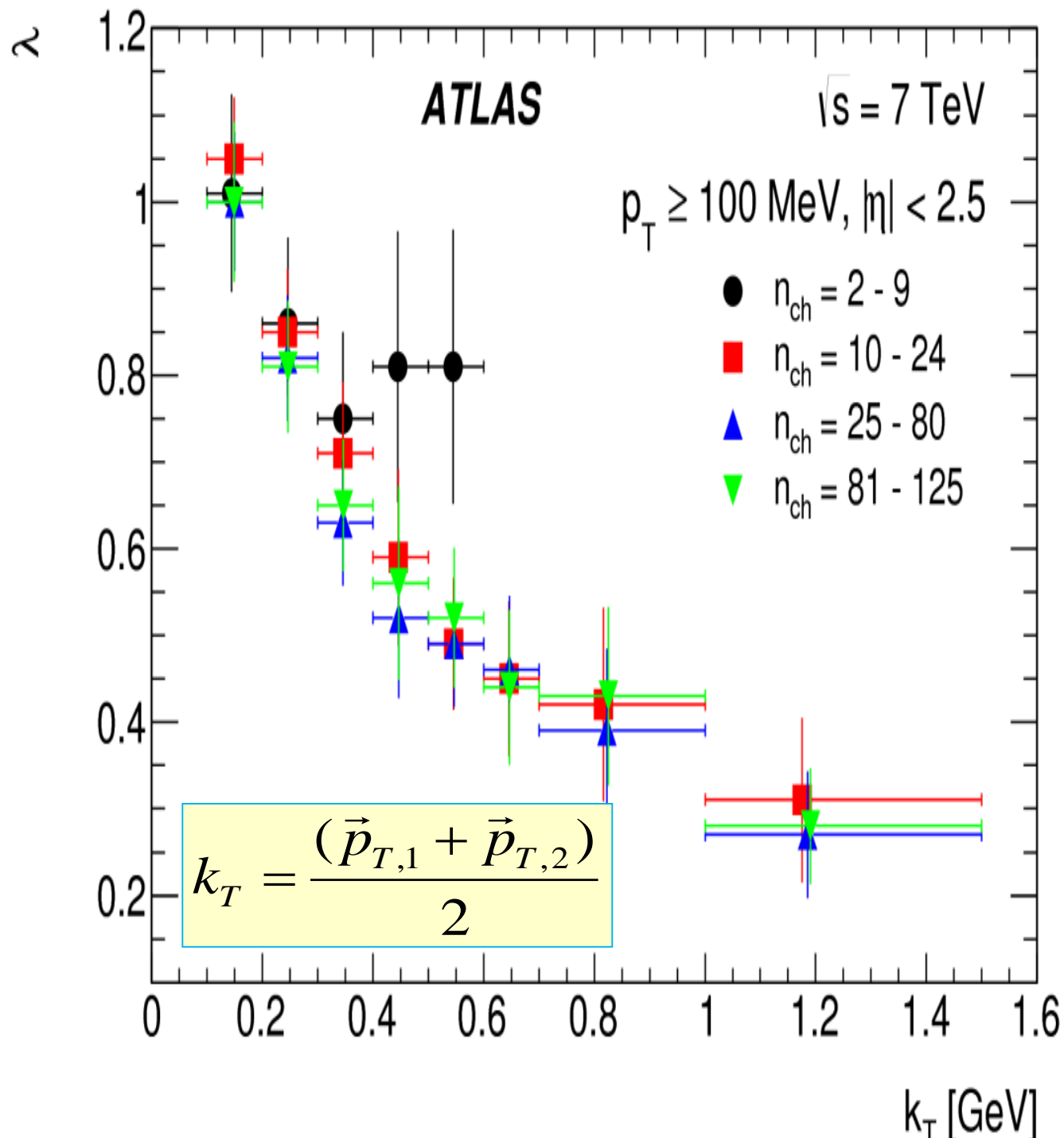
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- ▶ λ and R are energy-independent within uncertainties
- ▶ λ and R decrease exponentially with k_T
- ▶ Good agreement with earlier (non-LHC) measurements

k_T DEPENDENCE OF λ AND R PARAMETERS

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- ▶ **No k_T –dependence of λ for different multiplicity intervals**
- ▶ **k_T –dependence of R shows R increasing with multiplicity interv.**

CONCLUSION

- The Bose-Einstein correlations results of identical charged particles pairs measured in $|\eta| < 2.5$, $p_T > 100 \text{ MeV}$ in pp collisions at 0.9 & 7 TeV with the ATLAS experiment are presented.
- *For the first time* the multiplicity dependence of the BEC is investigated up to *very high multiplicities* (≈ 240).
- *For the first time* a *saturation effect* in the multiplicity dependence of the extracted BEC radius parameter is observed: $R = 2.28 \pm 0.02 \pm 0.31 \text{ fm}$ for high multiplicity region, $n_{\text{ch}} > 55$.
- The n_{ch} & k_T dependences of λ parameter are well described by the exponential function.
- The dependence of the BEC parameters on k_T is investigated for different multiplicity regions up to high multiplicity and is observed to be well described by the exponential function.
- The k_T dependence of R is obtained to increase with increasing of the multiplicity.

BACKUP SLIDES

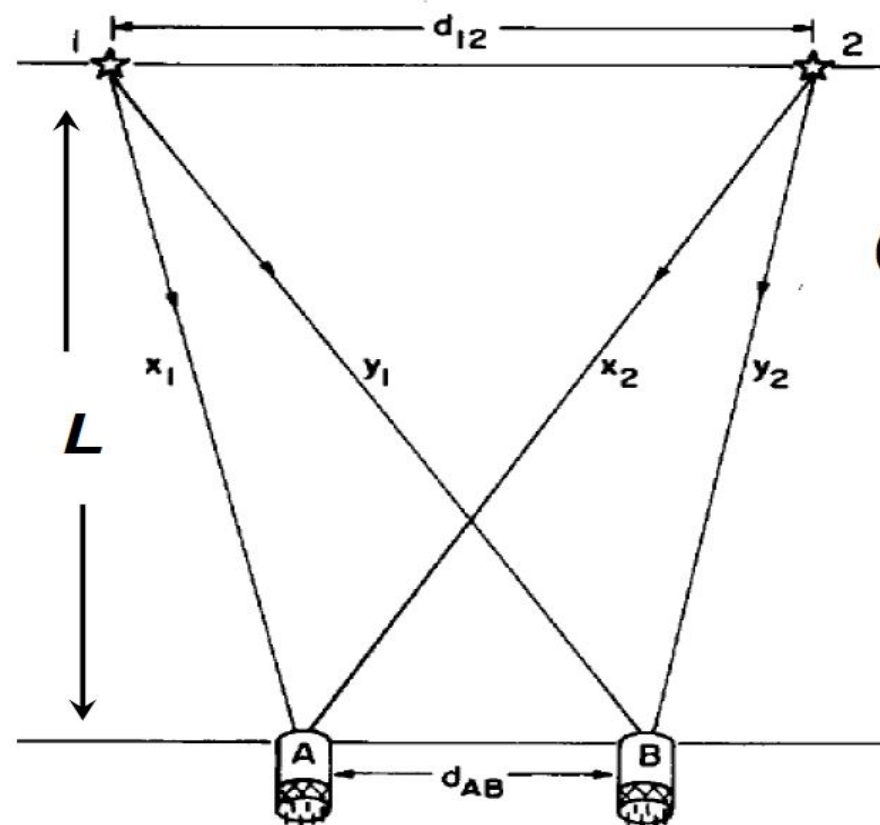
BOSE-EINSTEIN CORRELATIONS AND HANBURY BROWN – TWISS INTERFEROMETRY

BEC are often considered to be the analogue of the Hanbury Brown and Twiss effect in astronomy, describing the interference of incoherently-emitted identical bosons.

Intensity interferometry of photons in radio-astronomy: measures angular diameter of two stars, so the physical size of the source



Figure 10.1 The first stellar intensity interferometer; the pilot model of the stellar intensity interferometer at Jodrell Bank in 1955. Two Army searchlights were used to make the first measurement of the angular diameter of a main sequence star (Sirius).



$$C(d) = \frac{\langle I_1 I_2 \rangle}{\langle I_1 \rangle \langle I_2 \rangle}$$

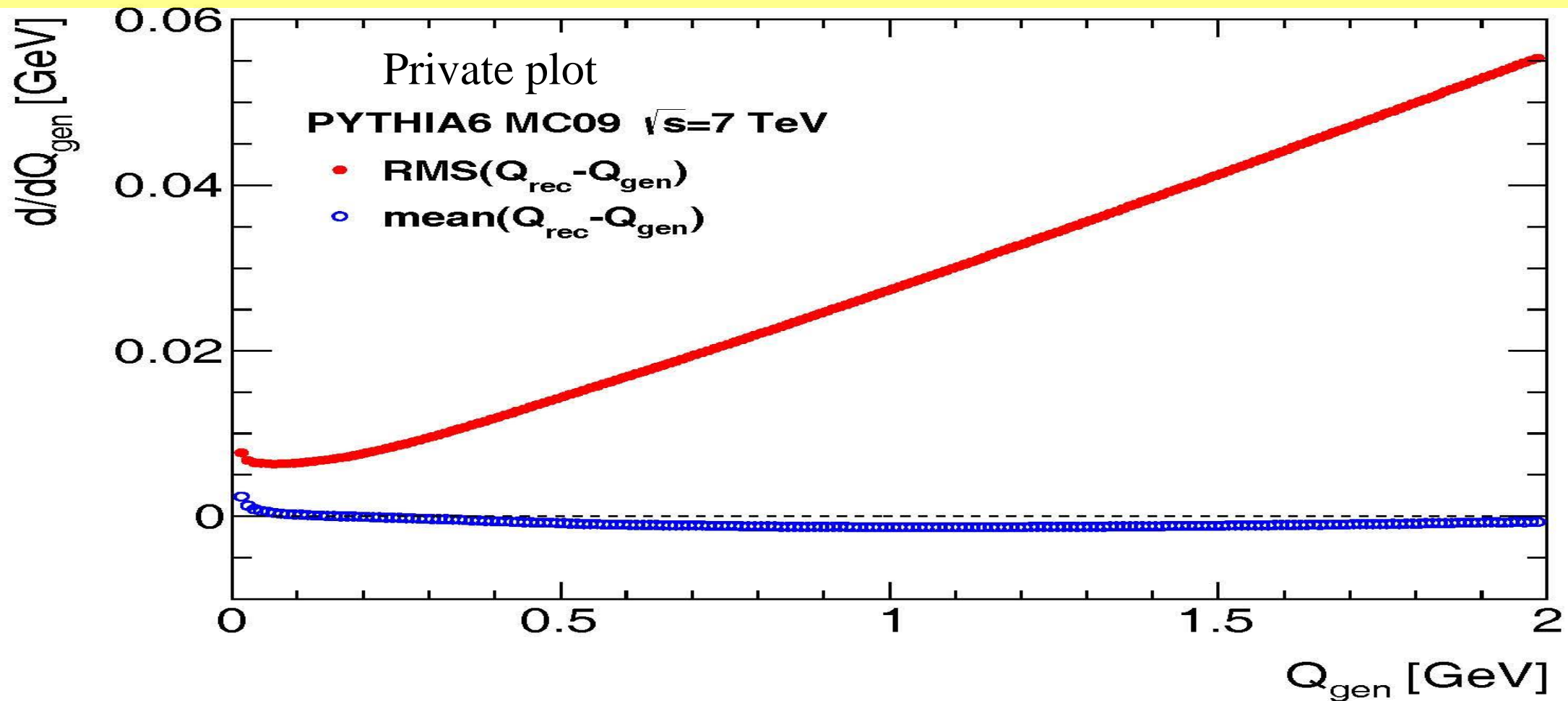
$$= 1 + A \cos(d_{AB})$$

$$d_{AB} = \lambda / \theta$$

$I_{1(2)}$ - intensities, $\langle x \rangle$ - averaging over random phases
 λ is the wavelength of the light, $\theta = d_{12}/L$

Varying d_{AB} one learns the angle, and using the individual wave vectors, the physical size of the source

Q-RESOLUTION



The estimated Q resolution and average bias of the reconstructed momentum difference as a function of the Q true generated value

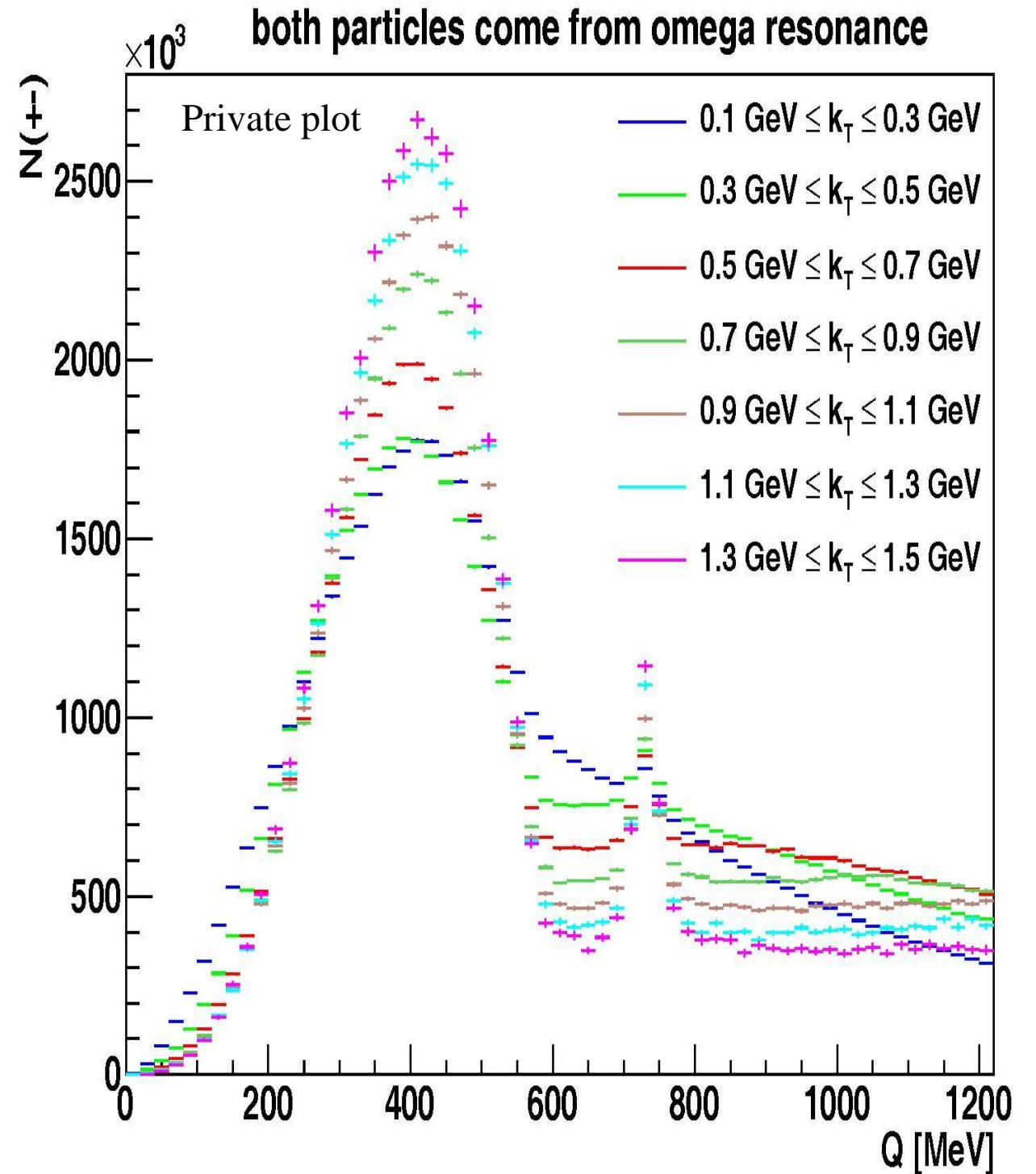
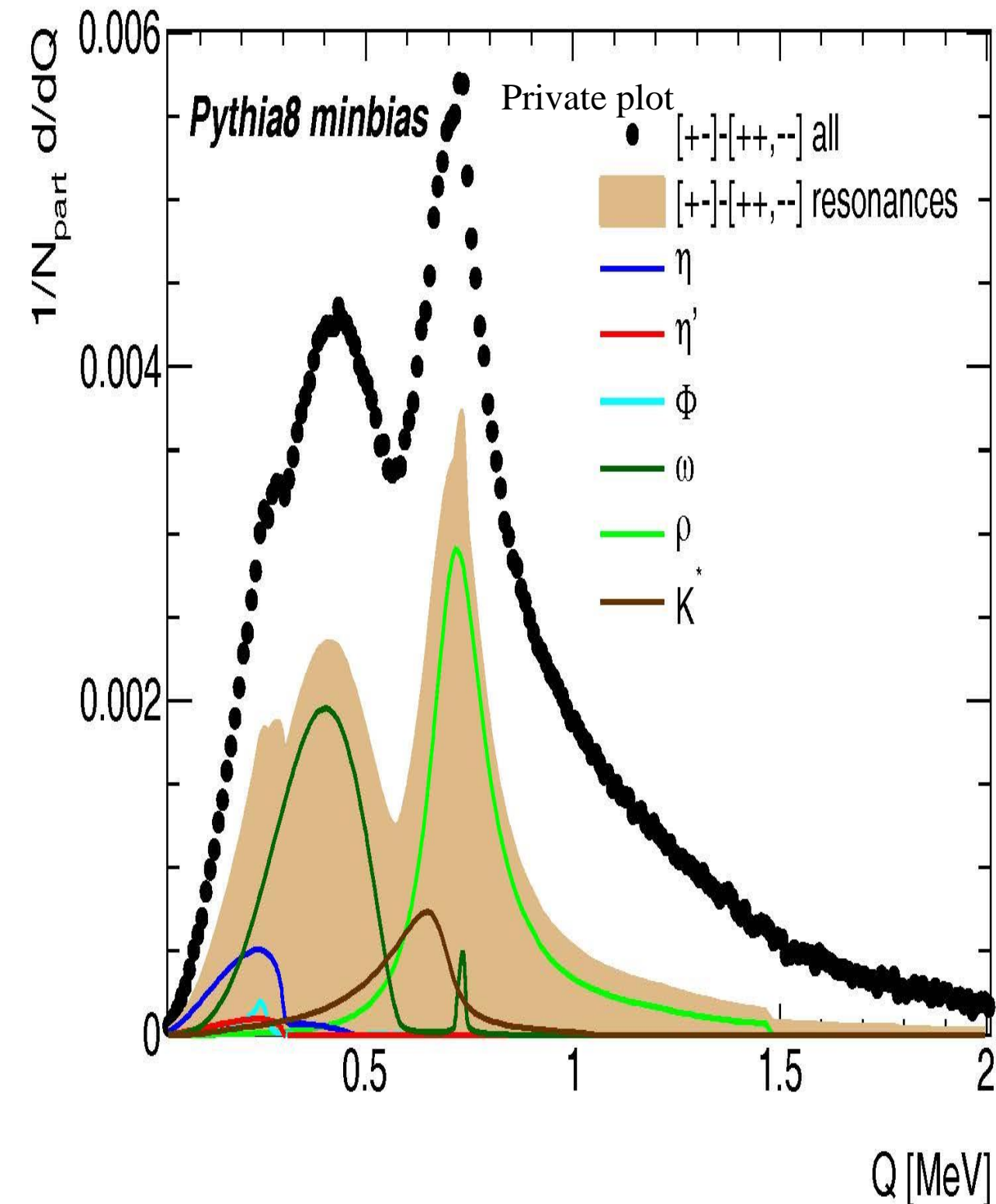
- ❑ The basic variable in which correlation functions are expressed is the scalar momentum difference Q .
- ❑ The ATLAS detector Q resolution is found using MC.
- ❑ To exclude the region of possible two track fake reconstruction a small Q threshold was introduced, $Q > 20$ MeV, as a minimal Q between two tracks.
- ❑ Using in fit of $R_2(Q)$ correlation functions.

$$R_2(Q) = \int_{Q_{\min}}^{Q_{\max}} R_2(Q') \frac{1}{\sqrt{2\pi}\sigma(Q')} e^{-\frac{(Q-Q')^2}{2\sigma(Q')^2}} dQ',$$

PYTHIA6 ATLAS MC09 TUNE

- ❑ Newest version of the color connection
- ❑ Bowler fragmentation function for heavy quarks
- ❑ Latest tune to **LO*PDF** set derived based on previous MC tunes → Increased cut-off scale (PARP(82)) of 2.3 GeV and rescale exponent (PARP(90)) to 0.25
- ❑ Significantly lower **UE** at the LHC energy than the previous version
- ❑ Tunes using **PROFESSOR** gives better description

THE RESONANCES STUDY

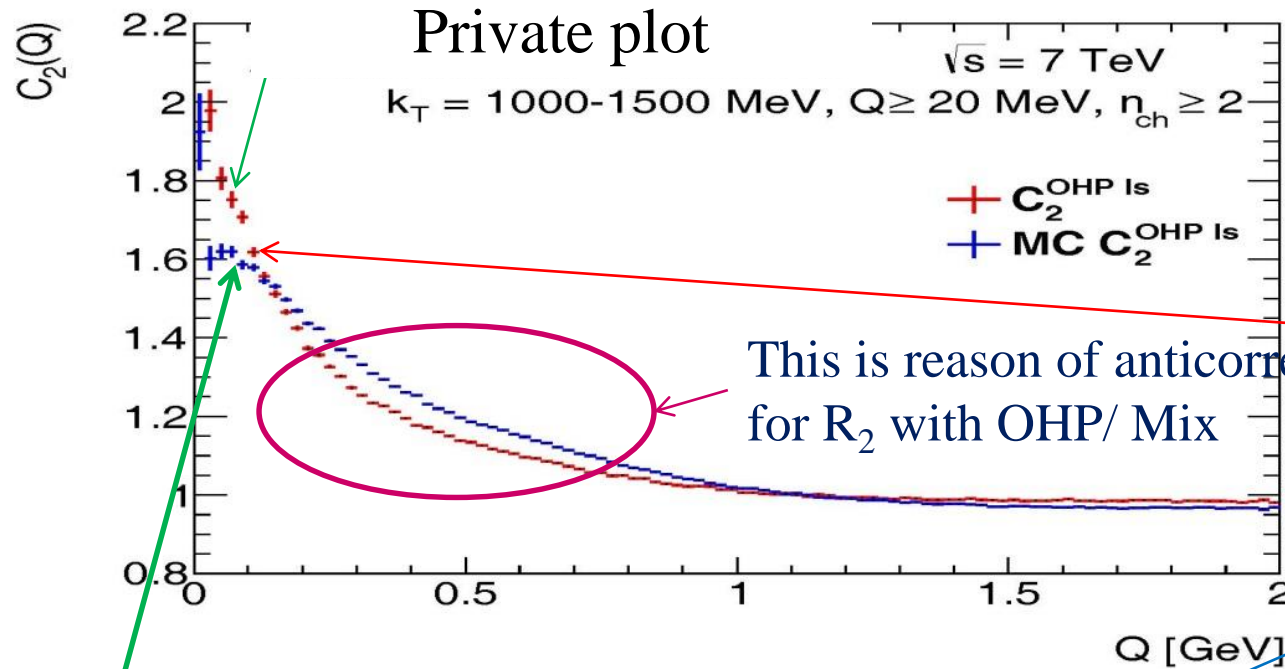


The ΔQ spectrum generated by Pythia8 and the decomposition of its resonant part into leading contributions at 7 TeV. The normalised ω -meson peak is increasing with k_T increase.

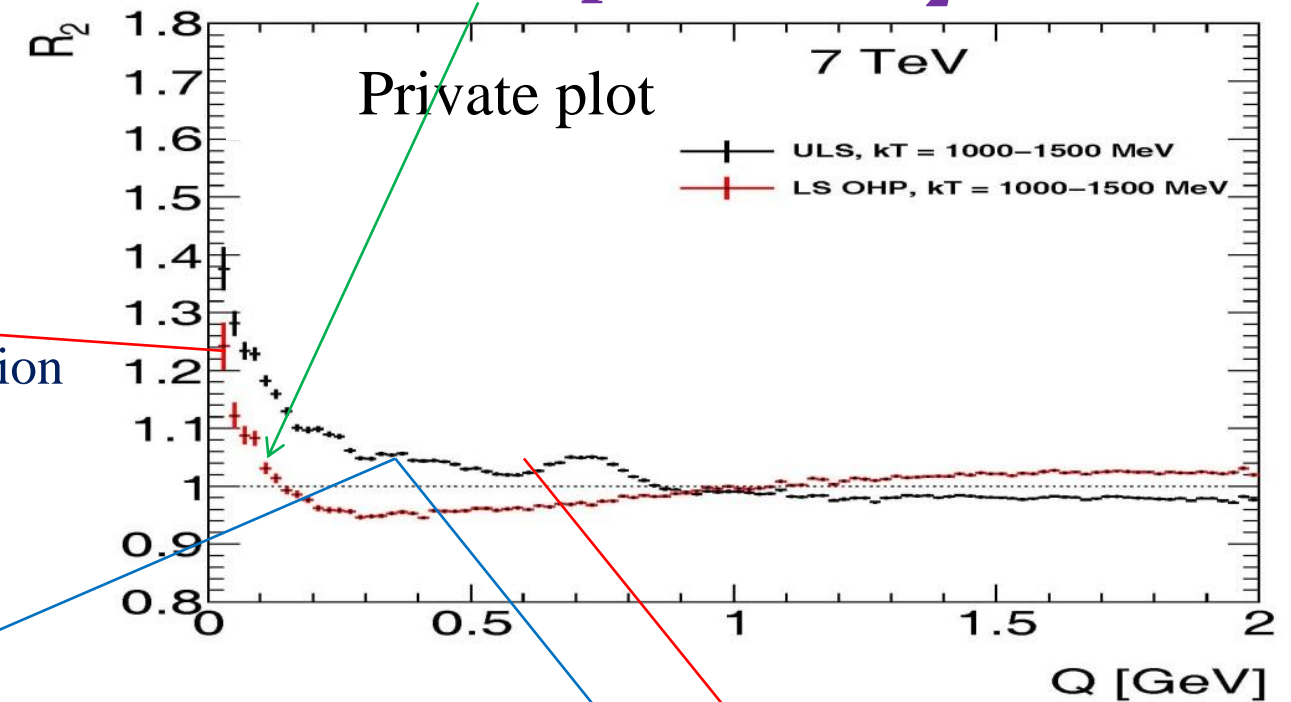
OHP (MIX) AND UCP REFERENCE SAMPLES: K_T

The two-particle correlation function $C_2(Q)$ at 7 TeV for different k_T intervals using the opposite-hemisphere reference sample for data (red) and MC (blue)

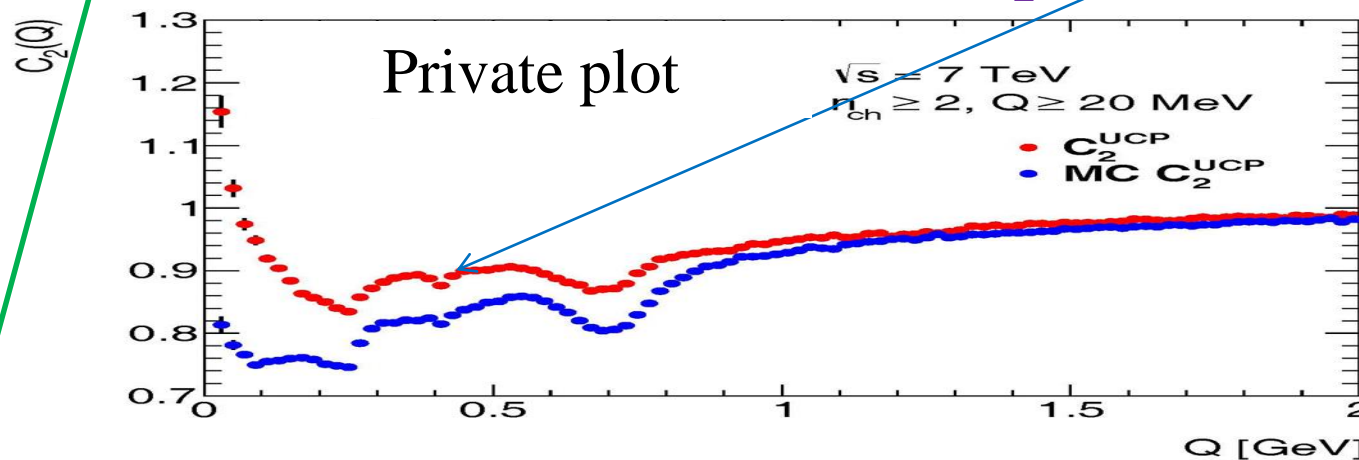
Artificial peak in C_2 in BEC region



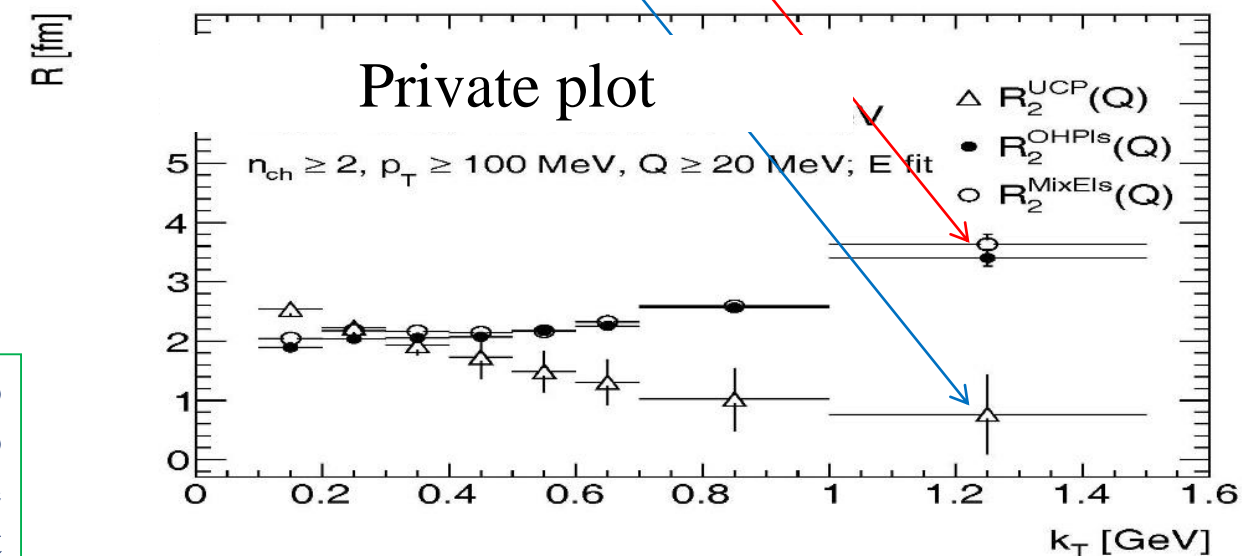
Small BEC peak for R_2 OHP



Reflection of resonances in C_2 UCP



k_T dependences for parameter R for R_2 with different reference samples



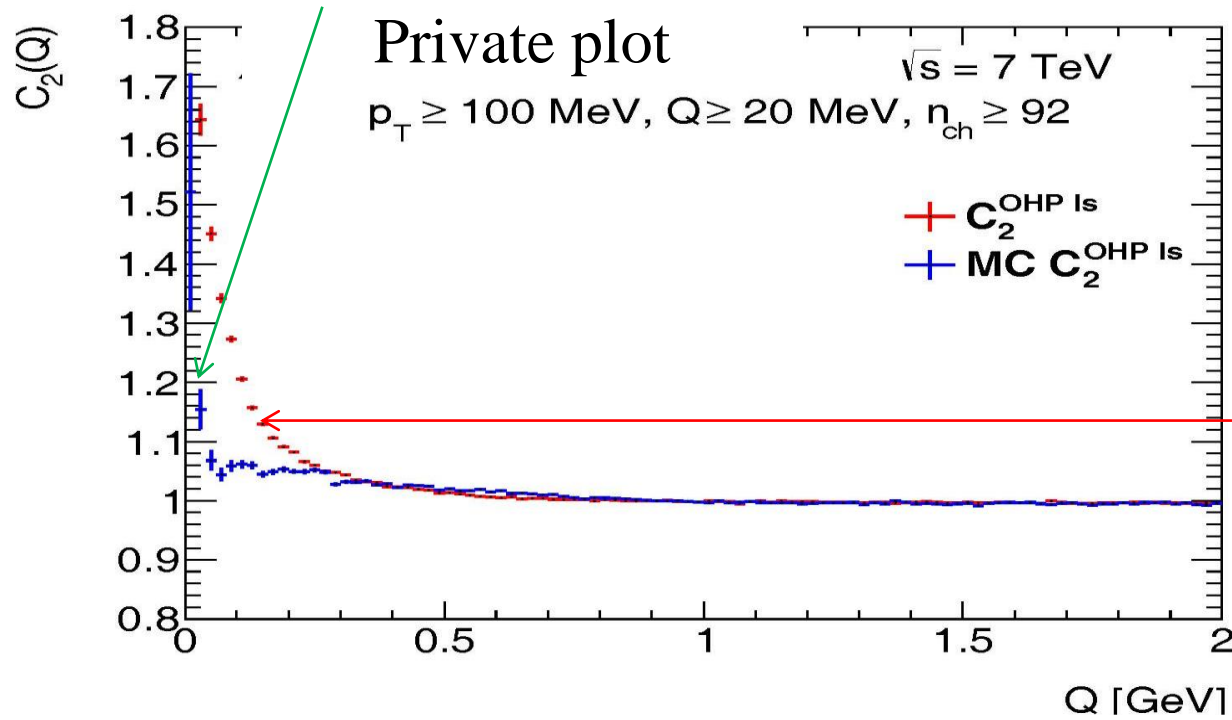
To note is that the slope in the MC can be explained by the fact that MC is tuned to the data and so reflects different dynamical constraints, but MC has no possibility to reproduce a peak at small Q as in the data but shows a broad enhancement.. The additional correlations in large multiplicity intervals seem to be due to multi-jet events in MC where the correlations between particles within the same jet can contribute to the region of low- Q s. In this case, the single-ratio $C_2(Q)$ correlation function numerator contains contributions from multi-jets, while the denominator does not have this effect as no correlations are expected in randomly paired particles.

Disadvantage for OHP/MIX: violation of energy-momentum constraint, event topology, destroying other features such as non-BEC etc.

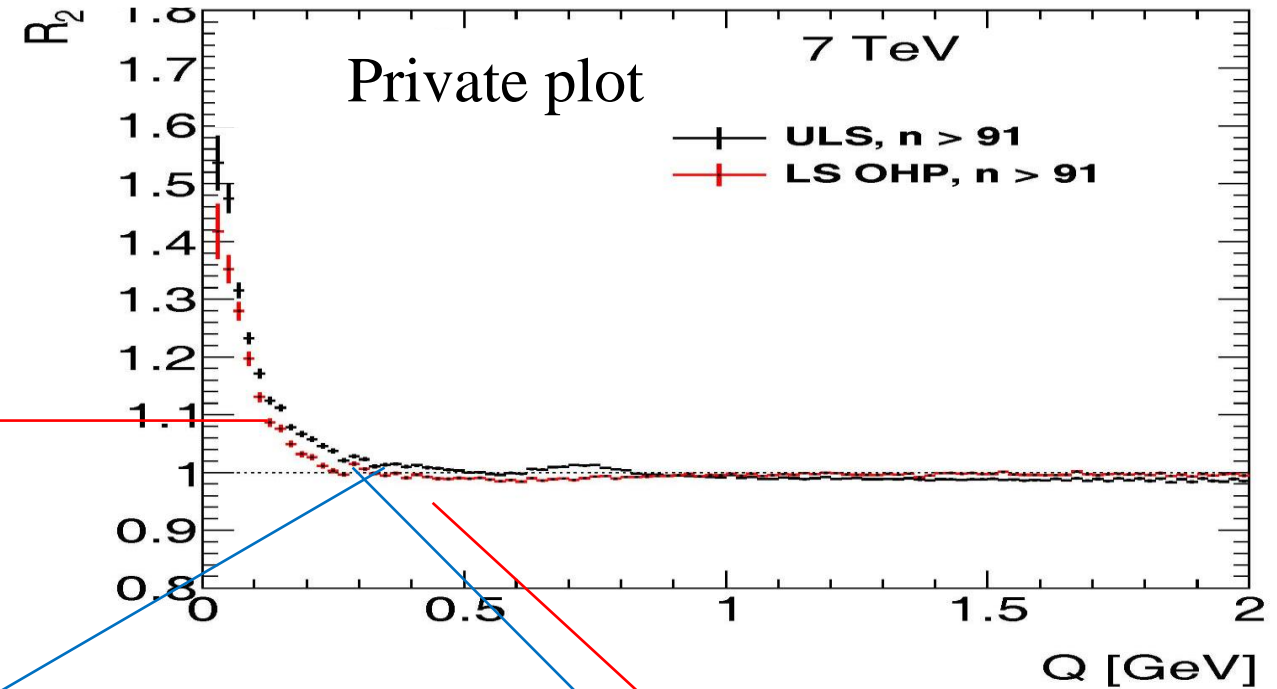
OHP(MIX) AND UCP REFERENCE SAMPLES: N_{CH}

The two-particle correlation function $C_2(Q)$ at 7 TeV for different n_{ch} intervals using the opposite-hemisphere reference sample for data (red) and MC (blue)

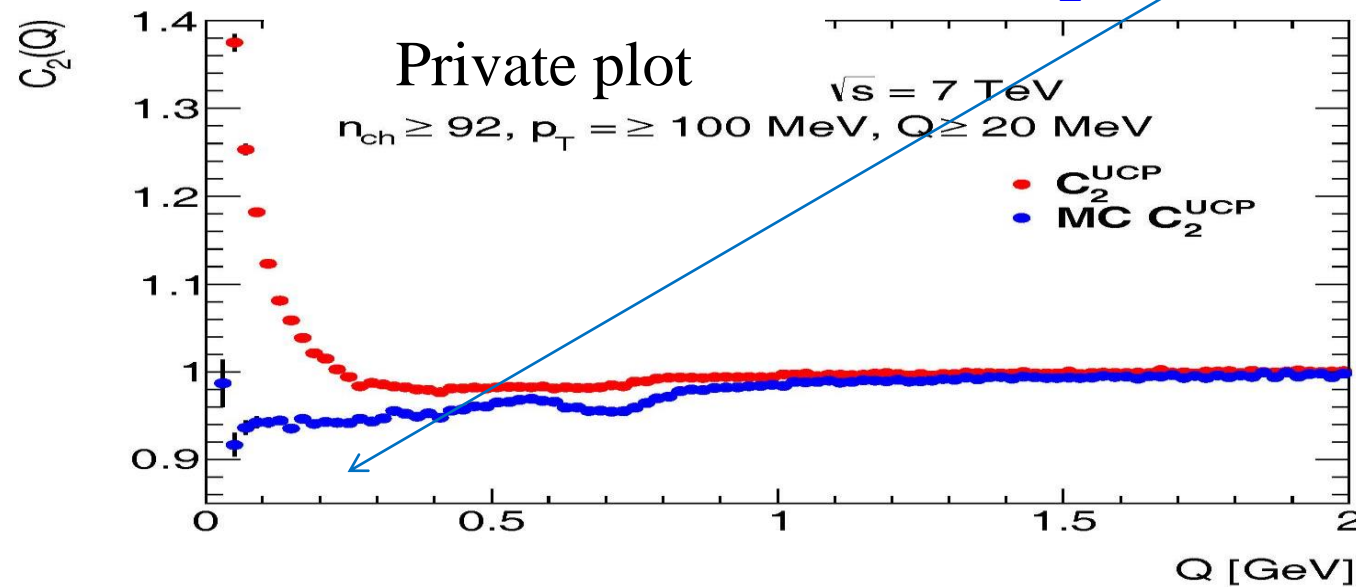
Artificial peak in C_2 in BEC region



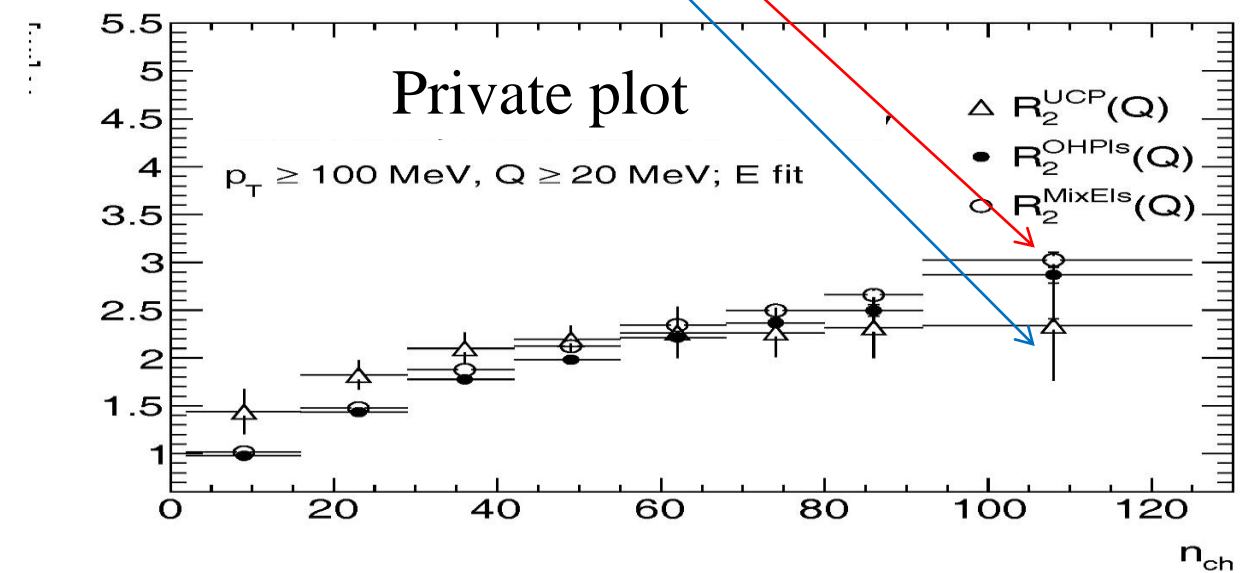
R_2 OHP & UCP



Reflection of resonances in C_2 UCP



The n_{ch} dependences for parameter R for R_2 with diff reference samples



Disadvantage: violation of energy-momentum constraint, event topology, destroying other features such as non-BEC etc.

RESULTS OF THE $R(N_{CH})$ & $R(K_T)$; AND $\lambda(N_{CH})$ & $\lambda(K_T)$ FITS

Table 4. Results of the fit for BEC parameter distributions for $\sqrt{s} = 0.9$ and 7 TeV. The $\sqrt[3]{n_{ch}}$ fit of $R(n_{ch})$ is applied to 7 TeV minimum-bias events at $n_{ch} \leq 55$ and to 0.9 TeV minimum-bias events. The constant fit of $R(n_{ch})$ is applied to 7 TeV minimum-bias events for $n_{ch} > 55$ and to 7 TeV high-multiplicity events. The exponential fit of $\lambda(n_{ch})$ is applied to 7 TeV minimum-bias and high-multiplicity events.

BEC param.	Fit function	0.9 TeV		7 TeV	
				Minimum-bias events	High-multiplicity events
$R(n_{ch})$	$p_0 \sqrt[3]{n_{ch}}$	$p_0 = 0.64 \pm 0.03 \pm 0.06$ fm		$p_0 = 0.60 \pm 0.02 \pm 0.05$ fm	
	p_0	—		$p_0 = 2.28 \pm 0.02 \pm 0.31$ fm	
$\lambda(n_{ch})$	$p_0 e^{-p_1 n_{ch}}$	$p_0 = 1.06 \pm 0.10 \pm 0.05$		$p_0 = 0.96 \pm 0.02 \pm 0.06$	
		$p_1 = 0.011 \pm 0.004 \pm 0.002$		$p_1 = 0.0039 \pm 0.0004 \pm 0.0005$	
$R(k_T)$	$p_0 e^{-p_1 k_T}$	$p_0 = 2.65 \pm 0.31 \pm 0.12$ fm		$p_0 = 2.88 \pm 0.15 \pm 0.36$ fm	
		$p_1 = 1.49 \pm 0.32 \pm 0.61$ GeV ⁻¹		$p_1 = 1.04 \pm 0.12 \pm 0.68$ GeV ⁻¹	
$\lambda(k_T)$	$p_0 e^{-p_1 k_T}$	$p_0 = 1.20 \pm 0.16 \pm 0.09$		$p_0 = 1.13 \pm 0.08 \pm 0.04$	
		$p_1 = 2.00 \pm 0.32 \pm 0.13$ GeV ⁻¹		$p_1 = 1.56 \pm 0.15 \pm 0.25$ GeV ⁻¹	