

Hadron Modifications in Dense Nuclear Matter

G. Musulmanbekov

JINR

genis@jinr.ru

Content

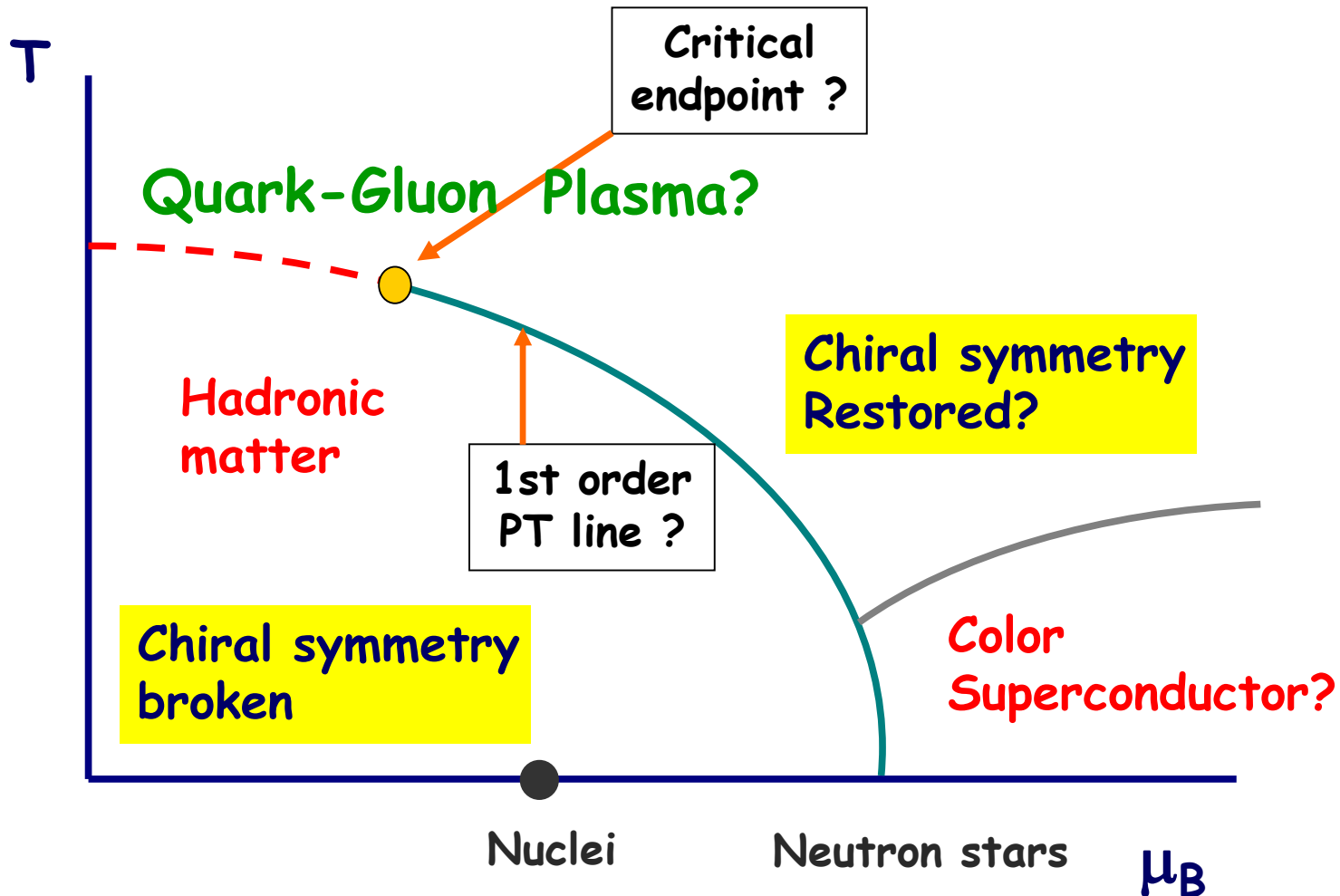
- Motivation
- Hadrons in Nuclei
 - Strongly correlated quark model (SCQM) of the hadron structure
 - Building the nuclear structure
- Hadron modifications in a dense nuclear matter
- Understanding of exp. effects in HIC
 - Enhanced strangeness production
 - Horn-effect
 - Enhancement of dilepton mass spectra in the range 0.2 – 0.6
- Conclusion

Motivation

- How nuclear matter behave under high compression?
- How hadron structures are modified in a dense matter?
- What observables are the possible signals of these modifications?

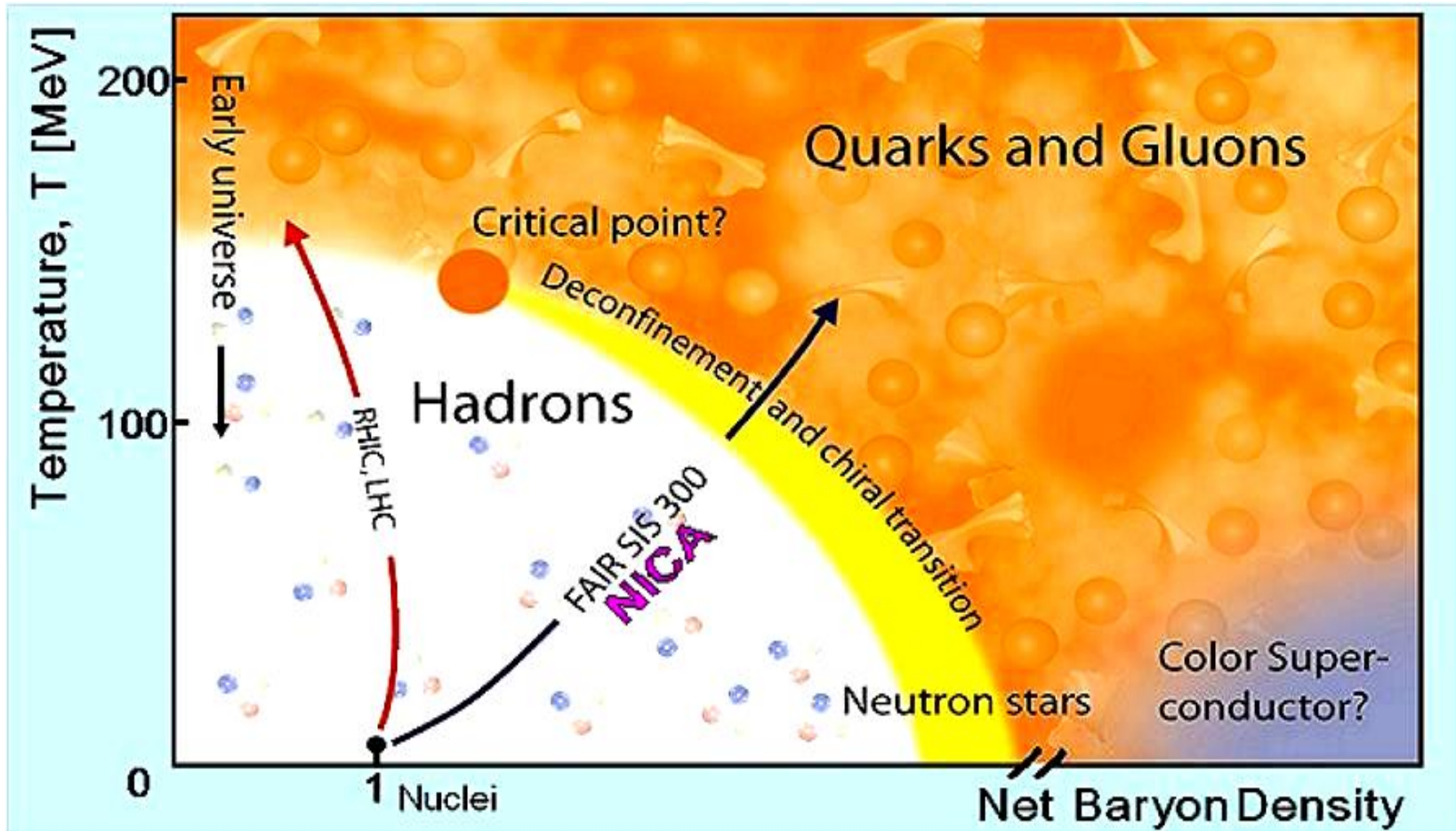
Motivation

- Does hadronic matter transit into QGP?



Motivation

- Does hadronic matter transit into QGP?



Toy Model:

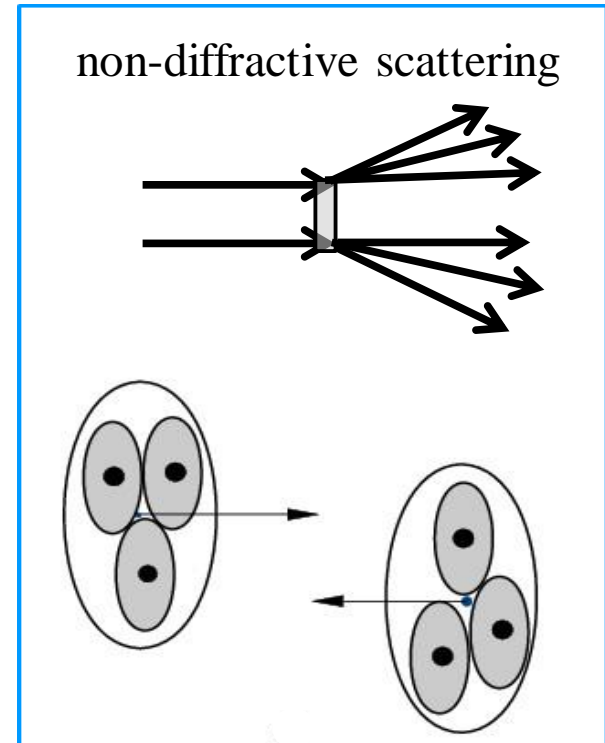
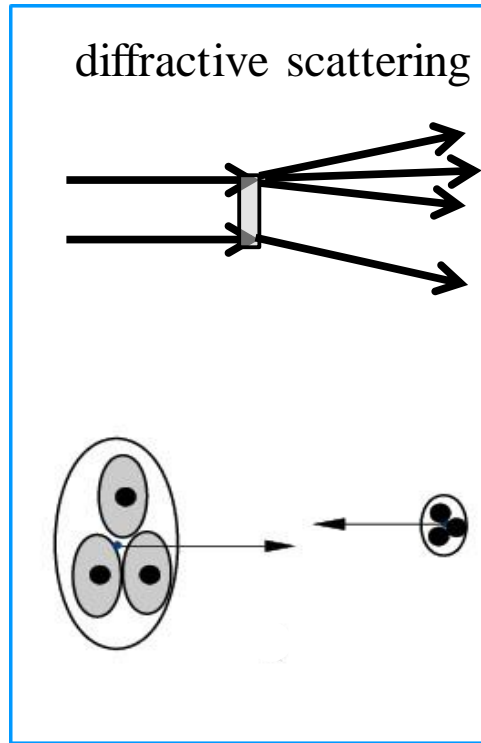
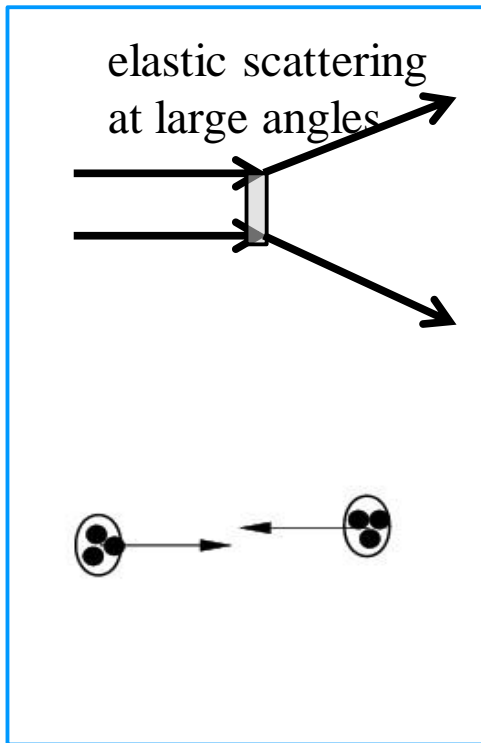
Strongly Correlated Quark Model

G. Musulmanbekov, 1995

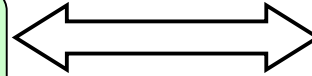
Strongly Correlated Quark Model

Motivations

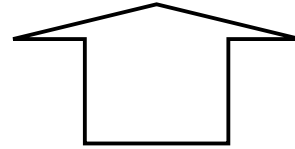
pp – interactions at high energies



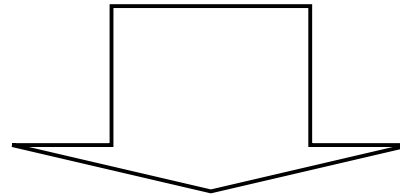
Constituent Quarks



Current Quarks

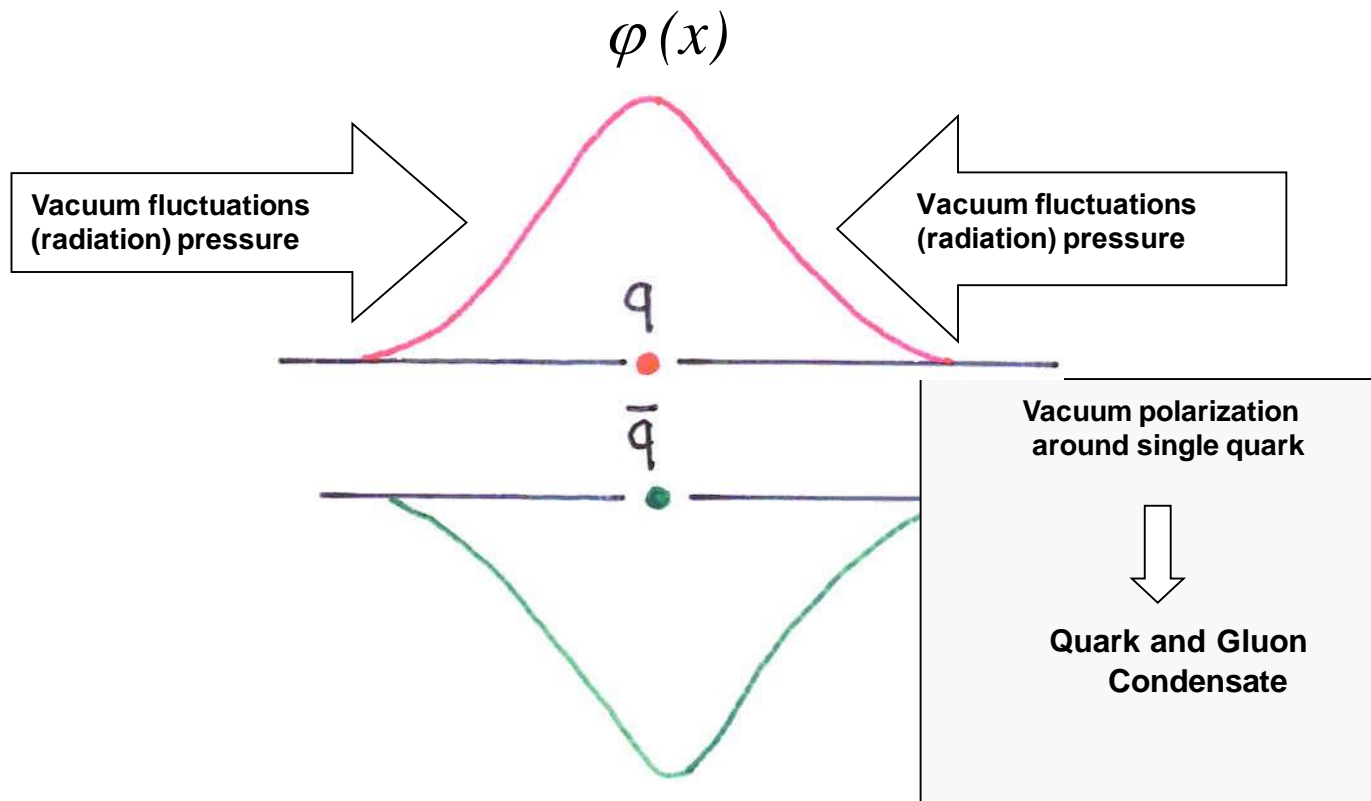


Quarks – Solitons!

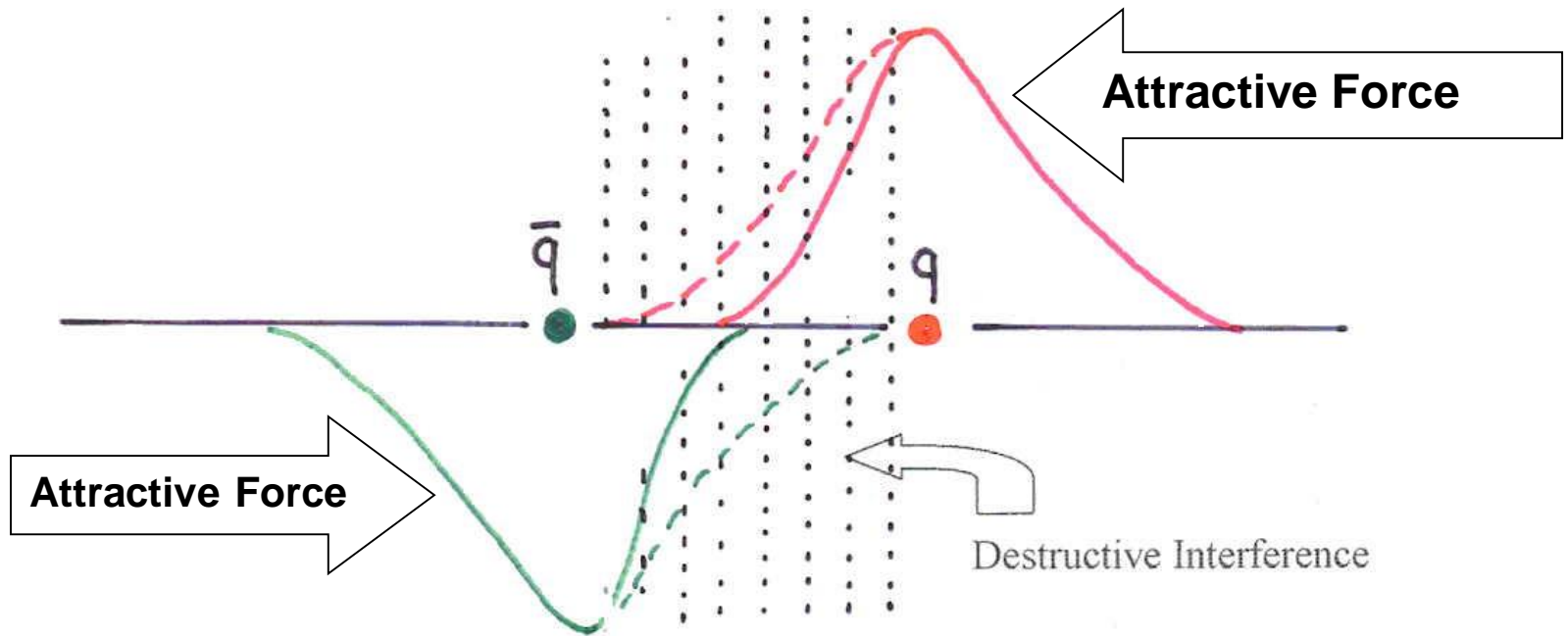


Strongly Correlated Quark Model

Strongly Correlated Quark Model (SCQM)



Strongly Correlated Quark Model (SCQM)



The Strongly Correlated Quark Model

Hamiltonian of the Quark – AntiQuark System

$$H = \frac{m_q^-}{(1 - \beta_q^{-2})^{1/2}} + \frac{m_q}{(1 - \beta_q^2)^{1/2}} + V_{qq}^-(2x)$$

m_q^- , m_q are the current masses of quarks,
 $\beta = \beta(\mathbf{x})$ – the velocity of the quark (antiquark),
 V_{qq}^- - is the quark–antiquark potential.

$$H = \left[\frac{m_q^-}{(1 - \beta_q^{-2})^{1/2}} + U(x) \right] + \left[\frac{m_q}{(1 - \beta_q^2)^{1/2}} + U(x) \right] = H_q^- + H_q$$

$U(x) = \frac{1}{2} V_{qq}^-(2x)$ is the potential energy of a single quark/antiquark.

Constituent Quarks – Topological Solitons

SCQM \equiv Breather Solution of Sine- Gordon equation

$$\partial_{\mu} \partial^{\mu} \phi(x, t) + \sin \phi(x, t) = 0$$

Breather – oscillating soliton-antisoliton pair:

$$\phi(x, t)_{s-as} = 4 \tan^{-1} \left[\frac{\sinh\left(ut / \sqrt{1-u^2}\right)}{u \cosh\left(x / \sqrt{1-u^2}\right)} \right]$$

$$\phi(x, t)_{s-as} = \frac{\partial \phi(x, t)_{s-as}}{\partial x}$$

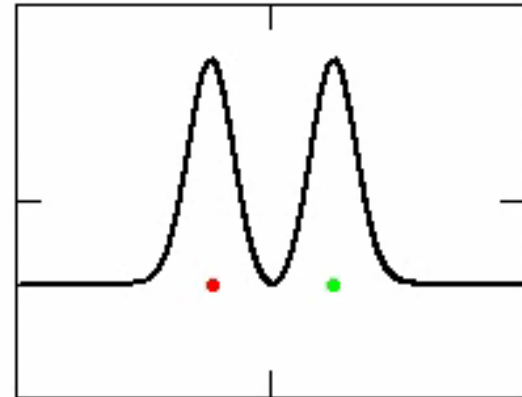
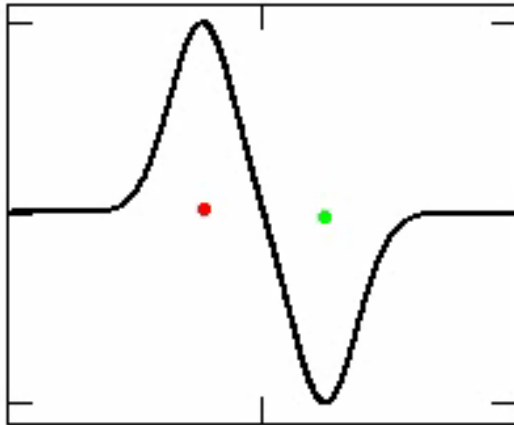
is identical to our quark-antiquark system;

Breather – quark-antiquark pair Meson

$\varphi(x,t)$

$\varepsilon(x,t)$

www.Bandicam.com



What is Chiral Symmetry and its Breaking?

- Chiral Symmetry

$$SU(2)_L \times SU(2)_R \quad \text{for } \psi_{L,R} = u, d$$

- The order parameter for symmetry breaking is quark or *chiral* condensate:

$$\langle \psi \psi \rangle \simeq - (250 \text{ MeV})^3, \quad \psi = u, d$$

- As a consequence massless valence quarks (u, d) acquire dynamical masses which we call constituent quarks

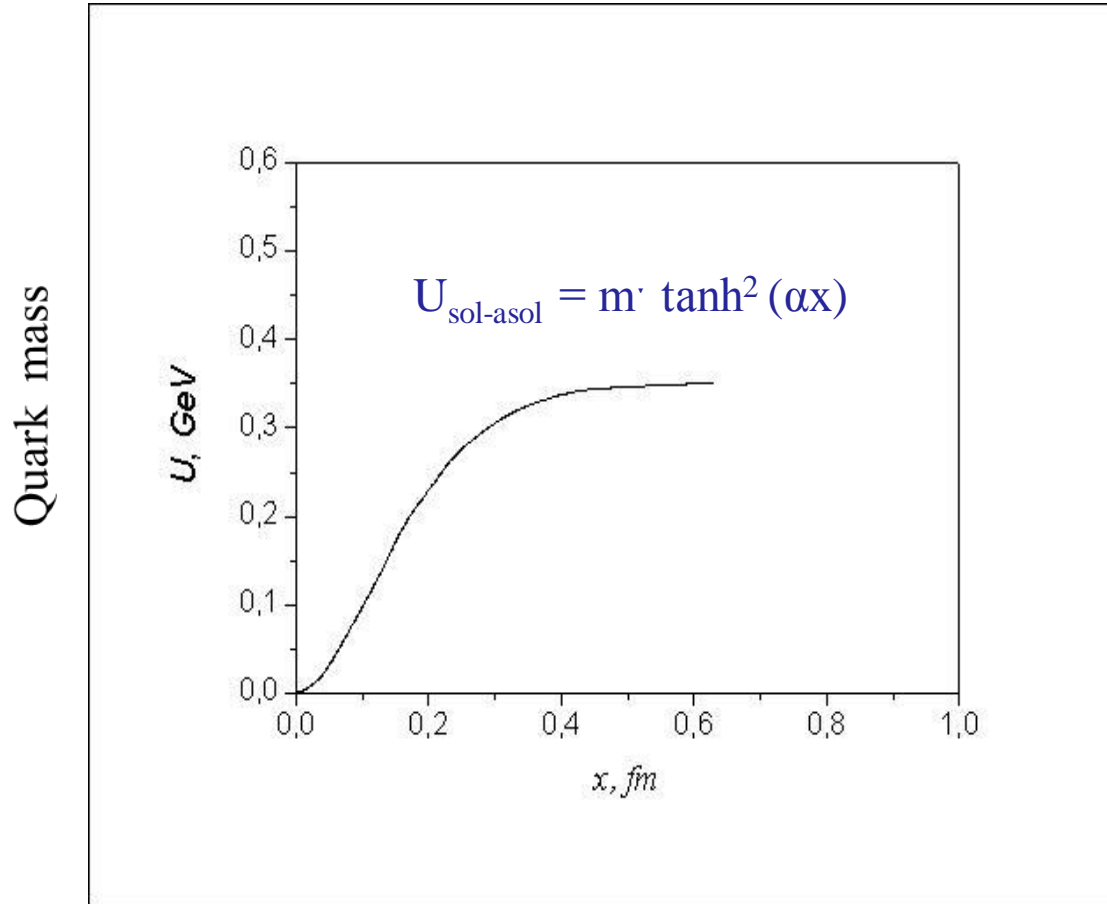
$$M_C \approx 350 - 400 \text{ MeV}$$

Quark Potential

Potential in soliton-antisoliton system: $U_{\text{sol-asol}} = m \cdot \tanh^2(\alpha x)$

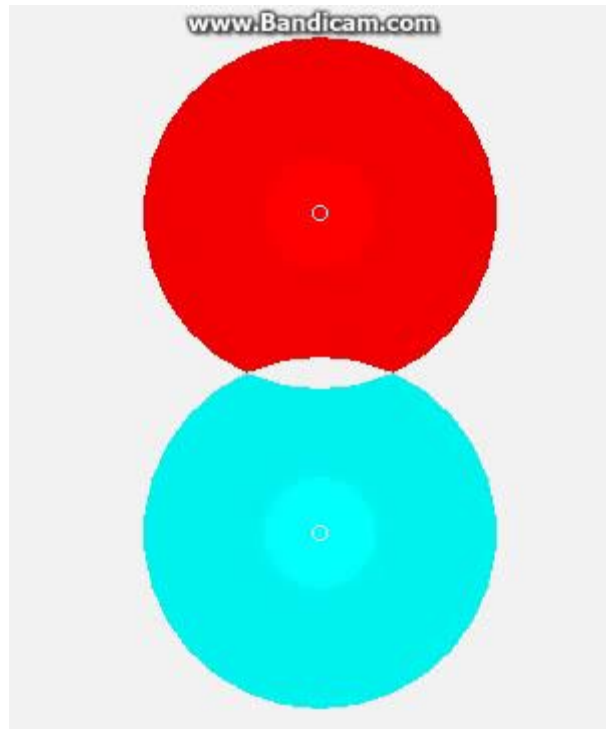
W. Troost, CERN Report, 1975;

P. Vinsareilly, Acta Phys. Aust. Suppl., 1976



quark-antiquark pair

Meson



Generalization to the 3 – quark system (baryons)

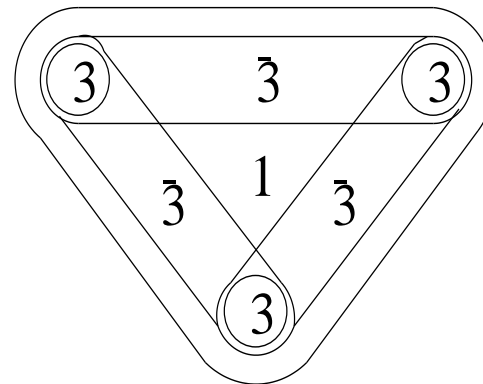
$SU(3)_{Color}$

$$q \Rightarrow SU(3) \Leftrightarrow RGB \quad \bar{q} \Rightarrow SU(\bar{3}) \Leftrightarrow CMY$$

$$\bar{q}q \Rightarrow \left(\begin{array}{ccc} \textcircled{\bar{3}} & 1 & \textcircled{3} \end{array} \right)$$

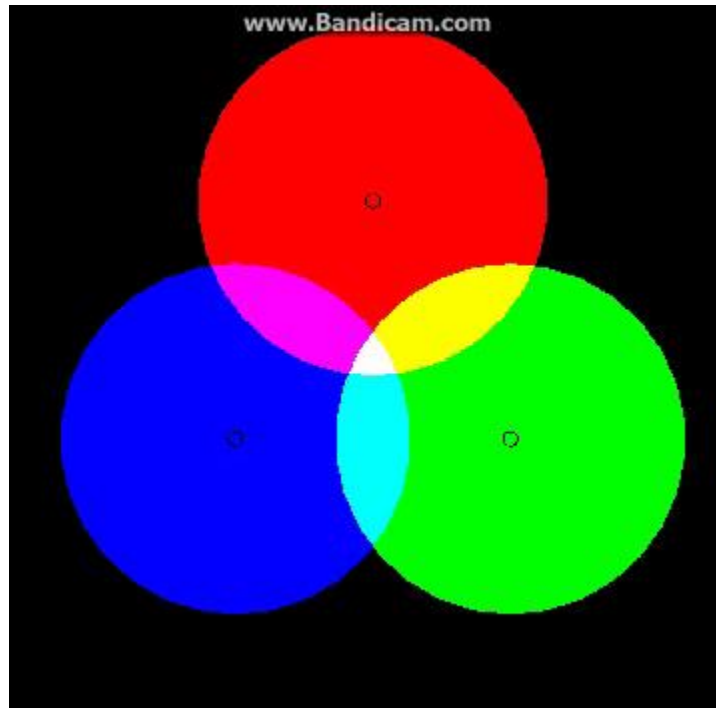
$$qq \rightarrow 3 \times 3 = 6 \oplus \bar{3} \quad \Rightarrow \quad \bar{q} \rightarrow qq$$

$$qqq \Rightarrow$$



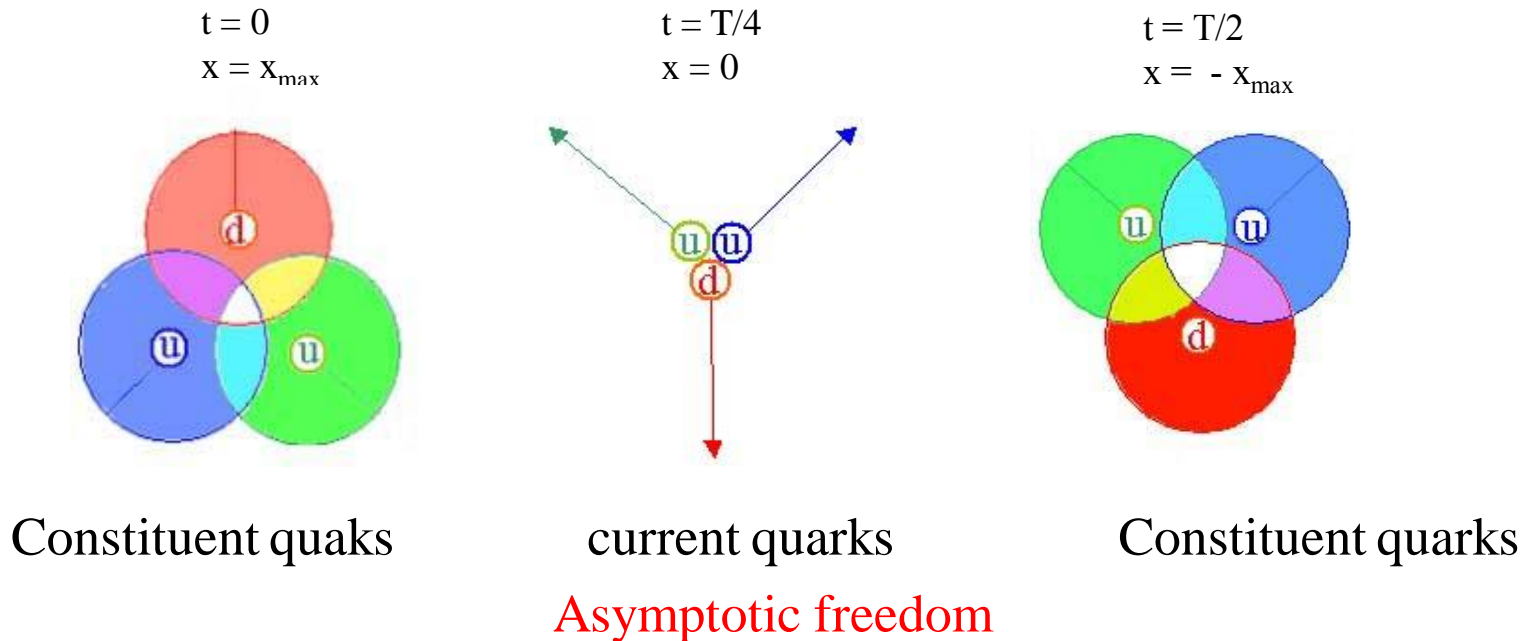
Nucleon - $SU(3)_{\text{color}}$ singlet

$SU(3)_{\text{color}}$ - RGB



Interplay between constituent and current quark states

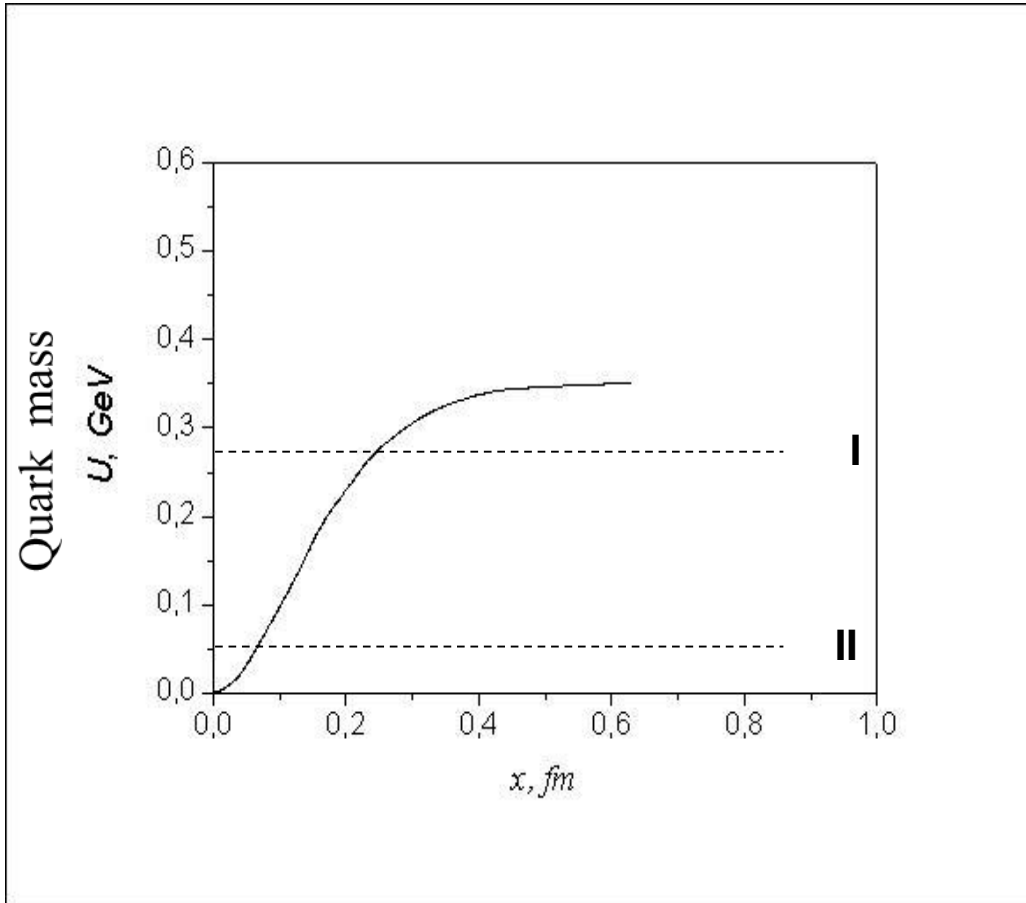
Chiral Symmetry Breaking \longleftrightarrow Restoration



During the valence quarks oscillations:

$$|B\rangle = a_1 |q_1 q_2 q_3\rangle + a_2 |q_1 q_2 q_3 \bar{q} q\rangle + a_3 |q_1 q_2 q_3 g\rangle + \dots$$

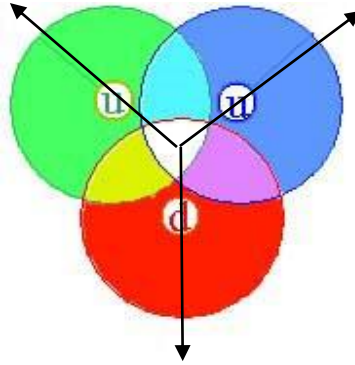
Quark Potential



$U(x) > I$ – constituent quarks

$U(x) < II$ – current (relativistic) quarks

Nucleon



Quark color wave function

$$\psi(x)_{Color} = \sum_{i=1}^3 a_i(x) |c_i\rangle$$

Where $|c_i\rangle$ are orthonormal states with $i, j = R, G, B$

$$\langle c_i | c_j \rangle = \delta_{ij}$$

Nucleon wave function

$$\psi(x)_{Color} \rightarrow \frac{1}{\sqrt{6}} \sum_{ijk} e_{ijk} |c_i\rangle |c_j\rangle |c_k\rangle$$

Parameters of SCQM for the Nucleon

- Parameters of Quark potential $U_{\text{sol-asol}} = m \cdot \tanh^2(\alpha x)$

1. Mass of Constituent Quark

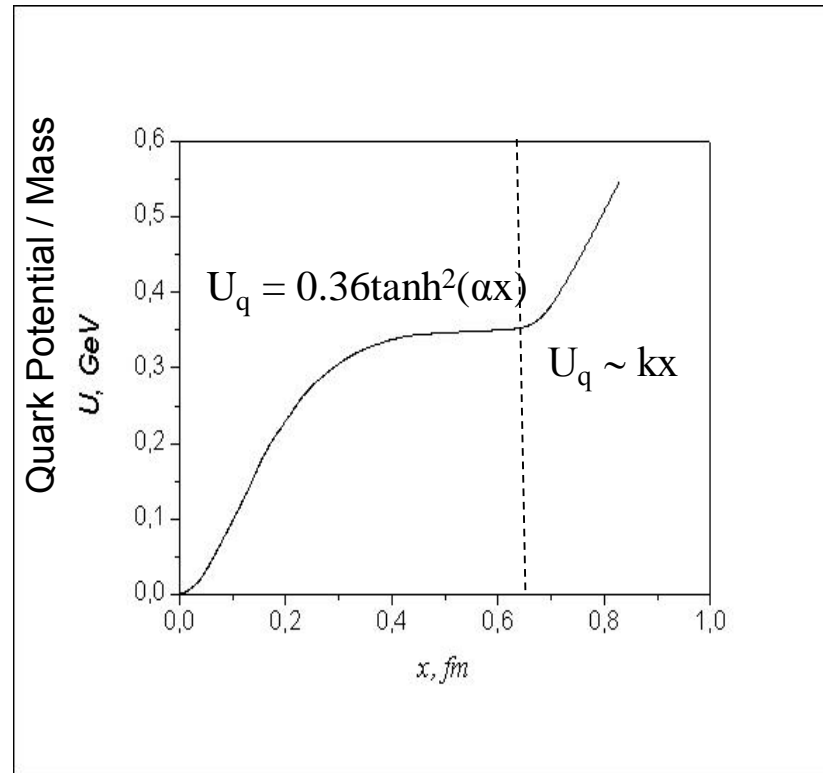
$$m = M_{Q(\bar{Q})}(x_{\text{max}}) = \frac{1}{3} \left(\frac{m_{\Delta} + m_N}{2} \right) \approx 360 \text{ MeV},$$

2. Amplitude of VQs oscillations : $\alpha = x_{\text{max}} = 0.64 \text{ fm}$,

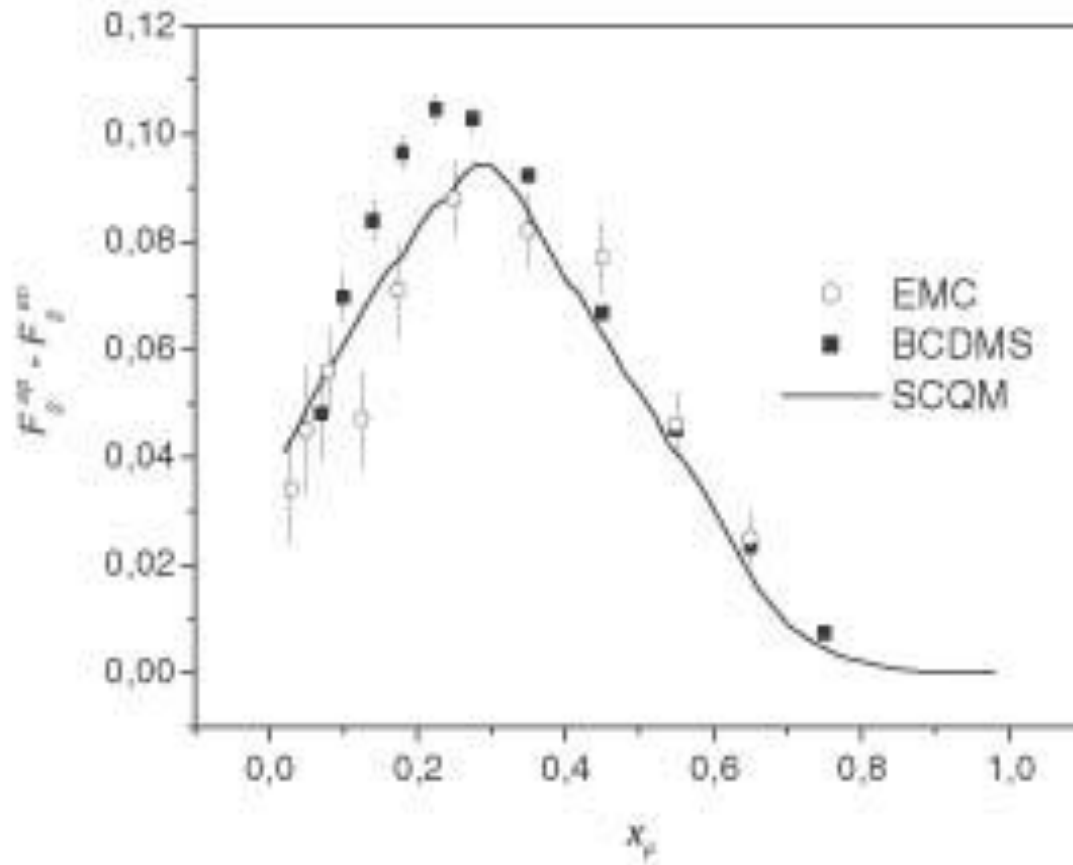
- Constituent quark sizes (parameters of gaussian distribution): $\sigma_{x,y} = 0.24 \text{ fm}$, $\sigma_z = 0.12 \text{ fm}$
- Parameters 2 and 3 are derived from the calculations of Inelastic Overlap Function (IOF) and σ_{tot} in $\bar{p} p$ and pp – collisions.

“The wave packet solution of time-dependent Schrodinger equation for harmonic oscillator moves in exactly the same way as corresponding classical oscillator”
E. Schrodinger, 1926

Quark Potential



Structure Function of Valence Quarks in Proton



SCQM \implies The Local Gauge Invariance Principle

Destructive Interference of color fields \equiv Phase rotation of the quark w.f. in color space:

$$\psi(x)_{Color} \rightarrow e^{ig\theta(x)}\psi(x)$$

Phase rotation in color space \implies quark dressing (undressing) \equiv the gauge transformation

$$A^\mu(x) \rightarrow A^\mu(x) + \partial^\mu\theta(x)$$

Therefore, during quark oscillation its

color charge

momentum

mass

**are continuously varying functions
of time.**

Relation SCQM to QCD

We reduce interaction of color quarks via **non-Abelian** fields to its **E-M** analog:

$$A_a^\mu(x) \rightarrow A^\mu(x)$$

$$F_a^{\mu\nu} = \partial^\mu A_a^\nu - \partial^\nu A_a^\mu - \lambda f^{abc} A_b^\mu A_c^\nu \rightarrow F_{ch}^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu$$

Spin in SCQM

Conjecture: spin of constituent quark is entirely analogous to the angular momentum carried by classical circularly polarized wave:

$$\mathbf{J}_Q = \mathbf{J}_g = \int_a^\infty d^3r [\mathbf{r} \times (\mathbf{E} \times \mathbf{B})]$$

Classical analog of electron spin – *F. Belinfante 1939; R. Feynman 1964; H. Ohanian 1986; J. Higbie 1988.*

Electron surrounded by proper \mathbf{E} and \mathbf{B} fields creates circulating flow of energy:

$$\mathbf{S} = \epsilon_0 \mathbf{c}^2 \mathbf{E} \times \mathbf{B}$$

Total angular momentum created by this Poynting's vector

$$\mathbf{s} = \mathbf{L} = (\dots) \int_a^\infty d^3r [\mathbf{r} \times (\mathbf{E} \times \mathbf{B})]$$

is associated with the entire spin angular momentum of the electron.

Spin in SCQM

1. Now we accept that

$$A^\mu = \{\varphi, \mathbf{A}\}$$

and intersecting \mathbf{E}_{ch} and \mathbf{B}_{ch} create around VQ color analog of Pointing's vector (circulating flow of energy)

$$\mathbf{S} = \epsilon_0 \mathbf{c}^2 \mathbf{E}_{\text{ch}} \times \mathbf{B}_{\text{ch}}.$$

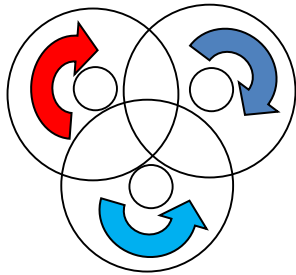
2. Total angular momentum created by this Pointing's vector

$$\mathbf{s}_Q = \mathbf{L}_g = (\dots) \int_a^\infty d^3 r [\mathbf{r} \times (\mathbf{E}_{\text{ch}} \times \mathbf{B}_{\text{ch}})]$$

is associated with the intrinsic spin of the constituent quark.

Quarks – Oscillating Vortices

$t = 0$
 $x = x_{max}$



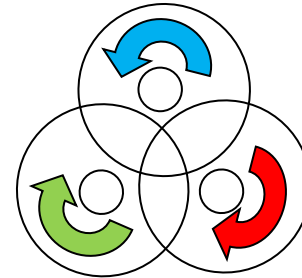
Constituent Quarks

$t = T/4$
 $x = 0$



Current Quarks

$t = T/2$
 $x = x_{max}$



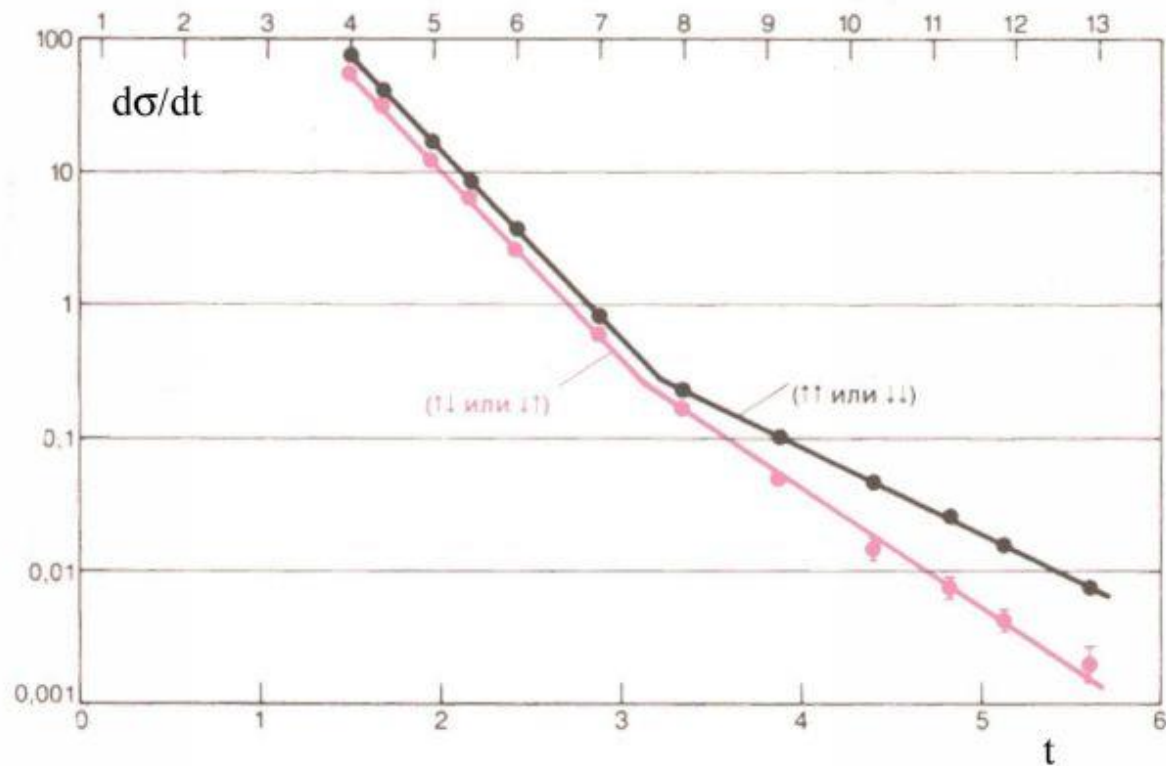
Constituent Quarks

- In the current quark state E_{ch} and B_{ch} are concentrated in a **small radius shell** around VQ.
- And so is for the vortices around VQs.

Polarized proton-proton collisions

— anti-parallel spins

— parallel spins

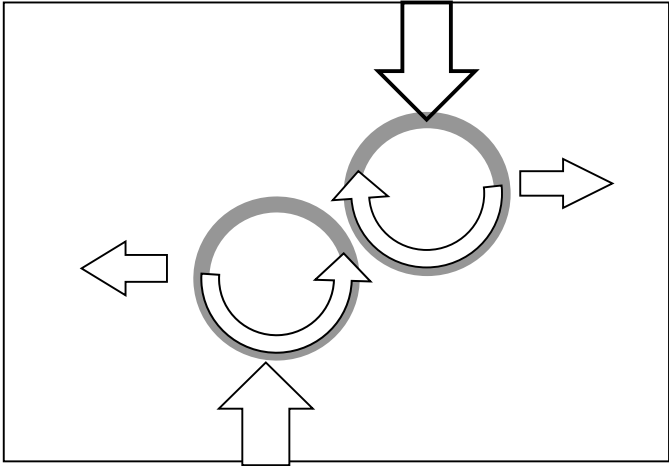


“Krish” - effect

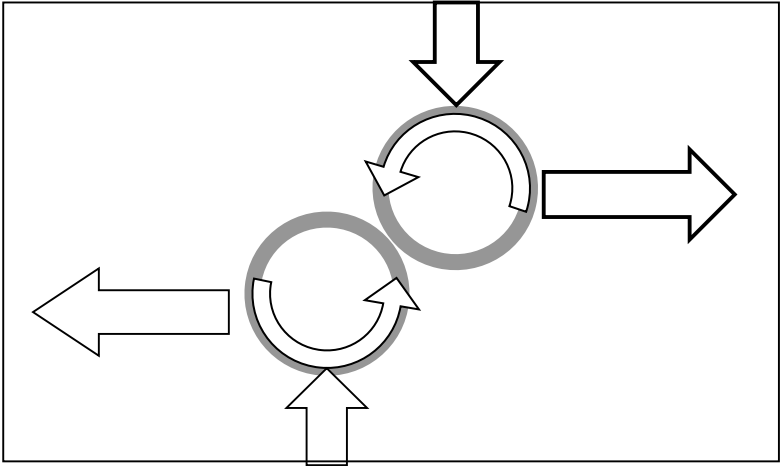
Collision of Vorticing Quarks

Spin Asymmetries

Anti-parallel Spins



Parallel Spins



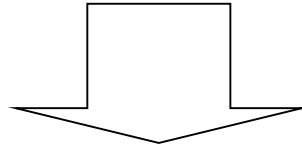
“Krish” - effect

Summary on Quarks in Hadrons

- Constituent quarks are identical to solitons.
- Quarks and gluons inside nucleons are strongly correlated;
- Hadronic matter distribution inside hadrons is fluctuating quantity;
- Nucleons are not spherically symmetric (deformed).

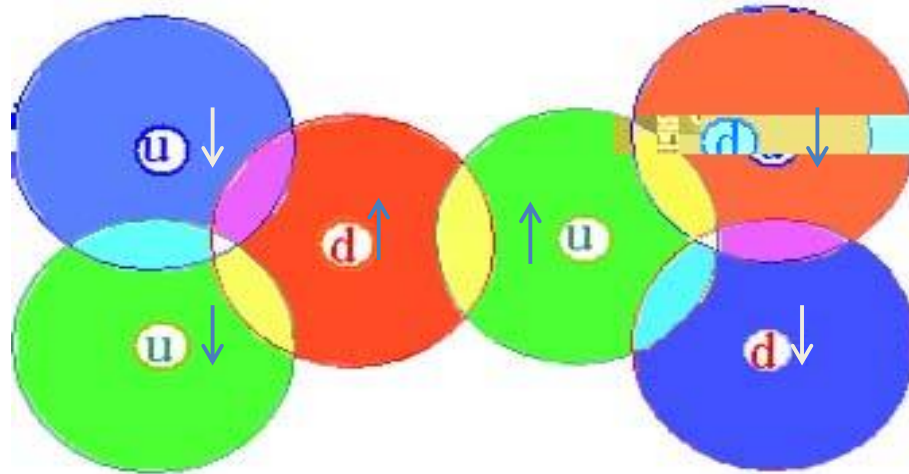
Quark Arrangement inside Nuclei

Strongly Correlated Quark Model



Crystal-like arrangement of Nuclear Structure

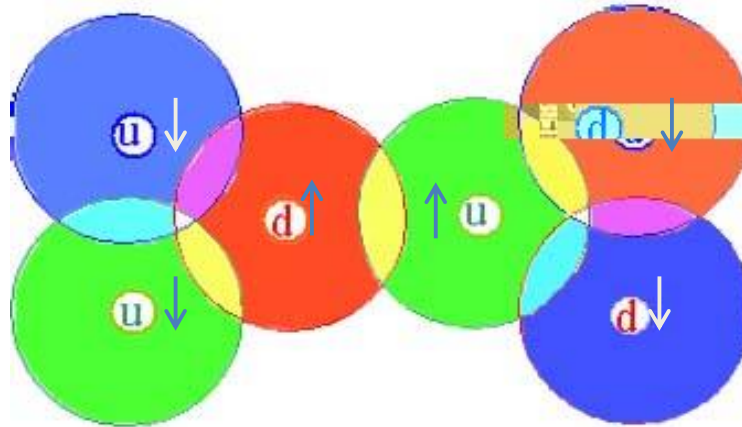
Two Nucleon System in SCQM



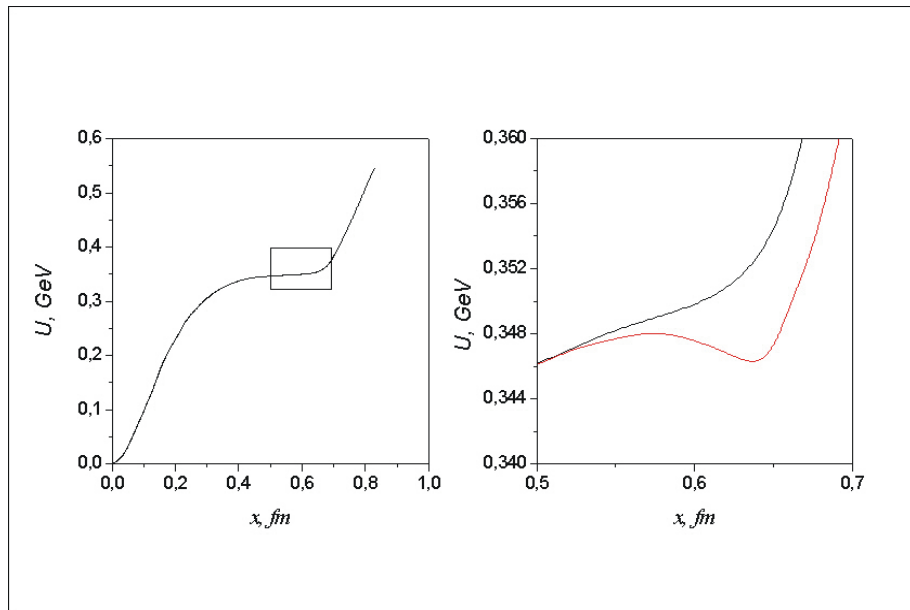
Selection rules for binding two quarks of neighboring nucleons at a junction:

- $SU(3)_{\text{Color}}$ – of different colors
- $SU(2)_{\text{Flavor}}$ – of different flavors
- $SU(3)_{\text{Spin}}$ – of parallel spins

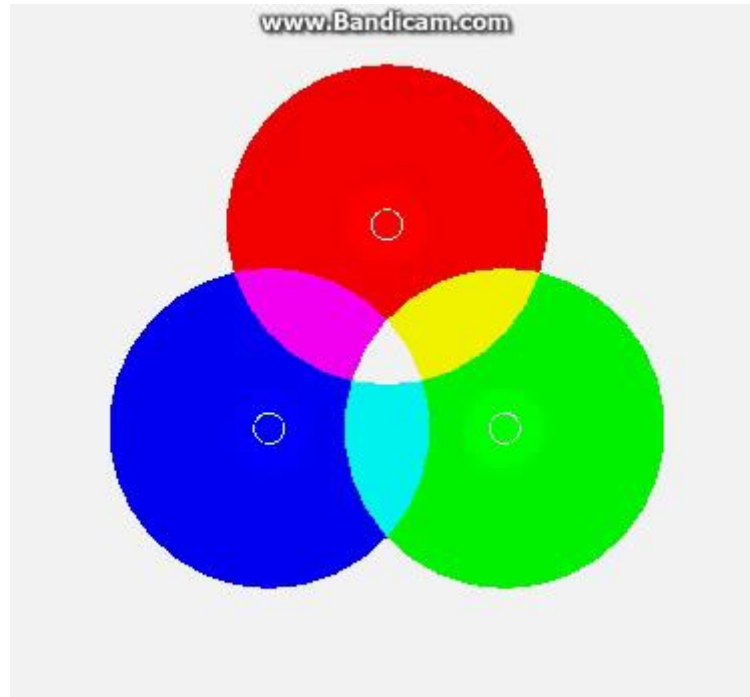
Two Nucleon System in SCQM



Quark Potential Inside Nuclei

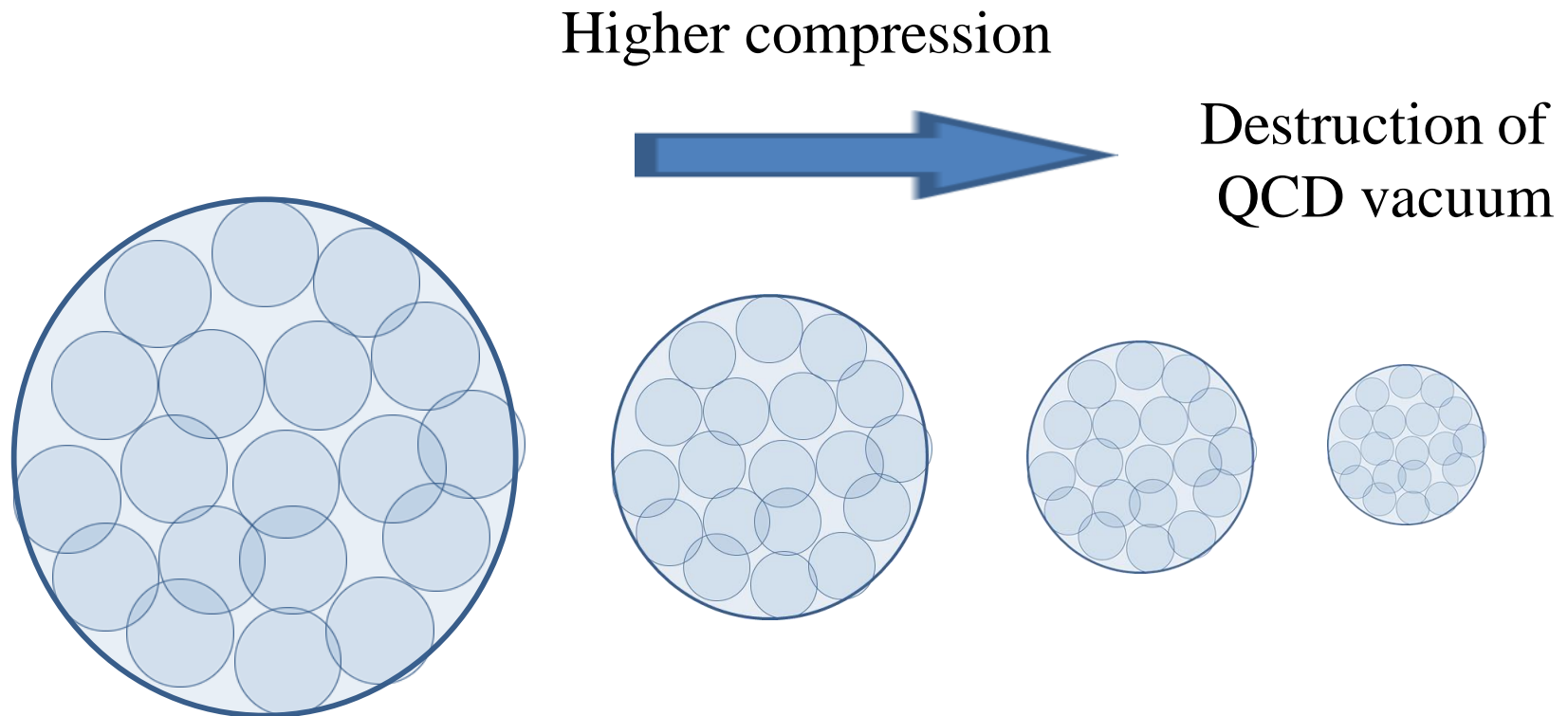


Quarks inside nucleus



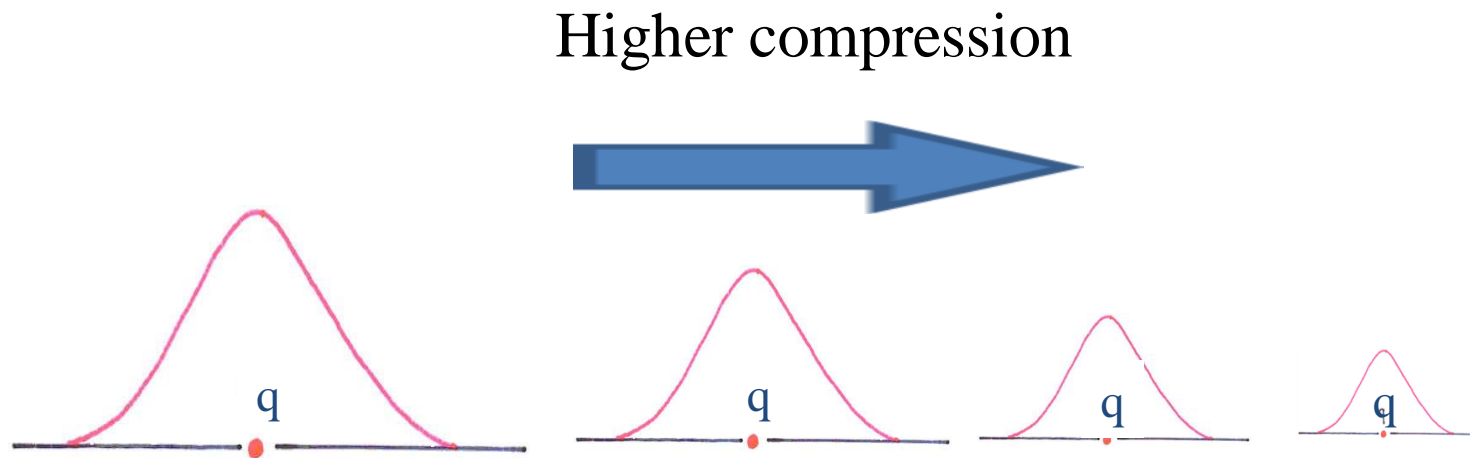
Hadron modifications in a dense nuclear matter

1. Baryonic matter under compression



Hadron modifications in a dense nuclear matter

1. Baryonic matter under compression

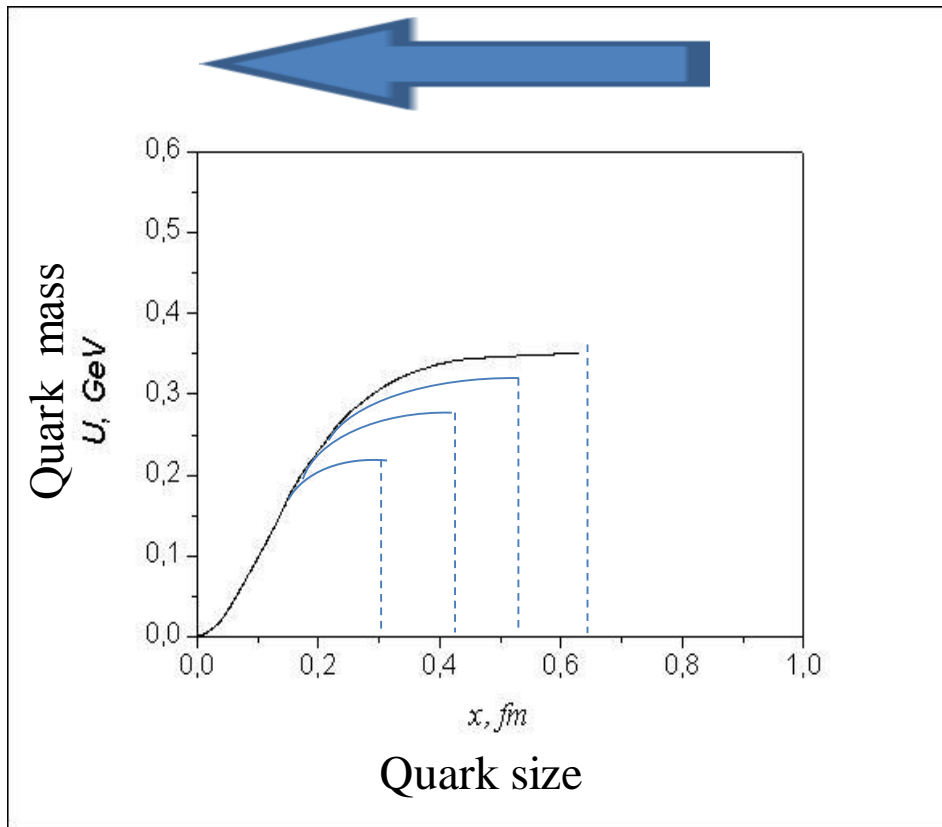


Quark condensate around valence quarks, $\langle \psi \psi \rangle$

Hadron modifications in a dense nuclear matter

1. Baryonic matter under compression

Higher compression



Decreasing nucleon **dimension**



Decreasing quark **condensate**



Decreasing quark/nucleon **mass**



Up to repulsive nucleon core

Baryonic matter under compression

Higher compression makes nucleons to convert into

- **delta-isobars**
- **hyperons**
- **higher mass resonances**

Nucleon Transition into Hyperon Phase

How can nucleons be converted into hyperons?

- Inside highly compressed nuclear matter a strange quark-antiquark condensate is created.

And:

- u and d quarks in nucleons are replaced by s -quarks,
- s -antiquarks together with those u and d form kaons:
 $p, n \rightarrow \text{hyperons} + \text{kaons}$
- the **heavier** quark content of a baryon, the **less** spatial dimensions it occupies

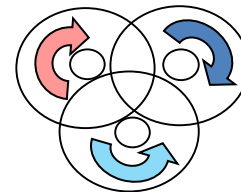
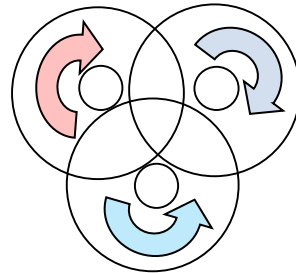
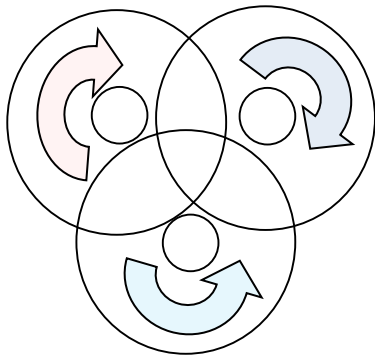
Scenario to avoid collapse

Higher compression



$n, p \Rightarrow \Delta$

$u, d \Rightarrow s, c, \dots \quad n, p \Rightarrow \Lambda, \Sigma, \Xi, \Omega, \dots$



Hadronic gas

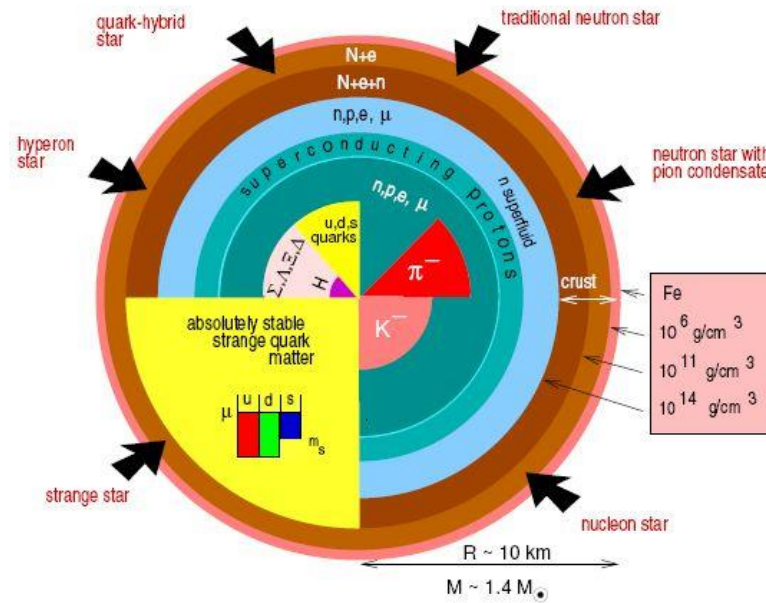


Hadronic liquid

Neutron star

Gravitational
compression

NS core



compression

Neutrons?

Δ - isobars

Hyperons

...

Hadron modifications in a dense nuclear medium

1. Hadronic matter at high density and temperature

Particle production in a hot and dense fireball

- π -production **is suppressed**
- vector mesons: $\rho, \omega, \varphi, K^*, \dots$ - **incompressible** (effective **cores**)
- ρ, ω - ‘**melting**’: mass dropping and width-widening;
dilepton spectra
- Fireball ‘cooling’ \rightarrow increased π - yield

Hadrons in a high dense and temperature medium

Model Consequences

1. **Baryons** transform to isobars then to hyperons
2. π -**production** is suppressed
3. Particle generation inside hot and dense fireball is realized mainly via **vector mesons** ρ , ω , φ , K^* , ...
4. ρ , ω – ‘**melting**’: mass dropping and width-widening;
5. **Fireball evolution:**
Hadron-Resonance Liquid \rightarrow Hadron-Resonance Gas

Hadrons in a high dense and temperature medium

1. Hadrons – topological solitons?

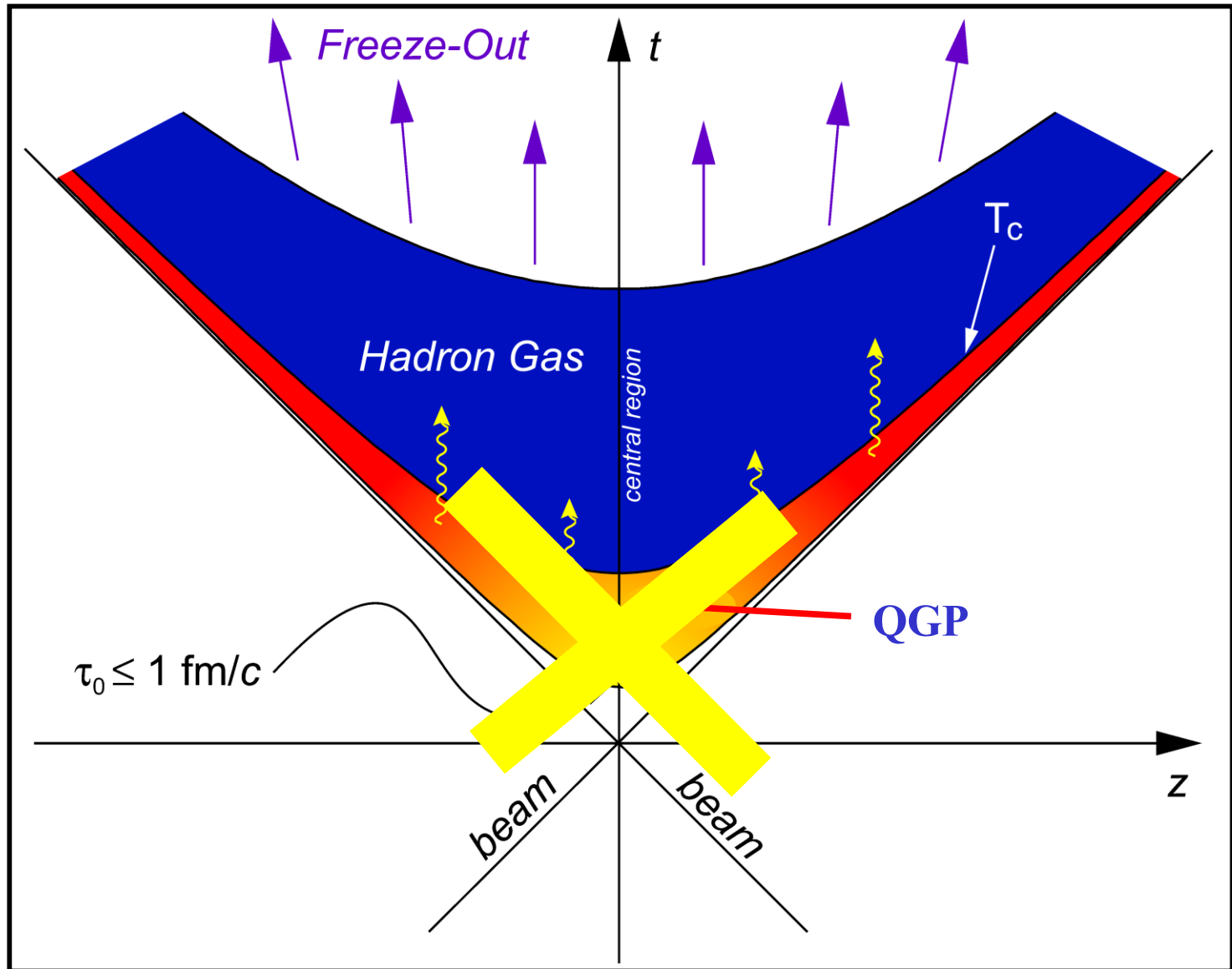


2. Conservation of topological charge

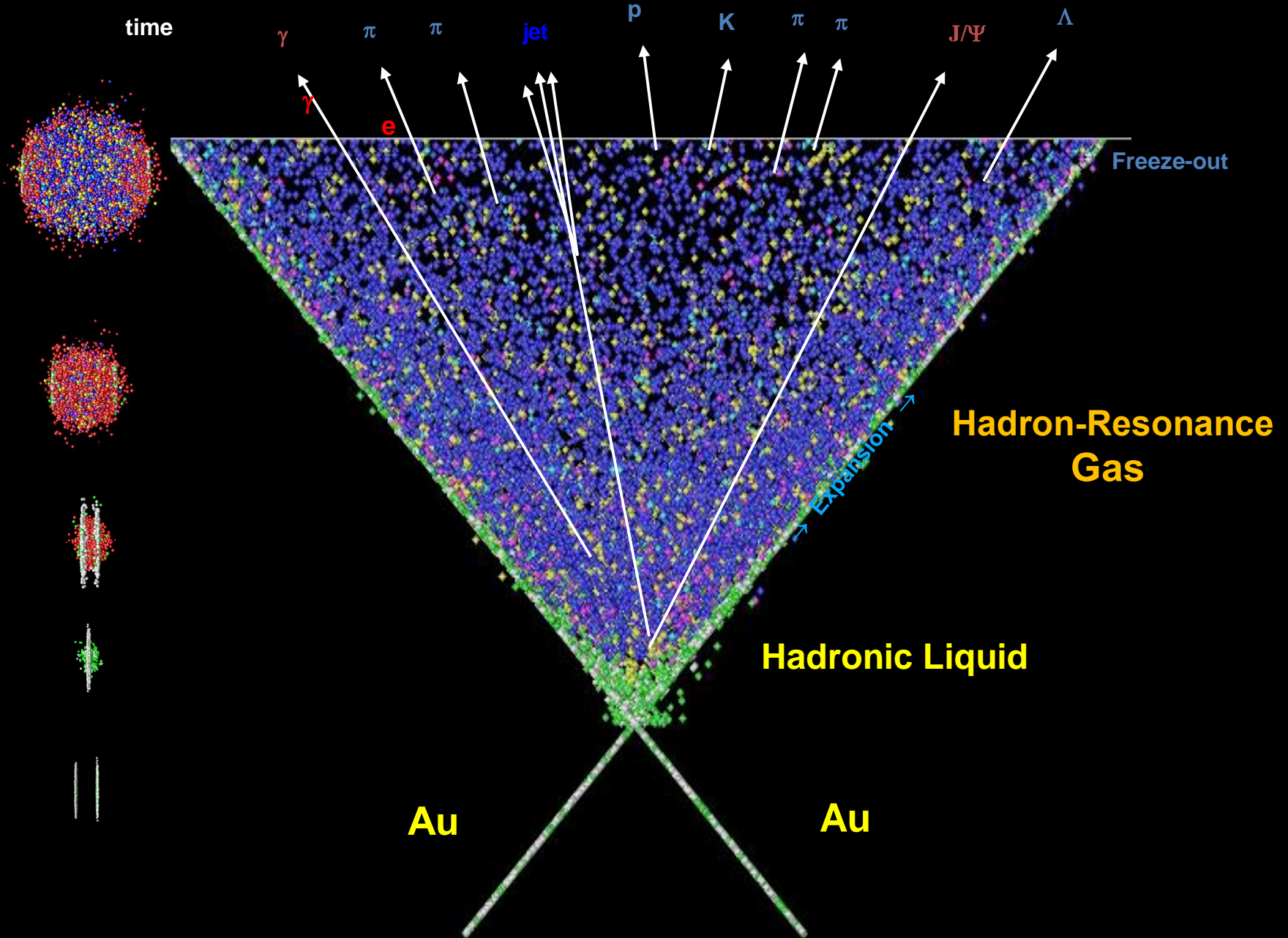


3. Deconfinement **is forbidden** → no room for QGP

Space-time Evolution of HIC



Space-time Evolution of HIC



Experiments

Energy range: $\sqrt{s} = 3 - 11 \text{ GeV}$ *most interesting!*

- **Enhanced yield of K^+ , ϕ , (multi)strange baryons**

experiments: KaoS, AGS, NA49 at low energies of SPS, (BES RHIC)

- **Horn-effect – irregular behaviour of K^+/π^+**

experiments: NA49, STAR (BES RHIC)

- **Dilepton production**

experiments: DLS, HADES, CERES, PHENIX

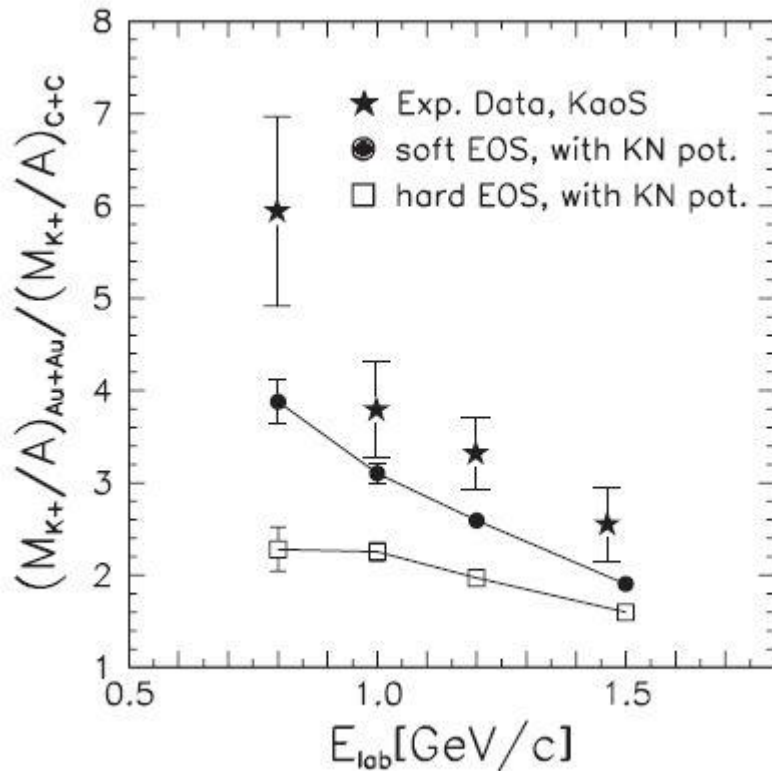
- **Projects:** FAIR/CBM, NICA/MPD, BM&N

Enhanced yield of K^+ in subthreshold kaon production

KaoS at SIS

Transport models with NN-interactions

- **underestimate** yield of K^+ by a **factor of 6**
- **overestimate** yield of K^-



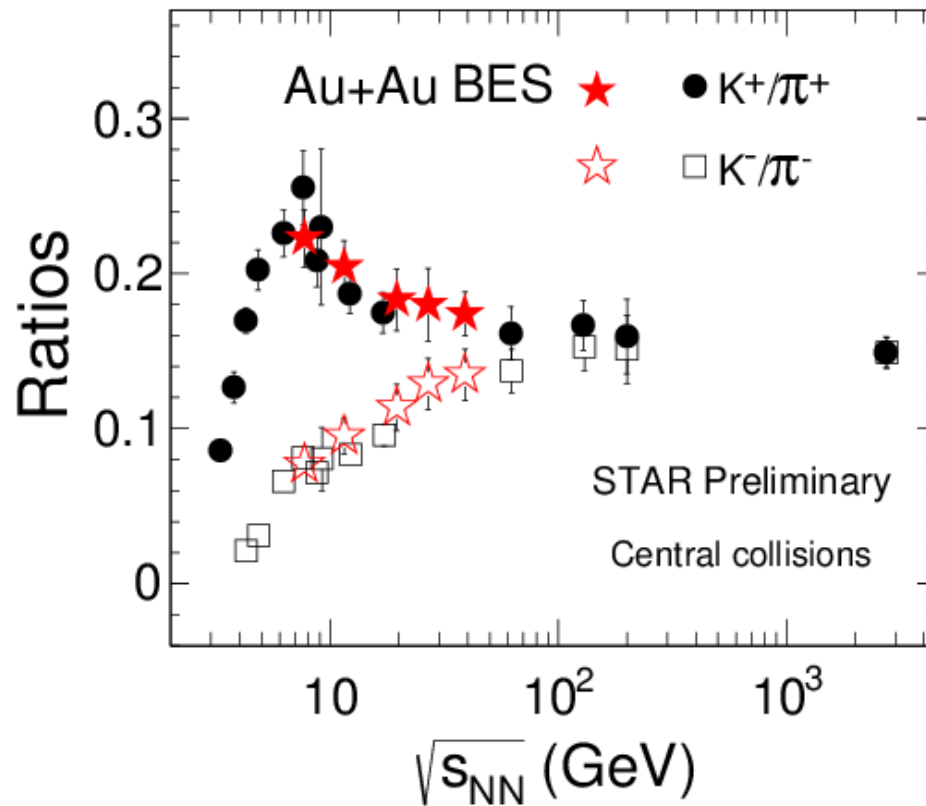
J. Phys. G: Nucl. Part. Phys.
27 (2001) 275

RQMD:

- K^+ N repulsive potential
- K^- N attractive potential
- Momentum dependent Skyrme forces
- Compression parameter
 - ✓ soft ~ 200 MeV
 - ✓ hard ~ 380 MeV

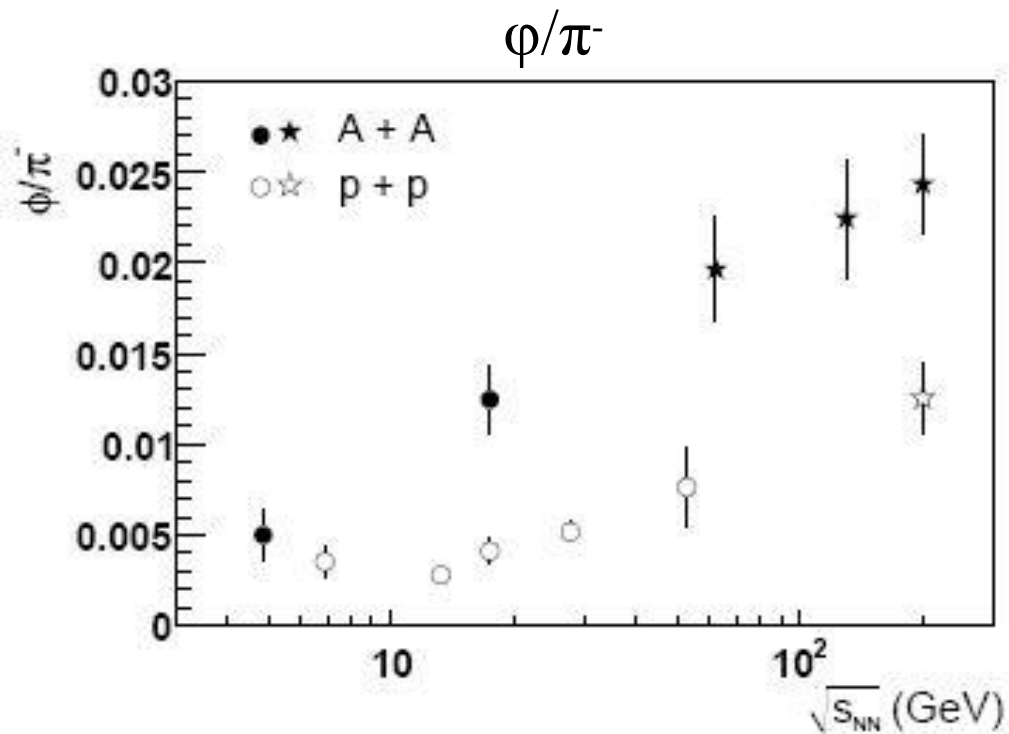
Enhancement of strangeness

- Clear evidence for “horn” structure in K^+/π^+ at ~ 30 A GeV !
- Non-horn structure in K^-/π^-



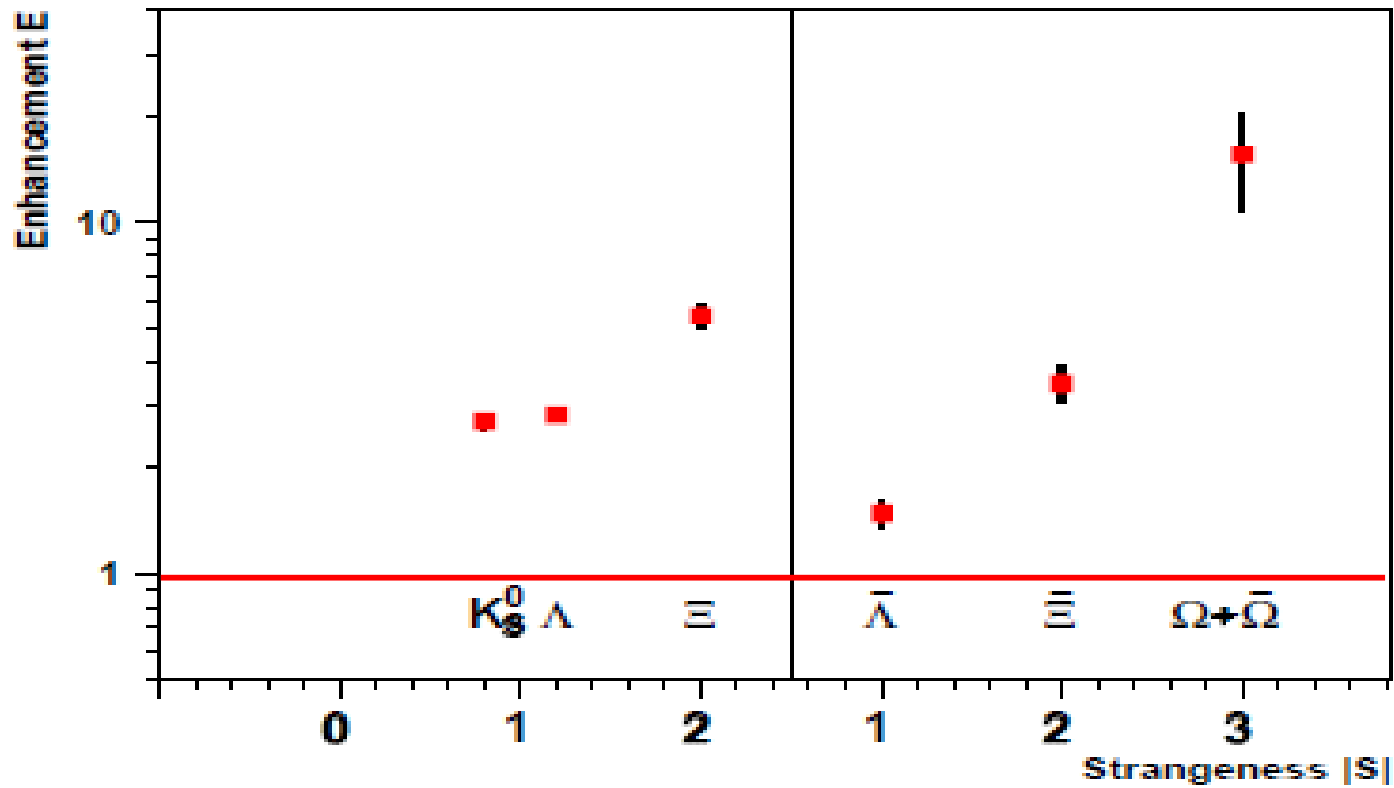
Enhancement of strangeness

ϕ -mesons



Enhanced yield of hyperons

PbPb vs pBe SPS



Thank you for your attention!