

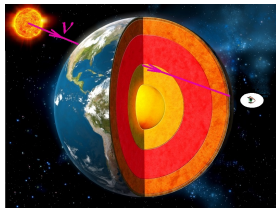
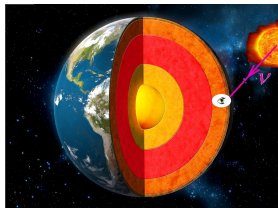
Reception of modulated neutrino day-night effect signal by neutrino detectors

Oleg Kharlanov*

*Moscow State University, Moscow, Russia

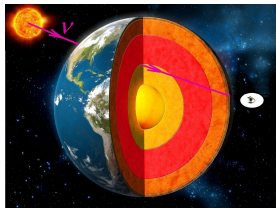
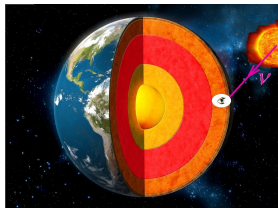
AYSS-2018, Joint Institute of Nuclear Research
23 April, 2018

Why day-night effect? Why modulation?



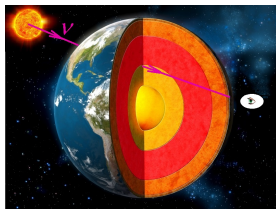
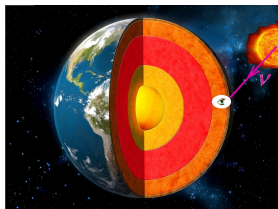
- Day-night effect (day-night asymmetry, DNA): flavor composition of nighttime solar neutrinos \neq that of daytime solar neutrinos \leftarrow regeneration inside the Earth [Carlson, 1986; Baltz, Weneser, 1986]
- Observed for ${}^8\text{B}$ ν 's at SK [PRL, 2014]: $\mathcal{A}_{\text{dn}} \sim 3 - 5\%$

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 - Extraction of DNA needs a long-term observation \sim integration over time of the effect, which, actually, depends on the nadir angle $\Theta_{\text{N}}(t)$, i.e., is modulated
- We probably lose the 'carrier' on $\int dt \dots$. Can we 'receive' it by 'smart demodulation'?

Oscillations in matter & Earth regeneration [1]



The theory of DNA is quite conventional: for each neutrino trajectory $r = r(x; \Theta_N(t))$,

$$i\lambda\partial_x R(x, x_0) = H(x)R(x, x_0), \quad R(x_0, x_0) = \mathbb{1};$$
$$H(x) = \left(-\cos 2\theta_0 + \frac{2EV(x)}{\Delta m^2} \right) \sigma_1 + \sin 2\theta_0 \sigma_3,$$

$R_{f,f'}(x, x_0) \equiv \langle \nu_f(x) | \nu_{f'}(x_0) \rangle$ is the **flavor evolution matrix** ($f, f' = e, \mu$)

$V(x) = \sqrt{2}G_F n_e(x)$ is the Wolfenstein potential $n_e(x)$ is the electron density
 $\lambda = \Delta m^2/4E = \pi/\ell_{\text{osc}}$, $\ell_{\text{osc}} \sim 300$ km E is the ν energy
 $\sin^2 2\theta_0 \approx 0.86$, $\Delta m^2 \approx 7.6 \times 10^{-5}$ eV² x goes along the ν ray

Oscillations in matter & Earth regeneration [2]

$$i\lambda\partial_x R(x, x_0) = \left\{ \left(-\cos 2\theta_0 + \frac{2EV(x)}{\Delta m^2} \right) \sigma_1 + \sin 2\theta_0 \sigma_3 \right\} R(x, x_0)$$

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In a number of papers [e.g., D'Olivo, 1992; D'Olivo *et al.*, 2008; de Holanda, Wei Liao, Smirnov, 2004], this equation was solved leading to ν_e observation probabilities

$$P_e(\text{day}) = \frac{1}{2} + \frac{1}{2} \cos 2\theta_{\text{Sun}} \cos 2\theta_0,$$

$$P_e(\text{night}) = \frac{1}{2} + \frac{1}{2} \cos 2\theta_{\text{Sun}} \left\{ \cos 2\theta_n^- + 2 \sin 2\theta_0 \sum_{j=1}^{n-1} \Delta\theta_j \cos 2\Delta\psi_{n,j} \right\}$$

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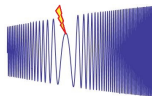
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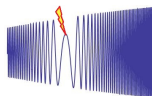
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$$\blacktriangleright \int_a^b F(t) e^{i\lambda S(t)} dt = \sqrt{\frac{2\pi i}{\lambda S''(t_0)}} F(t_0) e^{i\lambda S(t_0)} + \frac{F(t) e^{i\lambda S(t)}}{i\lambda S'(t)} \Big|_a^b + O(\lambda^{-3/2}), \lambda \rightarrow +\infty$$

The miracles of the stationary points [1]



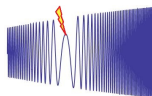
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- The contribution of the stationary point t_0 is **localized**, i.e., does not depend on the **observation time window** $[a, b]$, unless one gets under the **localization scale** δt s.t. $|S(t_0 + \delta t) - S(t_0)| \sim 2\pi/\lambda$

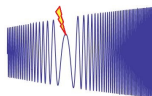
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- For the DN effect, $\int dt \rightsquigarrow \int d\zeta d\tau$ ($\zeta = \text{season}$, $\tau = \text{time of day}$), and the **stationary points** occur at **midnights** for $\int d\tau$ and on the **two solstices** for $\int d\zeta$

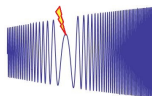
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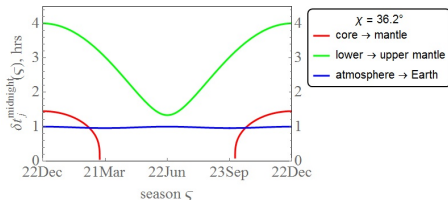


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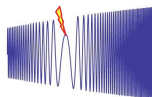
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- The winter solstice contribution gets considerably augmented for a tropical detector ($\chi \sim 23.5^\circ$) [Aleshin, O.K., Lobanov, PRD2013]
- Despite the Sun spends little time shining through the core, for the localized contribution, it may not be a problem

The miracles of the stationary points [1]: loc. scales

| Latitude χ | Interfaces crossed | r_j , km | $\Delta t_j^{\text{season}}$, month | $\delta t_{-1,j}^{\text{solstice}}$ (winter solstice), month for $E = 10$ MeV ($E = 862$ keV) | $\delta t_{+1,j}^{\text{solstice}}$ (summer solstice), month for $E = 10$ MeV ($E = 862$ keV) |
|-----------------|---------------------------|------------|--------------------------------------|---------------------------------------------------------------------------------------------------|---------------------------------------------------------------------------------------------------|
| 26.0° | atmosphere→crust | 6371 | 12.0 | 3.1 (0.9) | 0.8 (0.2) |
| | upper mantle→lower mantle | 5701 | 12.0 | 3.0 (0.9) | 0.7 (0.2) |
| | lower mantle→outer core | 3480 | 7.2 | 2.6 (0.8) | N/A |
| | outer core→inner core | 1221 | 3.3 | 1.7 (0.5) | N/A |
| | inner core→outer core | 1221 | 3.3 | 2.1 (0.6) | N/A |
| | outer core→lower mantle | 3480 | 7.2 | 4.7 (1.4) | N/A |
| | lower mantle→upper mantle | 5701 | 12.0 | 13 (3.7) | 1.7 (0.5) |
| 36.2° | atmosphere→crust | 6371 | 12.0 | 1.4 (0.4) | 0.7 (0.2) |
| | upper mantle→lower mantle | 5701 | 12.0 | 1.3 (0.4) | 0.6 (0.17) |
| | lower mantle→outer core | 3480 | 5.5 | 1.1 (0.3) | N/A |
| | outer core→lower mantle | 3480 | 5.5 | 2.0 (0.6) | N/A |
| | lower mantle→upper mantle | 5701 | 12.0 | 5.6 (1.6) | 0.9 (0.3) |
| 42.5° | atmosphere→crust | 6371 | 12.0 | 1.1 (0.3) | 0.7 (0.20) |
| | upper mantle→lower mantle | 5701 | 10.3 | 1.1 (0.3) | N/A |
| | lower mantle→outer core | 3480 | 4.4 | 0.9 (0.3) | N/A |
| | outer core→lower mantle | 3480 | 4.4 | 1.5 (0.4) | N/A |
| | lower mantle→upper mantle | 5701 | 10.3 | 4.4 (1.3) | N/A |

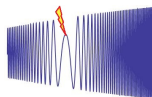


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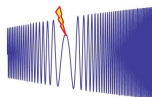
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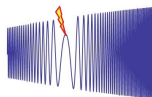


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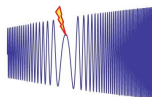
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- introduce a **weighted neutrino event number**

$$N_{\text{night}}^{(w)} = \sum_{k=1}^{N_{\text{obs}}} \vartheta(\pi/2 - \Theta_N(t_k)) w(t_k)?$$

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N.B.: the **noise** of a half-year $N_{\text{night}}^{(w)}$ is $\sqrt{2}$ times higher than that of the full-year one; the cumulative contribution to \mathcal{A}_{dn} is the same, while the localized one is multiplied by two! Thus, the **SNR** got even $\sqrt{2}$ times **better** for the localized DN effect!

Numerical simulation: definitions

The goal: magnify solstice/midnight contributions the DNA

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- **“Probabilistic” asymmetry factor** (does not depend on detector)

$$\hat{A}_{\text{dn}}^{(w)}(E) = \frac{2[\langle P(\text{night}; E) \rangle_w - P(\text{day}; E)]}{\langle P(\text{night}; E) \rangle_w + P(\text{day}; E)},$$
$$\langle P_{\nu_e}(\text{night}; E) \rangle_w \equiv \frac{\int_0^{\pi/2} P_{\nu_e}(\Theta; E) \epsilon_w(\Theta) d\Theta}{\int_0^{\pi/2} \epsilon_w(\Theta) d\Theta}$$

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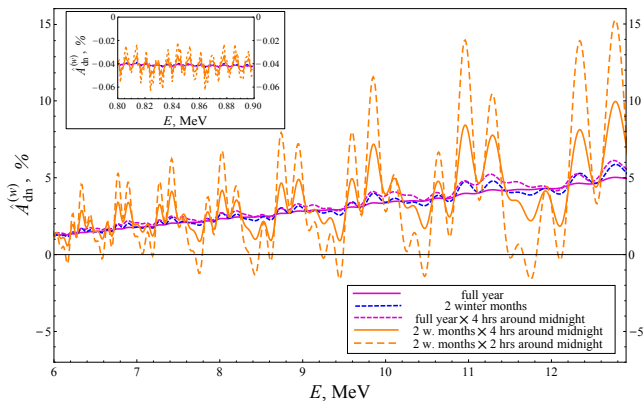
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- For a narrow electron recoil energy bin $[T, T + \Delta T]$, $\Delta T \rightarrow 0$,

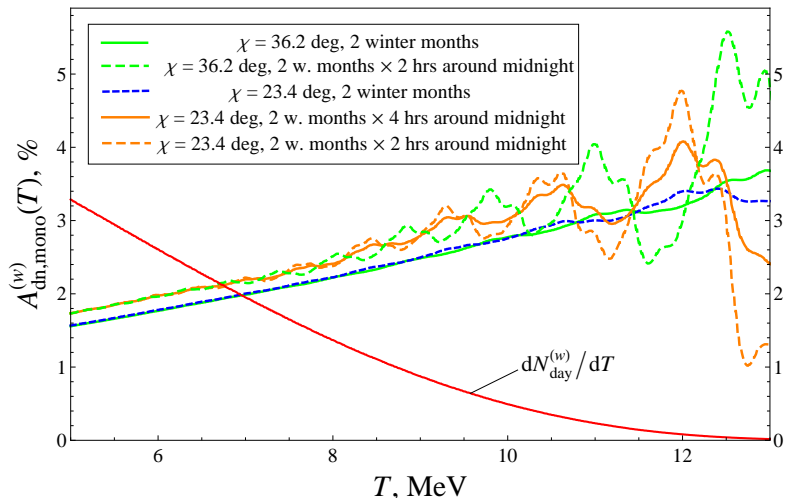
$$A_{\text{dn}}^{(w)}(T) \approx \frac{\int \Phi(E) dE \Delta \frac{d\sigma(E, T)}{dT} P_{\text{day}}(E) \hat{A}_{\text{dn}}^{(w)}(E)}{\int \Phi(E) dE \left\{ \Delta \frac{d\sigma(E, T)}{dT} P_{\text{day}}(E) + \frac{d\sigma_{\nu x}(E, T)}{dT} \right\}}$$

Numerical simulation: “probabilistic asymmetry factor”



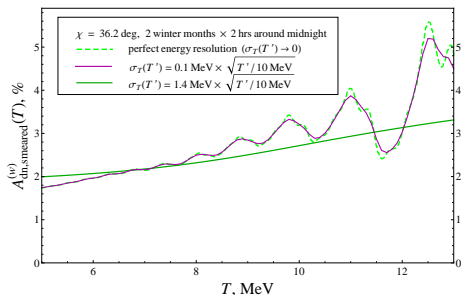
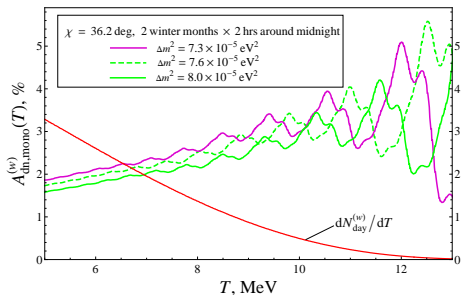
► The ‘demodulation’ of ${}^7\text{Be}$ neutrinos was studied by Ioannian, Smirnov, and Wyler [PRD2015]

Numerical simulation: electron recoil signatures



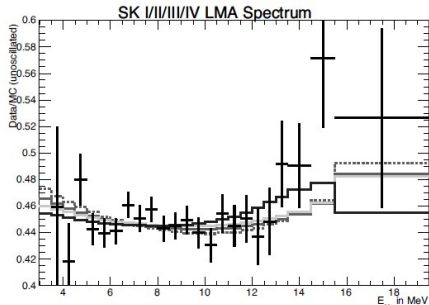
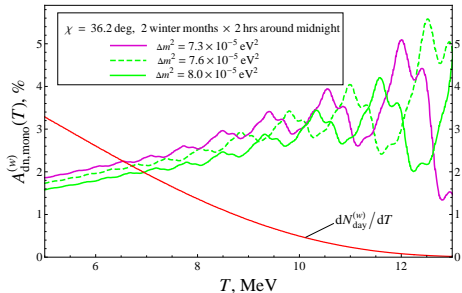
► The high-energy tail is affected; the signature depends strongly on the latitude

Numerical simulation: electron recoil signatures



► An 'interference' experiment for determining Δm^2 ?

Numerical simulation: a (possibly) rhetorical question



Measured recoil energy spectrum minus the prediction in a no-oscillations scenario [M.B. Smy for Super-Kamiokande Collaboration, *Terrestrial Matter Effects from Solar Neutrino Interactions in Super-Kamiokande*, Nucl. Part. Phys. Proceedings 265, 135 (2015)]

- ▶ The same 'wiggly' signature? Could it be made more statistically significant by temporal weighting?

Some conclusions

- Although oscillatory contributions to the physical effects are usually assumed to average out, this may be not quite so if, e.g., the oscillations have **stationary points**

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Some conclusions

- Although oscillatory contributions to the physical effects are usually assumed to average out, this may be not quite so if, e.g., the oscillations have **stationary points**
- The day-night effect has these points around midnights/solstices, and one can use their **localization** to amplify their contributions
- Signatures of these points may be present in the **high-energy tails** of the recoil energy distributions and are quite sensitive to the oscillation parameters

Acknowledgments & publications

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References

- [1] S. S. Aleshin, O. G. Kharlanov, and A. E. Lobanov, Analytical treatment of long-term observations of the day-night asymmetry for solar neutrinos, *Phys. Rev. D* **87**, 045025 (2013).
- [2] O. G. Kharlanov, Peculiar seasonal effects in the neutrino day-night asymmetry, [arXiv:1509.08073](https://arxiv.org/abs/1509.08073)[hep-ph].

Thank you for your attention!