

**The possibility of studying polarization observables  
in reactions of the production of pions by polarized  
beams of protons and deuterons at the JINR LHEP**

*A.G. Litvinenko*

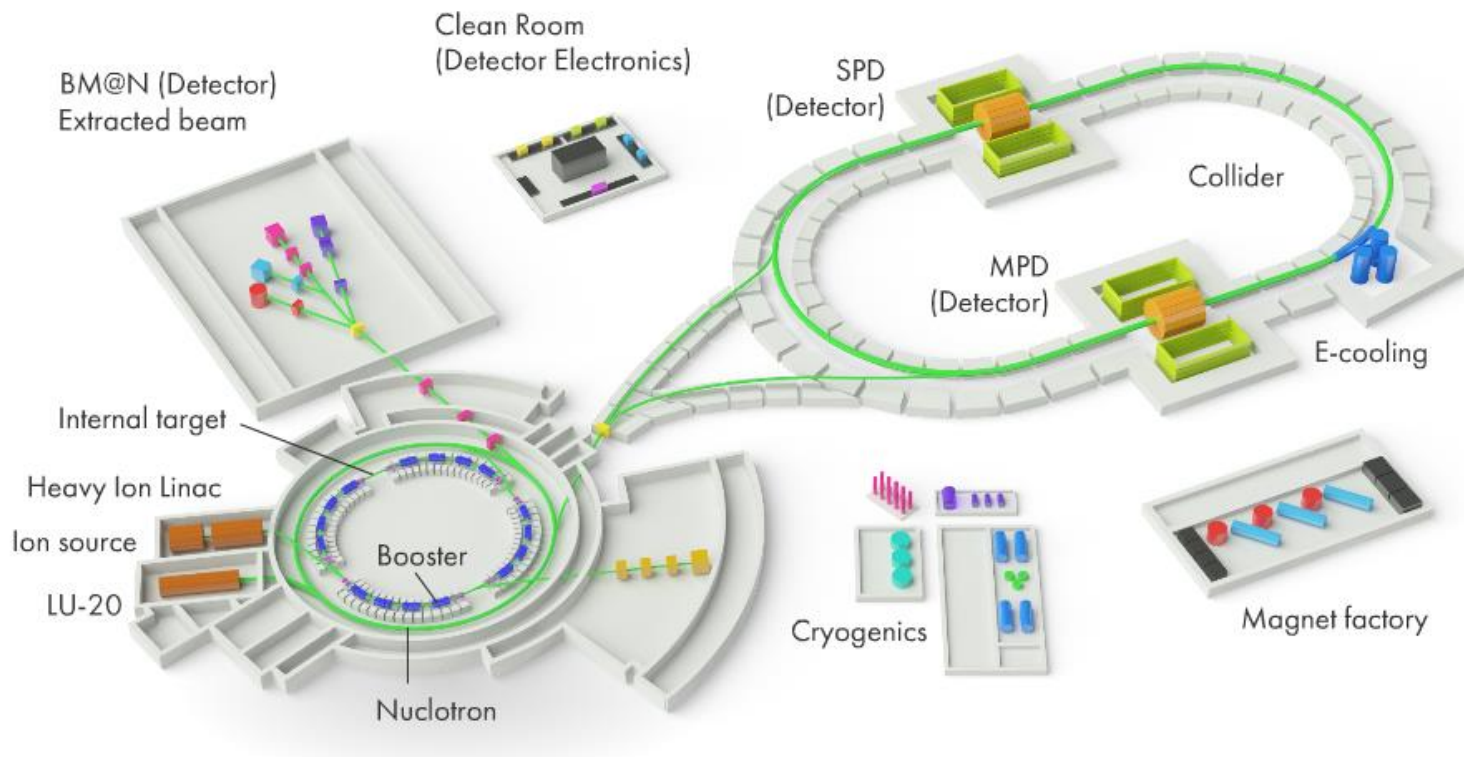
*([alitvin@jinr.ru](mailto:alitvin@jinr.ru))*

JINR, Dubna, Moscow region, Russia

University of Dubna, Dubna, Moscow. region, Russia

# NICA Complex

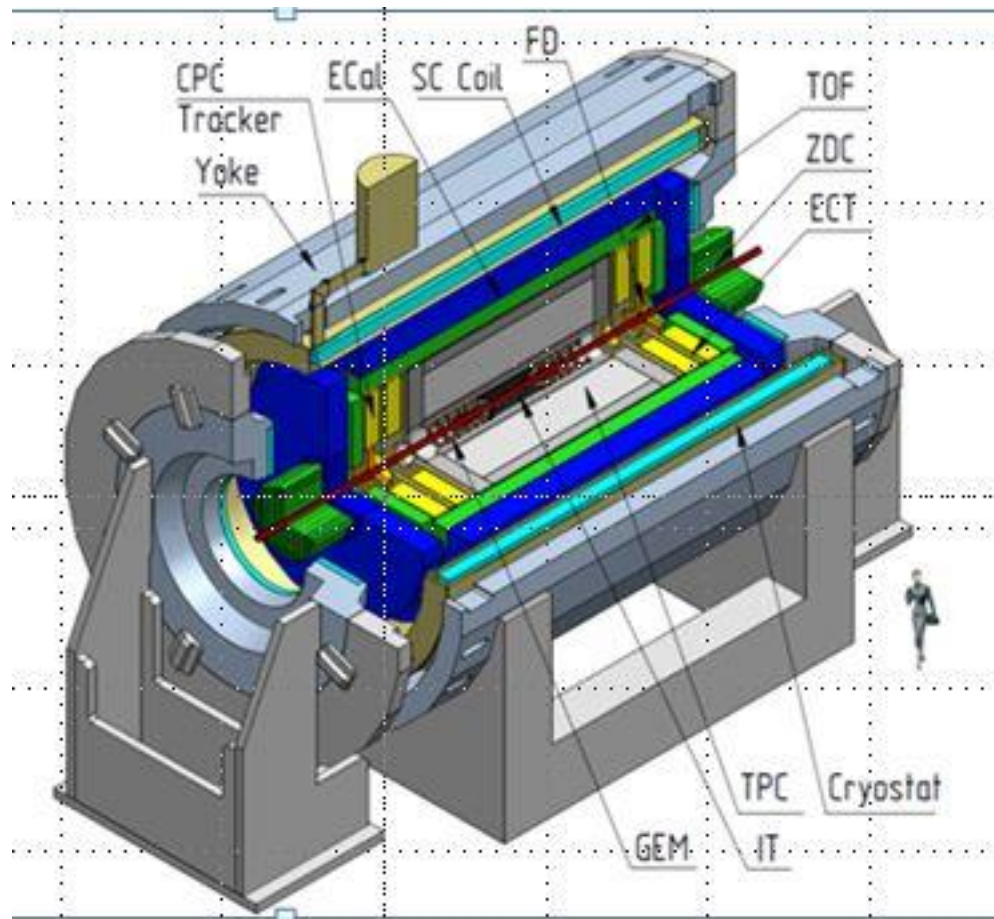
## *Nuclotron-based Ion Collider Facility (NICA)*



<http://nica.jinr.ru/complex.php>

# MPD

## Multi Purpose Detector (MPD)



<http://mpd.jinr.ru/mpd/>

## Motivation

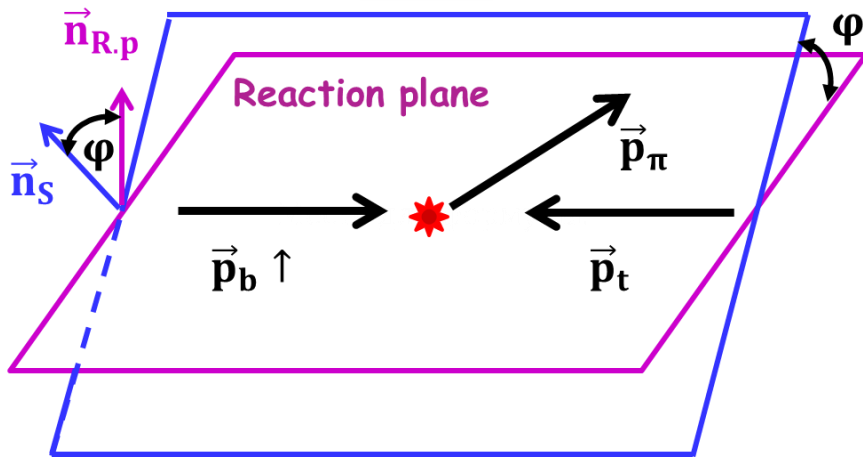
- ✓ Collisions of polarized beams are part of the NICA project
- ✓ Dubna has unique possibility to collide polarized proton and deuteron
- ✓ At the moment MPD@NICA is oriented to study heavy ion collision and has no proposal to use polarized beams
- ✓ One of the possibilities – to study transverse asymmetry of pion production

# Reaction



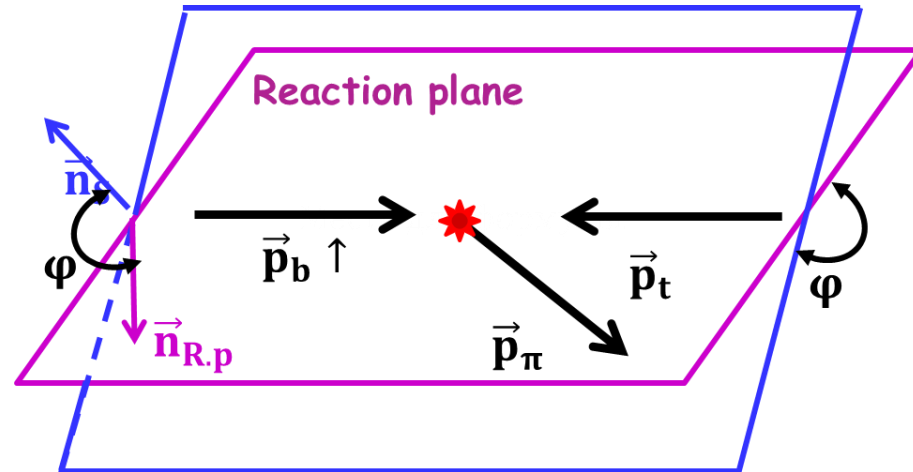
# Geometry

Left



$$(\vec{n}_s \vec{p}_b) = 0$$

Right



$$\vec{n}_{R,p} = [\vec{p}_b \vec{p}_\pi] / |[\vec{p}_b \vec{p}_\pi]|$$

*R. Machleidt, K. Holinde, and Elster. Phys. Rep., 149, 1, (1987).*

## Description and variables

$$E d\sigma/d\vec{p} = E d\sigma_0/d\vec{p} + (\vec{S}\vec{N})E d\sigma_S/d\vec{p}$$

$$E d\sigma/d\vec{p} = d\sigma_0(x_F, p_T, s) (1 + A_N(x_F, p_T, s) P \cos\varphi)$$

W.Haeberli. *Ann. Rev. Nucl. Sci.*, 17, 373, (1967)

$A_N$  - Single-Spin Asymmetry (analyzing power  $A_y$ )

Feynman variable  $x_F$

$$x_F = p_{c,l}/p_{c,max}$$

$p_{c,l}$  – momentum in c. m.

R.P.Feynman. *Phys. Rev.Lett.*, 23, 1311, (1969).

## Kinematic region

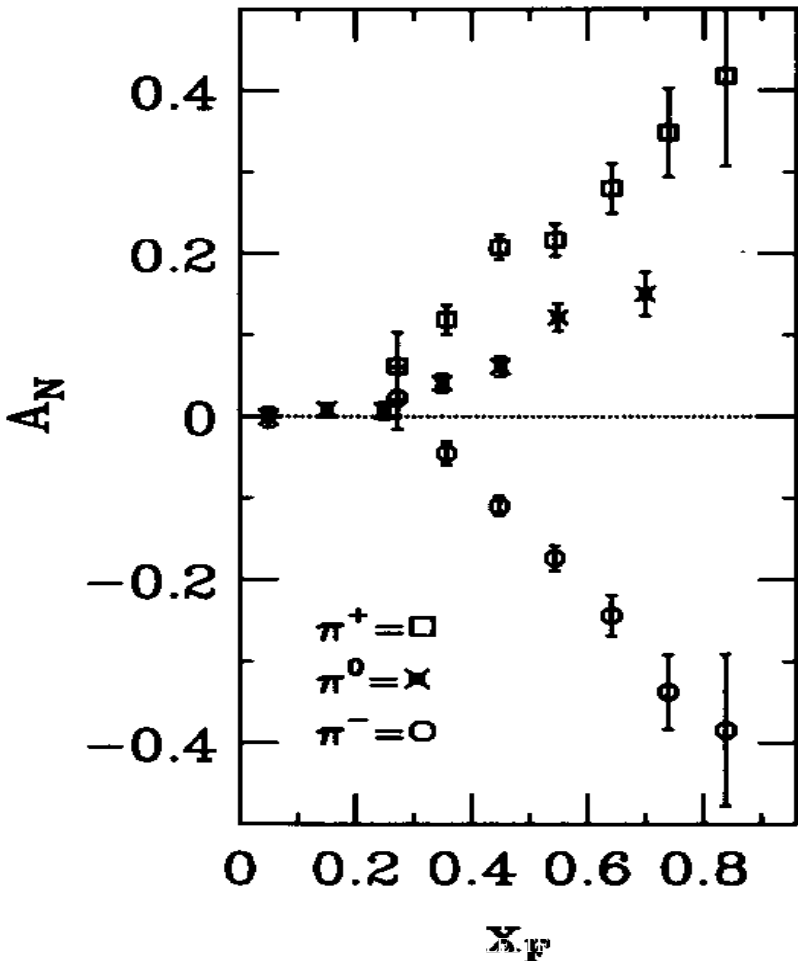
$$p \uparrow + p(A) = \pi + X$$

pion is detected in the fragmentation region of a polarized proton  $x_F > 0$

The experiments indicates a large asymmetry ( $A_N$  up to 40 %) in a wide range of collision energies

# High energy data

$(p \uparrow)_b(200 \text{ GeV}) + A_t = \pi + X; A_t = p; \sqrt{s} = 19.4 \text{ GeV}$



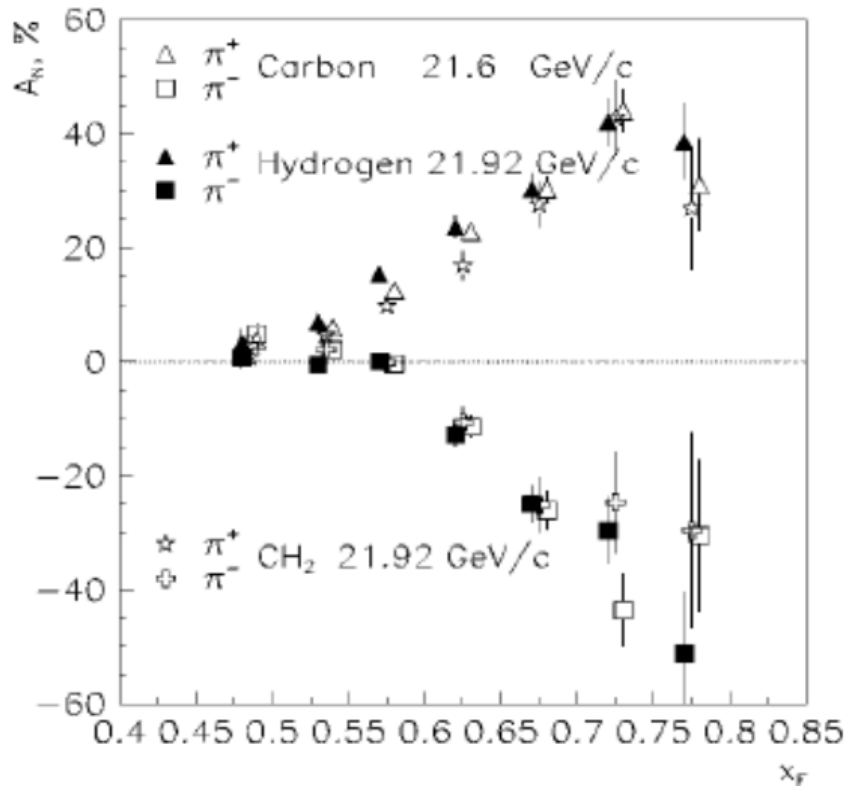
*FNAL E704, Phys. Lett. B 264, 462 (1991)*

$0.2 \text{ GeV}/c \leq p_T \leq 2 \text{ GeV}/c$



# Low energy data

$$(p \uparrow)_b(22 \text{ GeV}) + A_t = \pi + X; A_t = p; \sqrt{s} = 6.6 \text{ GeV}$$



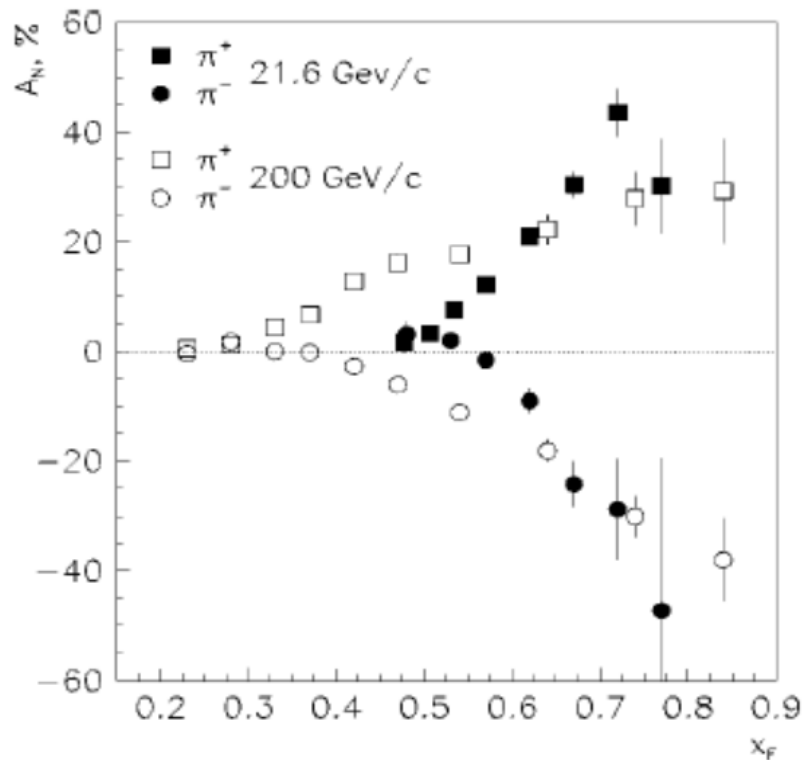
E925 Collaboration (AGS BNL)  
Phys.Rev. D 65 ,092008, (2002)

there is no dependence  
on the target for data in  
the fragmentation region  
of the polarized beam

$$0.2 \text{ GeV}/c \leq p_T \leq 2 \text{ GeV}/c$$

# High and Low energy data

$$p_b \uparrow + A_t = \pi + X$$



*D.L. Adams et al., Phys. Lett. B 264, 462 (1991) 200 GeV*

*C.E. Allgower et al., Phys. Rev. D 65, 092008, (2002) (22 GeV)*

$$0.2 \text{ GeV}/c \leq p_T \leq 2 \text{ GeV}/c$$

## Variables; distributions

- measurements with three polarization values  $P_+ > 0; P_- < 0; P_0 \equiv 0$
- number of recorded pions per polar angle interval  $dN_+/d\varphi; dN_-/d\varphi; dN_0/d\varphi;$
- corresponding luminosities  $L_+; L_-; L_0$
- normalized distributions along the polar angle  $n_i(\varphi) = (dN_i/d\varphi)/L_i$

$$n_i(\varphi) = \sigma_0(1 + P_i A_N \cos\varphi)$$

$$\langle n_i \rangle = \int_0^{2\pi} n_i(\varphi) d\varphi = \sum_{j=1}^{N_\varphi} n_i(\varphi_j) = 2\pi\sigma_{0,i}$$

$$\langle n_i \cos\varphi \rangle = \int_0^{2\pi} n_i(\varphi) \cos\varphi d\varphi = \sum_{j=1}^{N_\varphi} n_i(\varphi_j) \cos\varphi_j = \pi\sigma_{0,i} A_N P_i$$

## Asymmetry and experimental data

$$A_N = \frac{3}{P_+ P_-} \frac{P_- \langle n_+ \cos \varphi \rangle + P_+ \langle n_- \cos \varphi \rangle}{\langle n_- \rangle + \langle n_0 \rangle + \langle n_+ \rangle}$$

for simplicity

$$P_+ = -P_- = P$$

$$A_N = \frac{3}{P} \frac{\langle n_+ \cos \varphi \rangle - \langle n_- \cos \varphi \rangle}{\langle n_- \rangle + \langle n_0 \rangle + \langle n_+ \rangle}$$

## Errors estimations

$$n_i = (N_i)/(L_i)$$

$$N_0 \approx N_+ \approx N_-$$

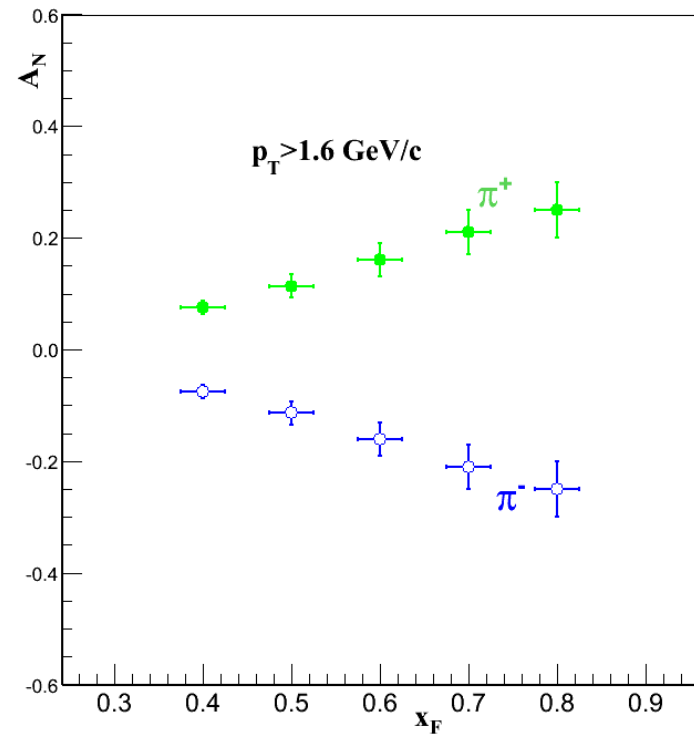
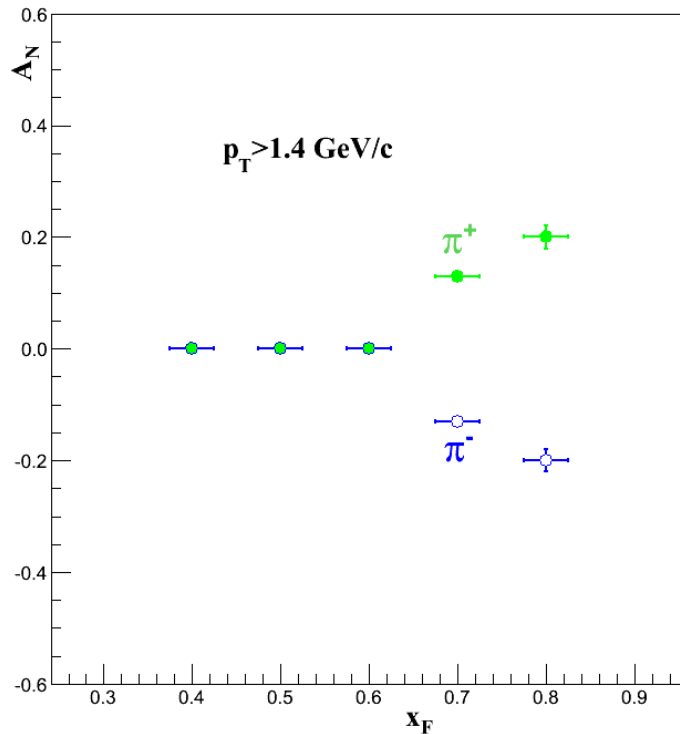
$$\Delta A_N \approx \frac{\sqrt{2}}{P} \sqrt{((\delta N_0)^2 + (\delta L_0)^2)}$$

$$\Delta A_N \approx A_N \delta P$$

# Asymmetry for $p \uparrow p \rightarrow \pi X$ at MPD

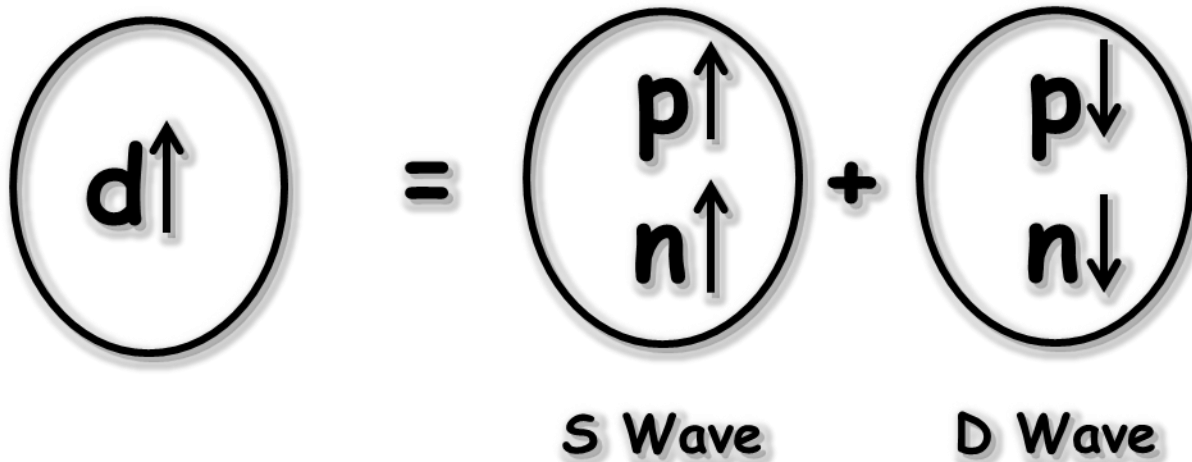
$L = 10^{31} \text{ cm}^{-2} \text{ s}^{-1}; T = 1 \text{ Month}$

$p \uparrow + p = \pi + X; \sqrt{s} = 4 \text{ GeV}, E_b = 7.6 \text{ GeV}$



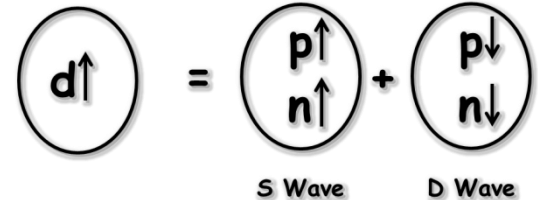
Unique possibility to obtain asymmetry  
of collision of polarized neutron  
from collision of polarised deuteron

$$d\uparrow + p(A) = \pi + X$$



# Structure of the cross section

$$\mathbf{D} \uparrow + \mathbf{p} = \pi + \mathbf{X}$$



cross section

$$d\sigma(d + p \rightarrow \pi X) = w_S d\sigma(d_S + p \rightarrow \pi X) + w_D d\sigma(d_D + p \rightarrow \pi X)$$

$$d\sigma(d_S \uparrow + p \rightarrow \pi X) = \alpha [d\sigma(p \uparrow + p \rightarrow \pi X) + d\sigma(n \uparrow + p \rightarrow \pi X)]$$

$$d\sigma(d_D \uparrow + p \rightarrow \pi X) = \alpha [d\sigma(p \downarrow + p \rightarrow \pi X) + d\sigma(n \downarrow + p \rightarrow \pi X)]$$

$\alpha \approx 0.85$  – shading factor



# Reid DWF

## Deuteron Wave Function (DWF)

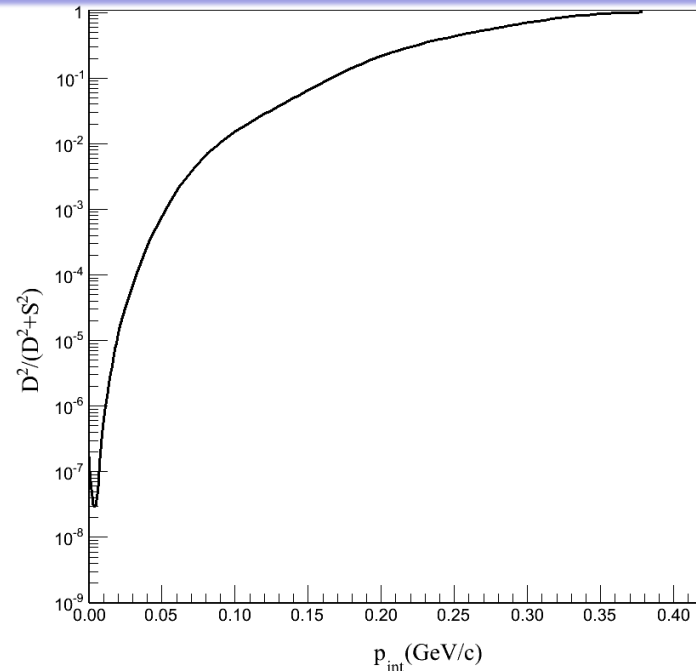
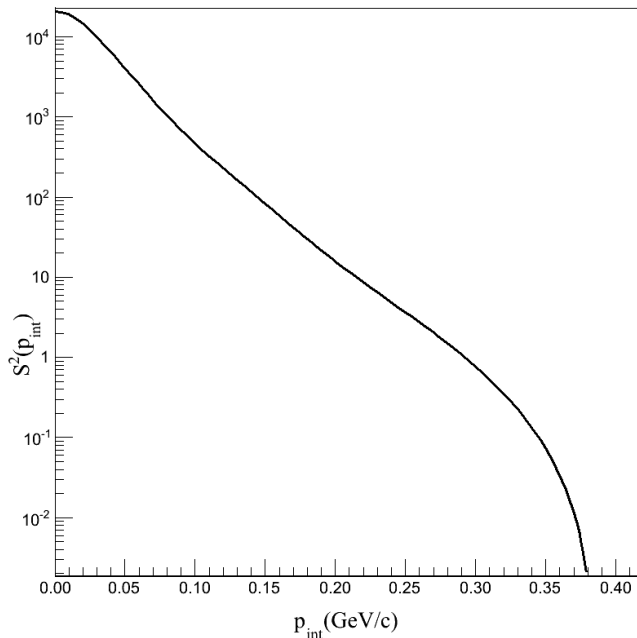
$$\sigma_p = \sqrt{\langle p^2 \rangle} = 0.09 \text{ GeV}/c$$



$$\Delta x_F \approx 0.004$$

$$w_S = \frac{\int S^2(p) p^2 dp d\Omega}{\int (S^2(p) + D^2(p)) p^2 dp d\Omega} = 0.953$$

$$w_D = \frac{\int D^2(p) p^2 dp d\Omega}{\int (S^2(p) + D^2(p)) p^2 dp d\Omega} = 0.047$$



Thus, we can ignore D-wave in cross section

$$(w_S > w_D) \rightarrow d\sigma(dp \rightarrow \pi X) = \alpha d\sigma(d_S p \rightarrow \pi X)$$

$$\begin{aligned} d\sigma_{+,d}(\varphi) &= d\sigma_{0,d}(1 + P_+ A_N(d) \cos\varphi) = \\ &+ \alpha d\sigma_{0,p}(1 + P_+ A_N(p) \cos\varphi) + \\ &+ \alpha d\sigma_{0,n}(1 + P_+ A_N(n) \cos\varphi) \end{aligned}$$

And we can obtain  $A_N(n)$  from  $A_N(d)$  and  $A_N(p)$

**Asymmetry  $A_N(\mathbf{n})$  for  $\mathbf{n} \uparrow + \mathbf{p} = \pi + \mathbf{X}$**

$$\sigma_0(\mathbf{d}/\mathbf{p}) = \frac{1}{3 \cdot 2\pi} [\langle \mathbf{n}_+(\mathbf{d}/\mathbf{p}) \rangle + \langle \mathbf{n}_0(\mathbf{d}/\mathbf{p}) \rangle + \langle \mathbf{n}_-(\mathbf{d}/\mathbf{p}) \rangle]$$

$$A_N(\mathbf{p}/\mathbf{d}) = \frac{3}{P_+ P_-} \frac{P_- \langle \mathbf{n}_+(\mathbf{p}/\mathbf{d}) \cos \varphi \rangle + P_+ \langle \mathbf{n}_-(\mathbf{p}/\mathbf{d}) \cos \varphi \rangle}{\langle \mathbf{n}_-(\mathbf{p}/\mathbf{d}) \rangle + \langle \mathbf{n}_0(\mathbf{p}/\mathbf{d}) \rangle + \langle \mathbf{n}_+(\mathbf{p}/\mathbf{d}) \rangle}$$

$$A_N(\mathbf{n}) = \frac{\sigma_0(\mathbf{d}) A_N(\mathbf{d}) - \alpha \sigma_0(\mathbf{p}) A_N(\mathbf{p})}{\sigma_0(\mathbf{d}) - \alpha \sigma_0(\mathbf{p})}$$

## Conclusion I

The report shows that the study of transverse asymmetry at MPD@NICA allows to obtain new data like  $A_N$  for  $n \uparrow + p(A) = \pi + X$ .

**But you need to know the absolute luminosity.**

Further studies can be connected with the central region and spin correlation measurements.

## Fixed target - luminosity

$$L(\text{cm}^{-2}\text{s}^{-1}) = N_{\text{bm}} \cdot N_A \frac{\rho_t(\text{g} \cdot \text{cm}^{-3})l_t(\text{cm})}{A_t} f(\text{s}^{-1})$$

$N_{\text{bm}}$  -the number of particles in the burst

$N_A = 6 \cdot 10^{23}$  - Avogadro number

$\rho_t$  -the density of the target

$l_t$  -target length along the beam

$f$  - repetition rate

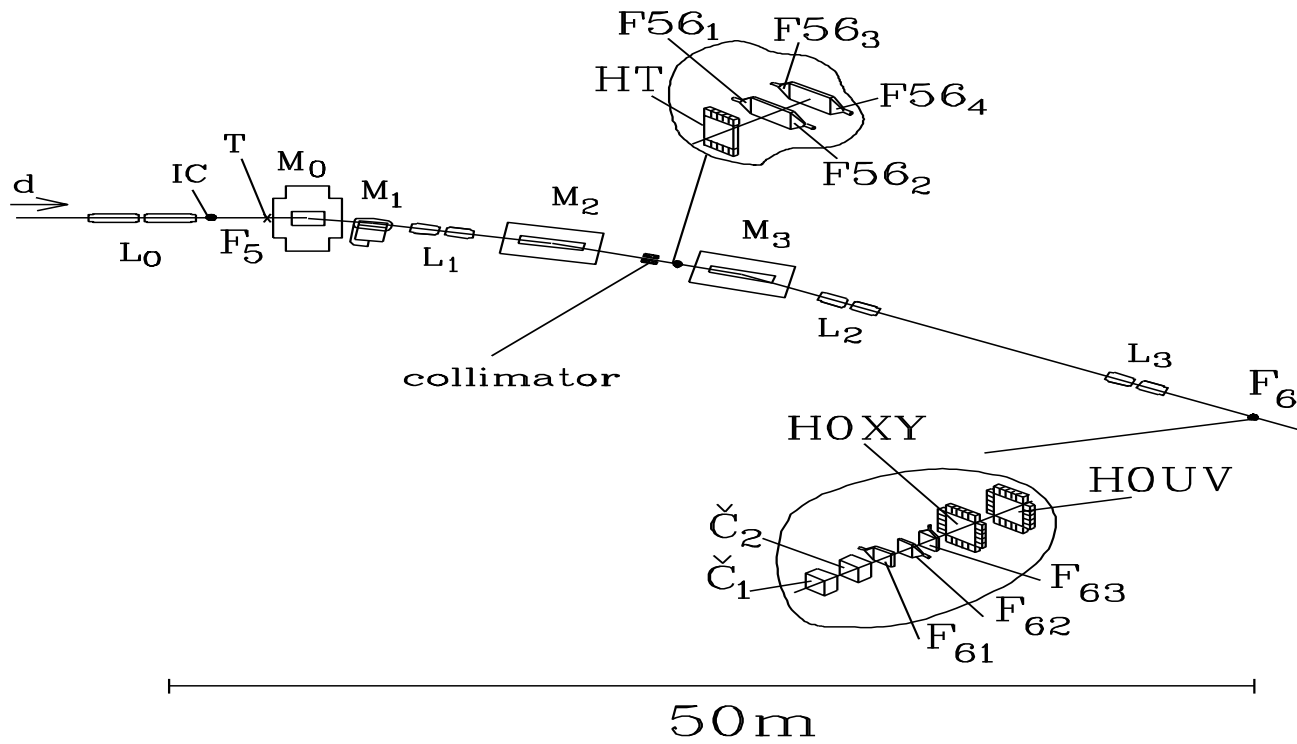
### estimations

✓ carbon target

$$\{N_{\text{bm}} = 10^{10}; l_t = 10 \text{ cm}, f = 0.1 \text{ s}^{-1}\} \rightarrow L \approx 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$$

✓ beryllium target

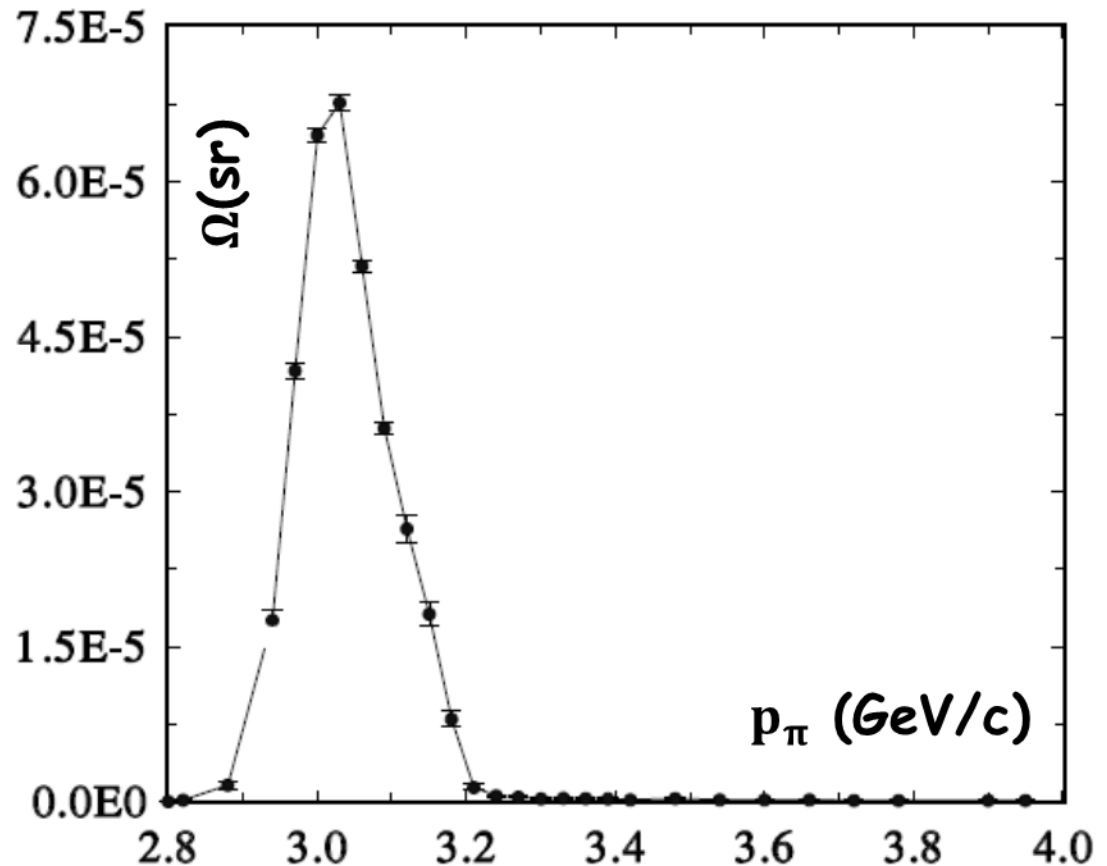
$$\{N_{\text{bm}} = 10^{10}; l_t = 10 \text{ cm}, f = 0.1 \text{ s}^{-1}\} \rightarrow L \approx 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$$



$F_i, F_{i,j}$   
 $\check{C}_1, \check{C}_2$   
 HT, HOXY, HOUV  
 $M_i$   
 $L_i$   
 IC  
 T

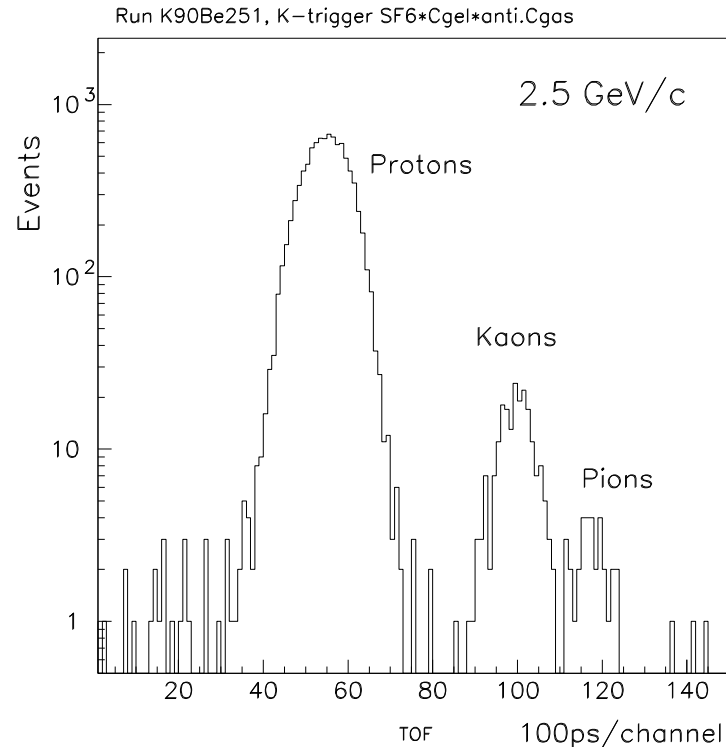
the scintillation counters  
 the threshold Cherenkov counters  
 the scintillation hodoscope  
 the bending magnets  
 a doublet of quadrupole lenses  
 ionization chamber  
 target

# Acceptance



$$\text{Acc}(4.5 \text{ GeV/c}) = \int \Omega(p) dp = 1.8 \cdot 10^{-5} \text{ sr} \cdot \text{GeV/c}; \Delta p/p = 2 \%$$

# Time-of-Flight spectrum



protons are suppressed by using aerogel Cherenkov counter ( $n=1.035$ ) and pions are suppressed by gas Cherenkov counter (anticoincidences,  $\text{CO}_2$ ,  $P = 12$  at.)



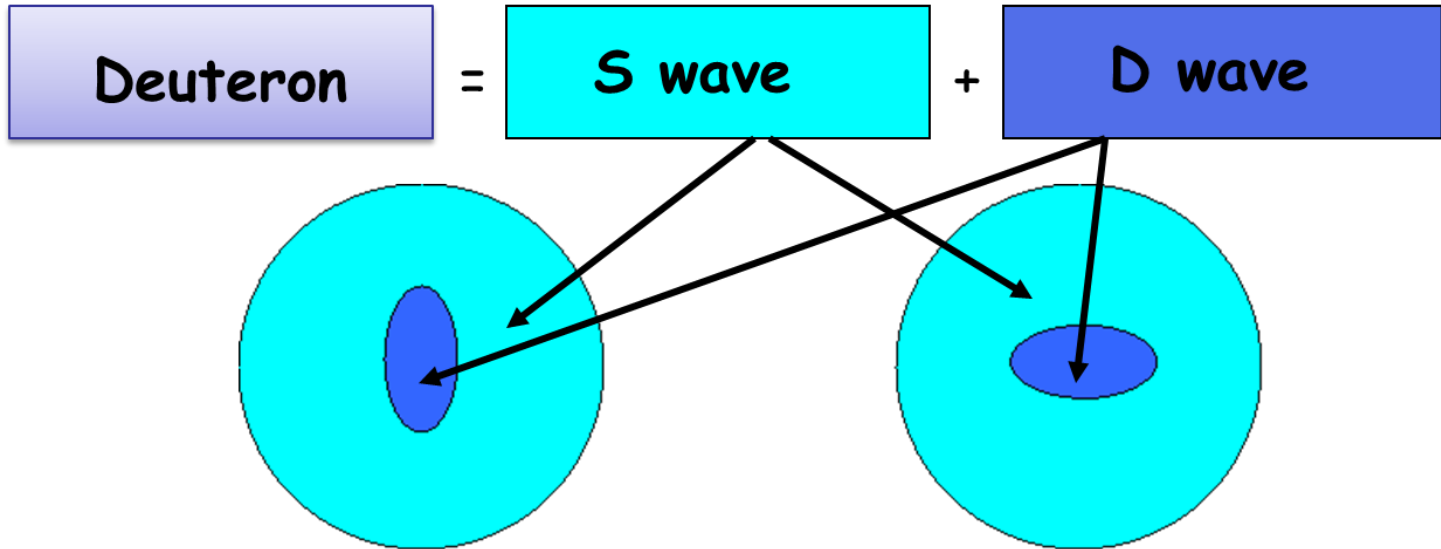
# Deuteron Tensor polarization

$$N_{+1} + N_0 + N_{-1} = 1$$

$$\tau_{20} = \frac{1}{\sqrt{2}}(1 - 3N_0) \rightarrow -\sqrt{2} \leq \tau_{20} \leq \frac{1}{\sqrt{2}}$$

$$p_{yy} = (1 - 3N_0) \rightarrow -2 \leq \tau_{20} \leq 1$$

$$\tau_{20} = \frac{1}{\sqrt{2}} p_{yy}$$



# Deuteron fragmentation into cumulative pions

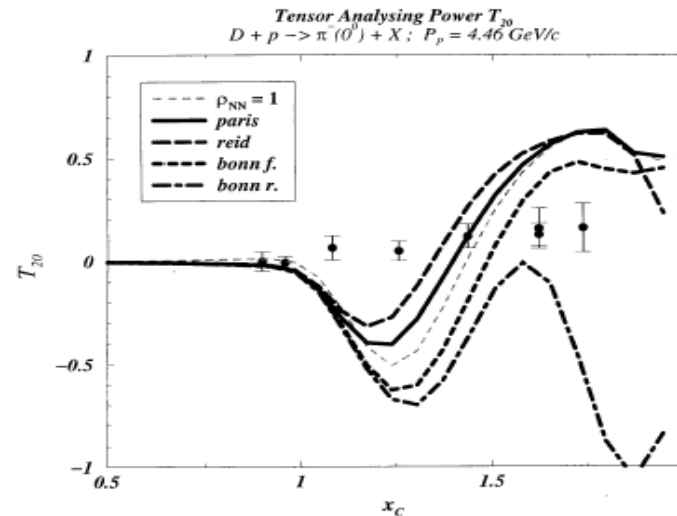
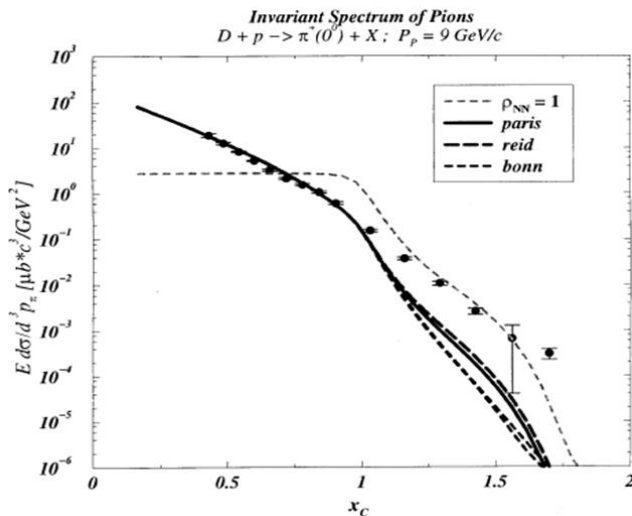
$$\mathbf{D} \uparrow + \mathbf{p}(\mathbf{A}) = \boldsymbol{\pi}(\boldsymbol{\theta}) + \mathbf{X}$$

$$d\sigma = d\sigma_0 \left( 1 + \frac{3}{2} A_y p_y + \frac{1}{2} A_{yy} p_{yy} \right)$$

$$T_{20} = -\frac{1}{\sqrt{2}} A_{yy}$$

# Deuteron fragmentation into cumulative particle Polarization observables

$$\vec{D} + p = \pi^-(0^0) + X$$

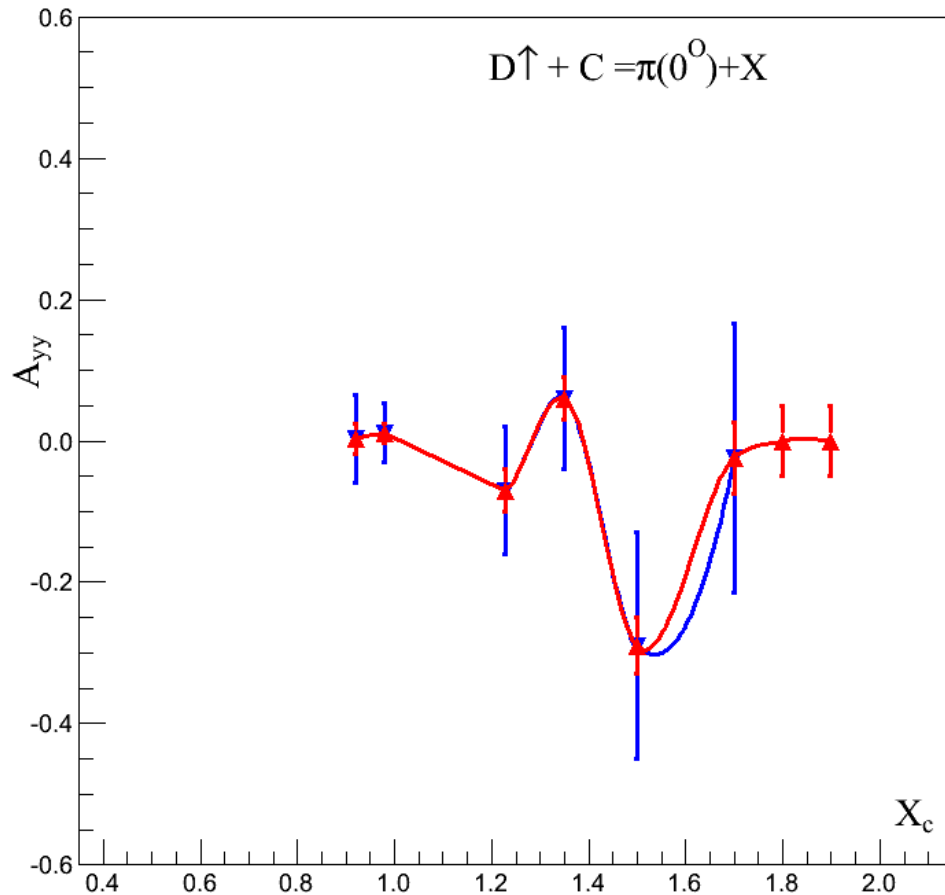


In Impulse Approximation  $d\sigma$  depends strongly on  $\rho_{NN} = d\sigma(NN \rightarrow c)$

In Impulse Approximation  $T_{20}$  depends weakly ( $\Delta T_{20} < 0.2$ ) on  $\rho_{NN} = d\sigma(NN \rightarrow c)$

# production of $\pi$ under a zero angle

$$X_c = \frac{m_N E_\pi - m_\pi^2 / 2}{p_\pi p_b \cos\theta_\pi - E_\pi E_b + E_b m_N - (m_N)^2}$$



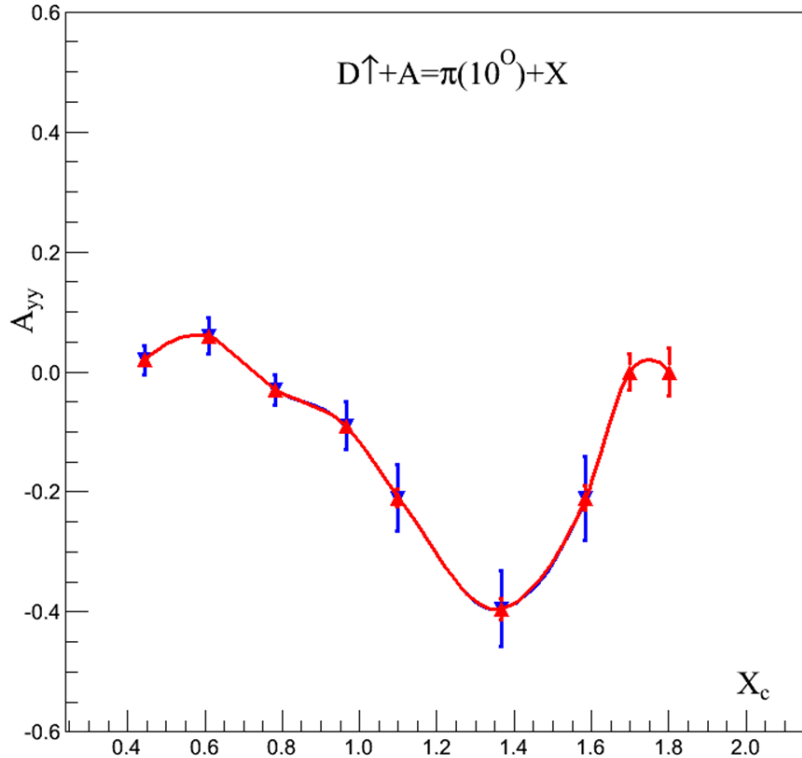
blue - exp. data;  
one week;  $I = 10^9$  per burst

*S.V.Afanasiev et al.,  
Nucl.Phys. A(625),p.817,(1997)*

red -red available in two weeks  
for  $I = 10^{10}$  per burst

# production of $\pi$ under a zero angle

$$X_c = \frac{m_N E_\pi - m_\pi^2 / 2}{p_\pi p_b \cos\theta_\pi - E_\pi E_b + E_b m_N - (m_N)^2}$$



blue - exp. data;  
one week;  $I = 10^9$  per burst

*S.V.Afanasiev et al., Phys.Lett.B,  
p.445,(1998)*

red - available in two weeks for  $I$   
 $= 10^{10}$  per burst

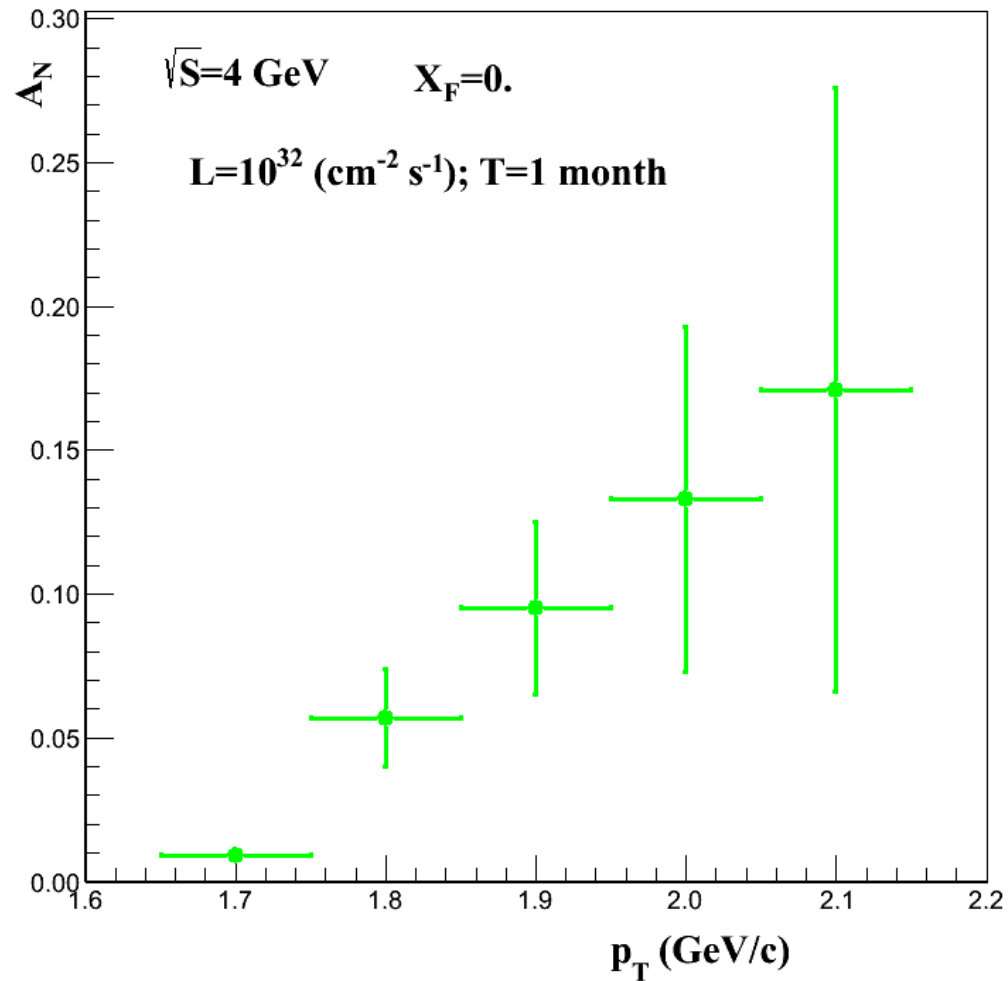
## Conclusion II

The study of tensor analyzing power with the beam intensity  $10^{10}$  would allow us:

1. to obtain more accurate data (in factor 3)
2. to reach to the region of higher internal momentum  $p_{int}$  up to  $0.75 \text{ GeV}/c$  ( $l_{NN} \cong 0.3 \text{ fm}$ )
3. to include into study new particles (positive and negative kaons)

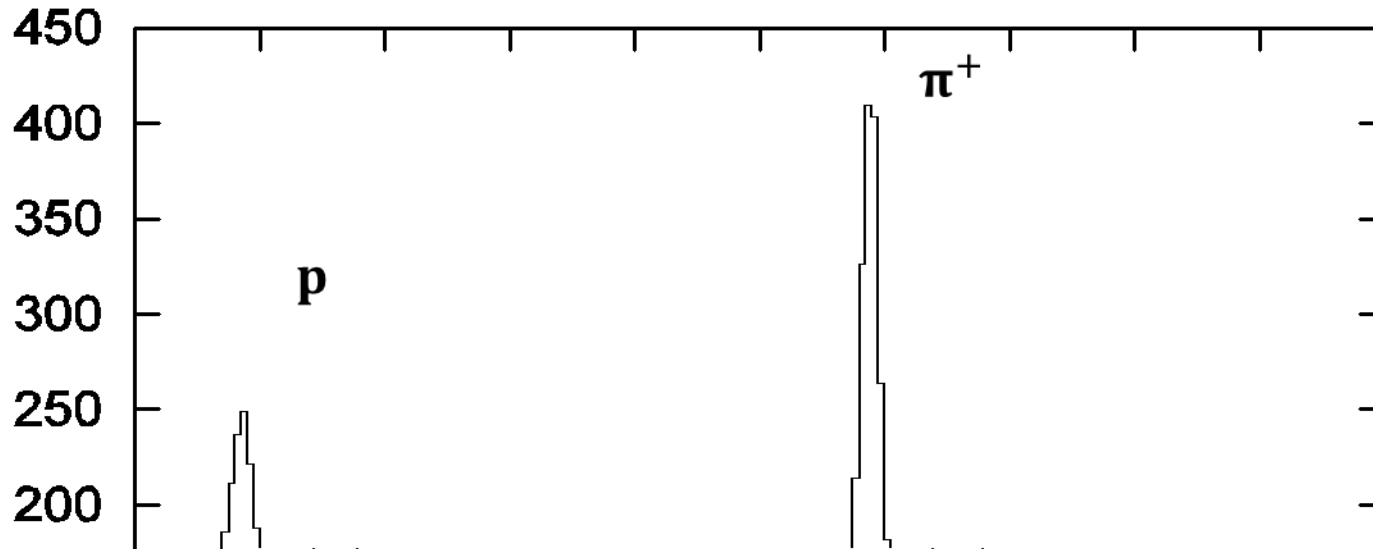
**Backup slides**

# Asymmetry in the central region





# Time-of-Flight spectrum

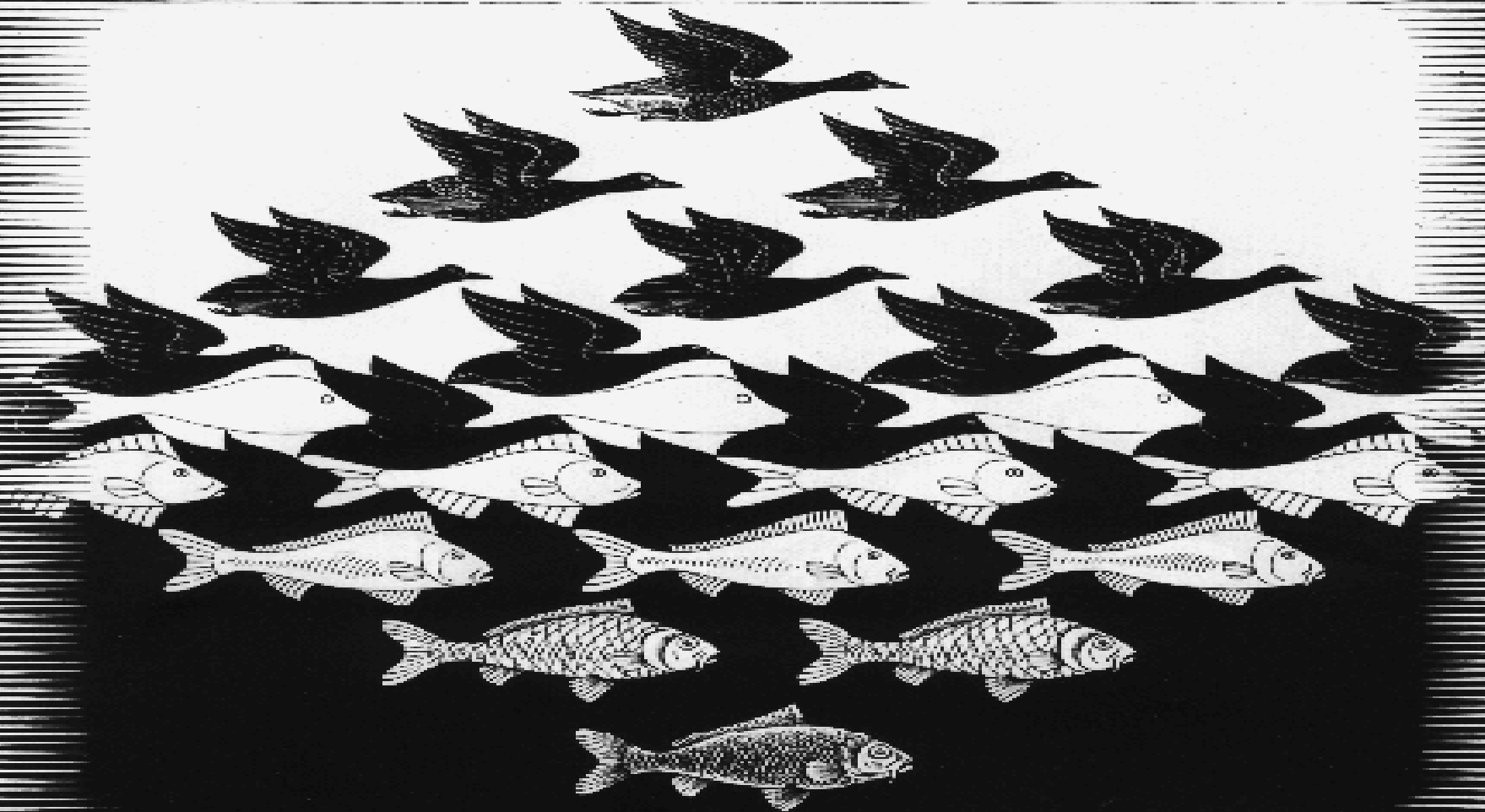


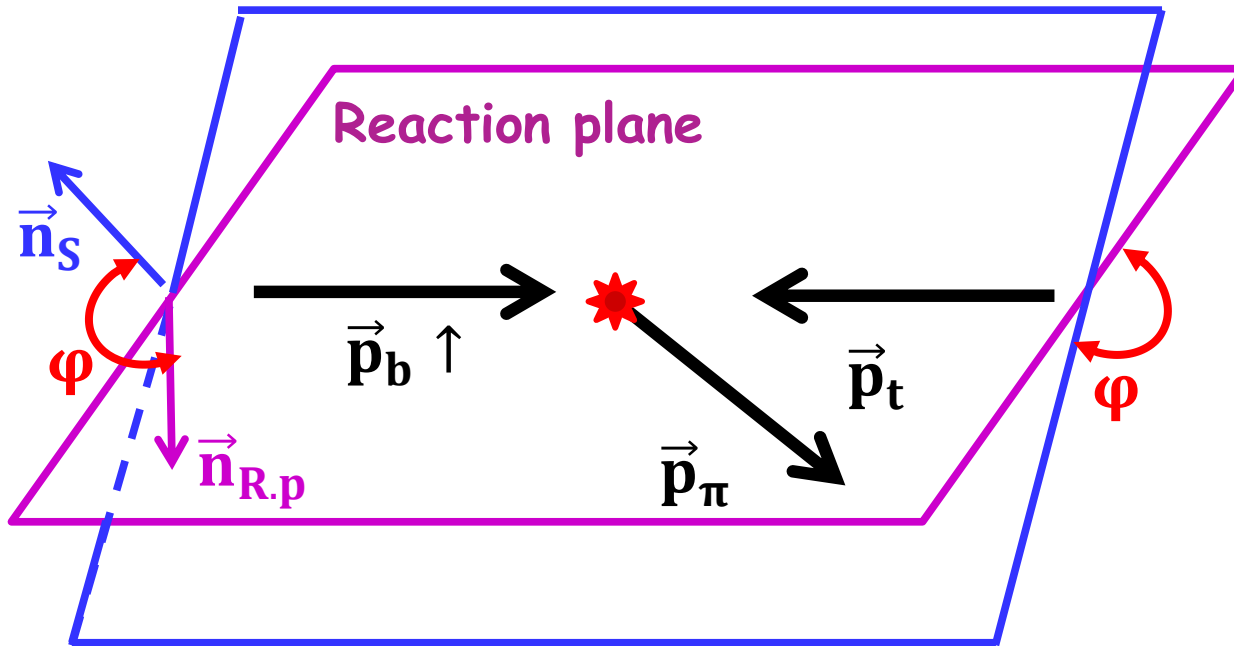
ns/ch)

) 780 800

protons are suppressed by 100 times  
using aerogel Cherenkov counter

# DUALITY – HOW IT IS LOOKS LIKE





$$(\vec{n}_s \vec{p}_b) = 0$$

$$\vec{n}_{R.p} = [\vec{p}_b \vec{p}_\pi] / |[\vec{p}_b \vec{p}_\pi]|$$

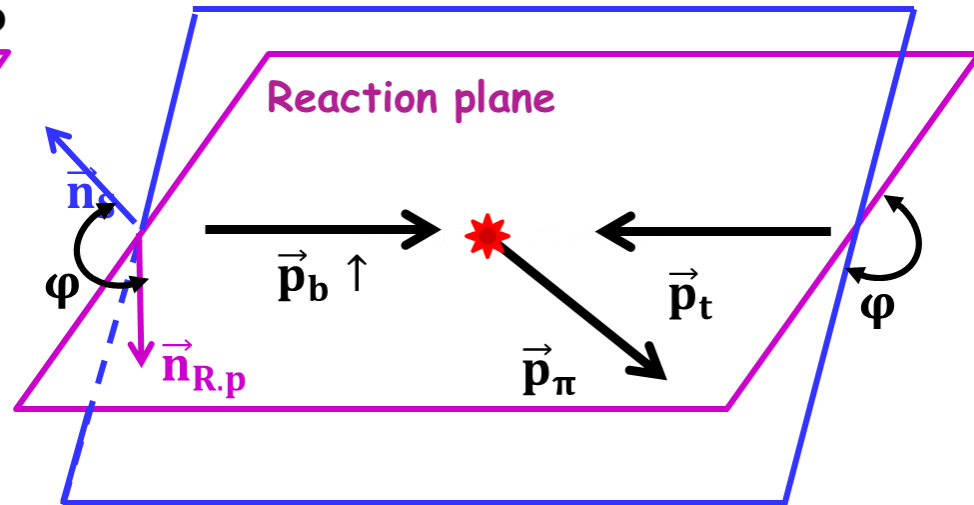
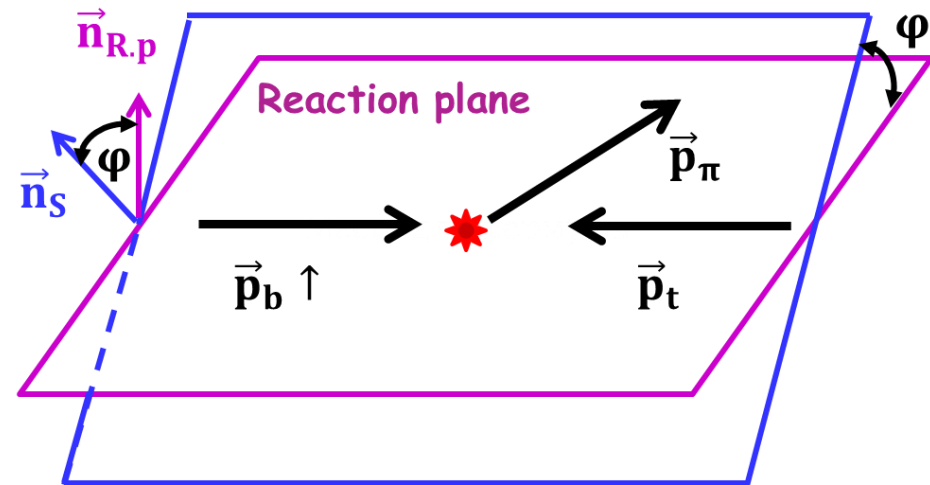
# Reaction



# Geometry

Left

Right



$$(\vec{n}_s \vec{p}_b) = 0$$

$$\vec{n}_{\text{R.p}} = [\vec{p}_b \vec{p}_\pi] / \left| [\vec{p}_b \vec{p}_\pi] \right|$$

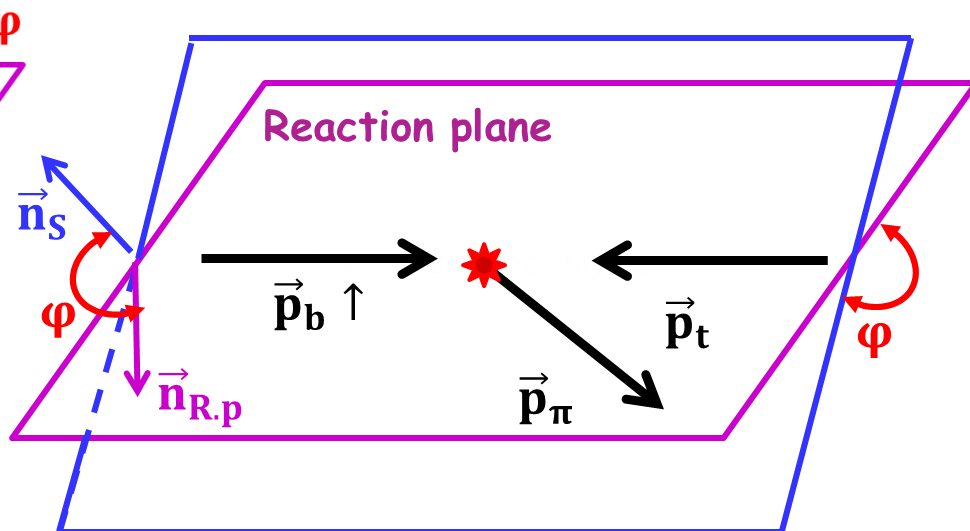
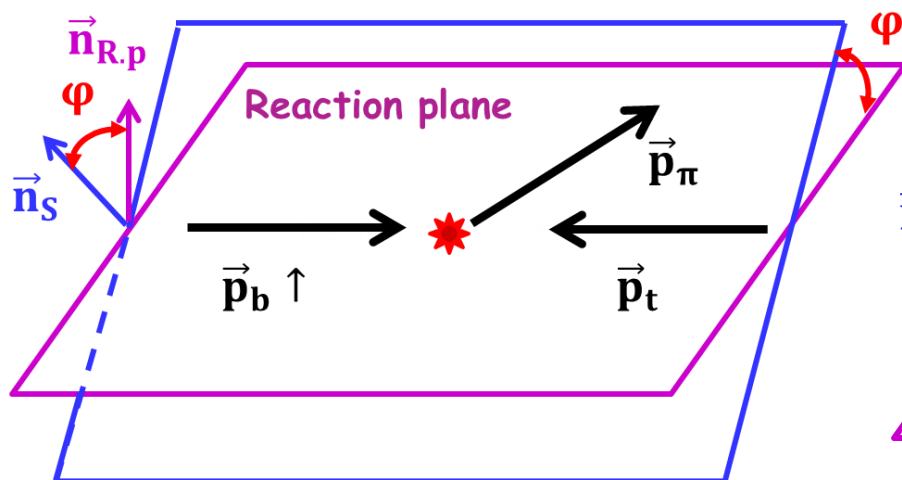
# Reaction



# Geometry

Left

Right



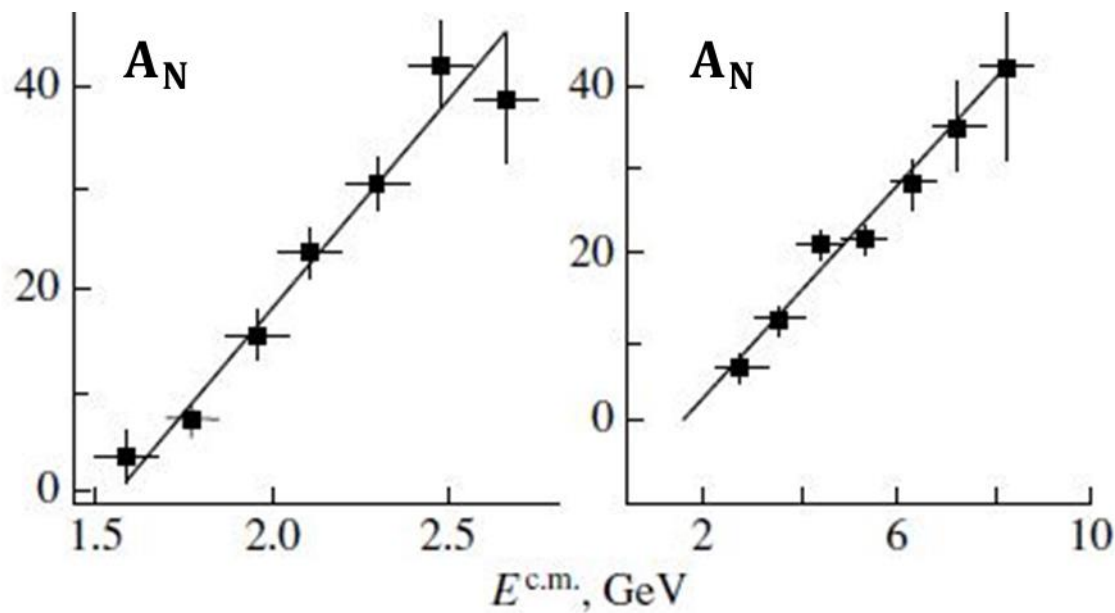
$$(\vec{n}_s \vec{p}_b) = 0$$

$$\vec{n}_{\text{R.p}} = [\vec{p}_b \vec{p}_\pi] / \left| [\vec{p}_b \vec{p}_\pi] \right|$$

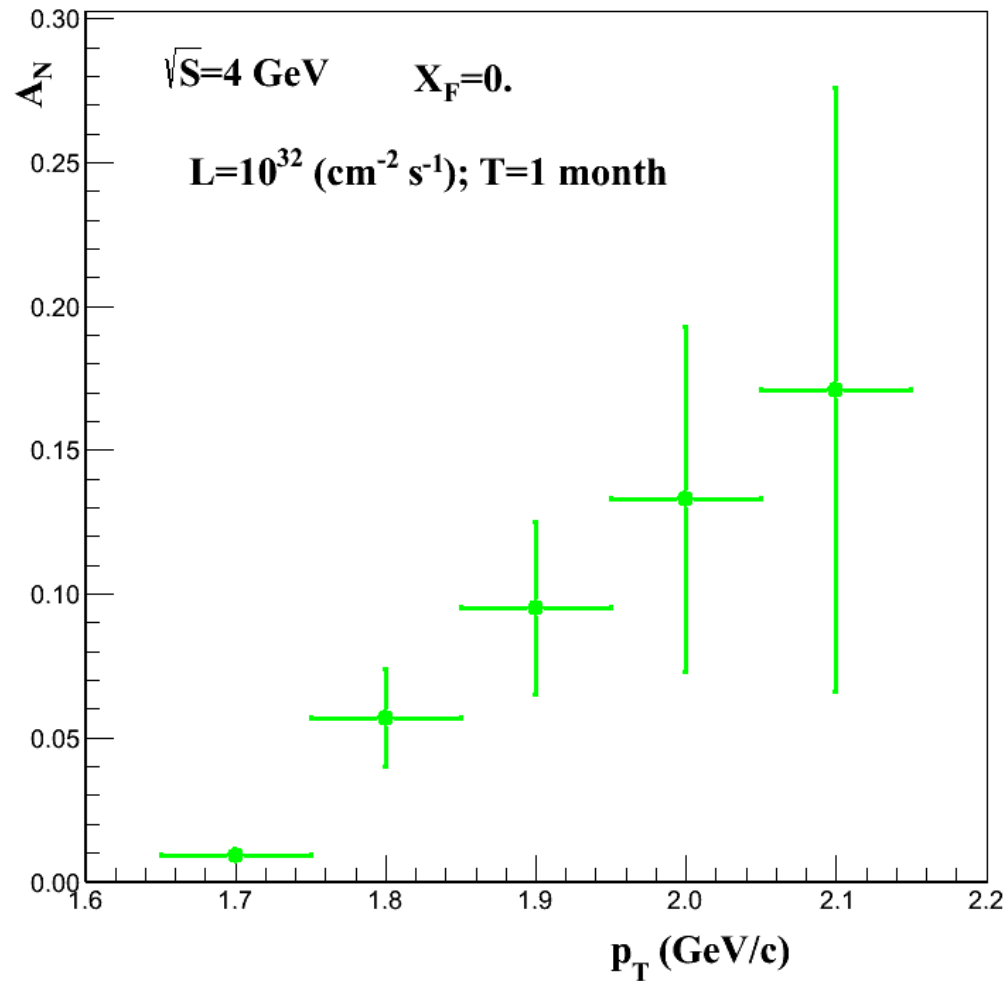
# Approximation from energy at the center of mass

A.N. Vasiliev and V. V. Mochalov,  
*Physics of Atomic Nuclei, Vol. 67, p. 2169 (2004)*

$$A_N = \begin{cases} = 0, E_{c.m.} < E_0; \\ = k(E_{c.m.} - E_0) \end{cases}$$



# Asymmetry in the central region



## Asymmetry and experimental data

$$A_N = \frac{3}{P_+ P_-} \frac{\int_0^{2\pi} \left( (P_- n_+(\varphi) + P_+ n_-(\varphi)) \cos\varphi \right) d\varphi}{\int_0^{2\pi} (n_+(\varphi) + n_0(\varphi) + n_-(\varphi)) d\varphi}$$

for simplicity and clarity

$$P_+ = -P_- = P$$

$$A_N = \frac{3}{P} \frac{\int_0^{2\pi} \left( (n_+(\varphi) - n_-(\varphi)) \cos\varphi \right) d\varphi}{\int_0^{2\pi} (n_+(\varphi) + n_0(\varphi) + n_-(\varphi)) d\varphi}$$

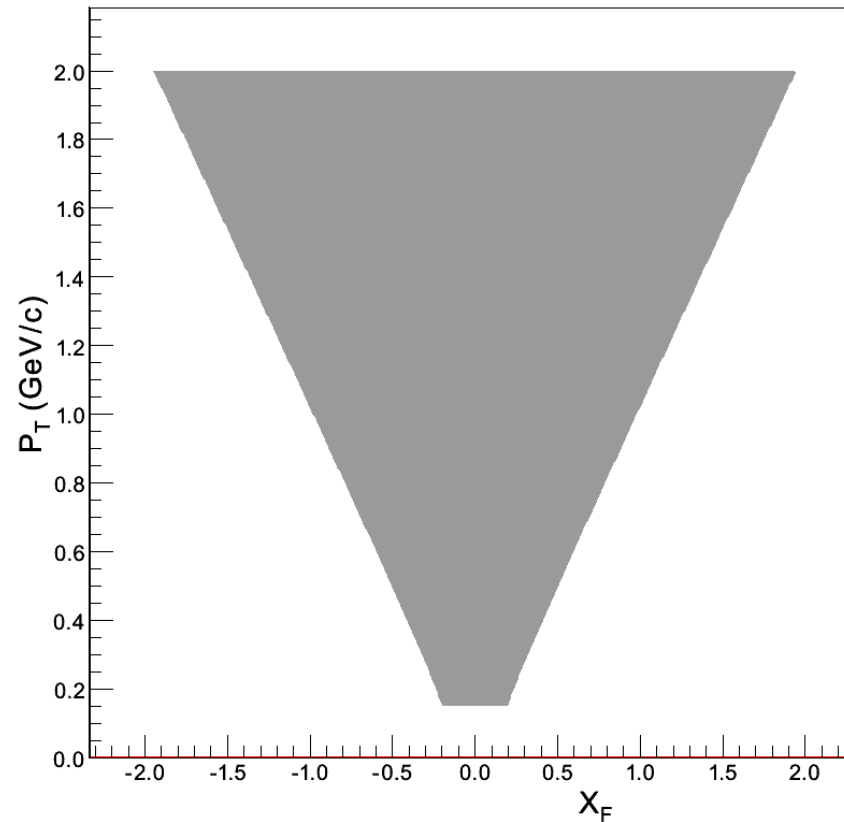


# MPD pion acceptance ( $x_F, p_T$ )

*Barrel part*



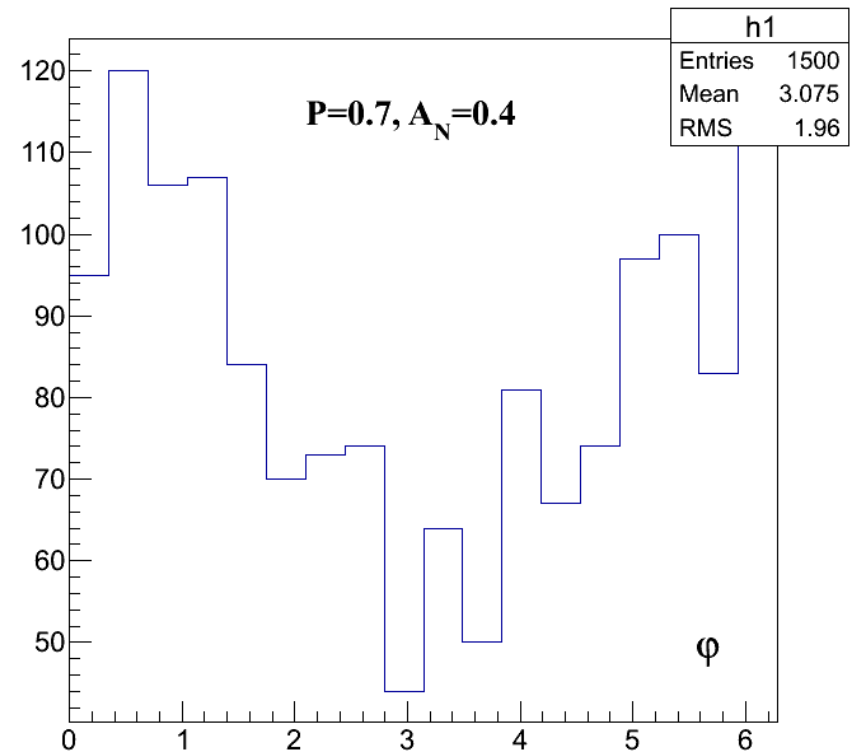
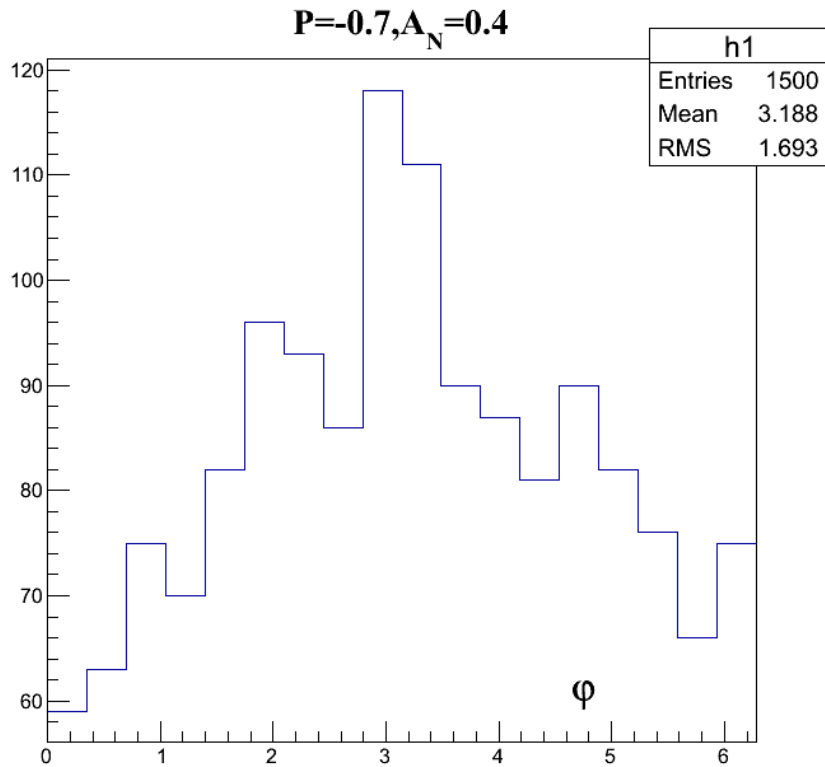
$\sqrt{s_{NN}} = 4 \text{ GeV}; P_b = 7.5 \text{ GeV}/c$



$$p_{\pi} (\text{max}) \approx \sqrt{s}/2$$

$$x_F \approx (2p_L)/\sqrt{s}$$

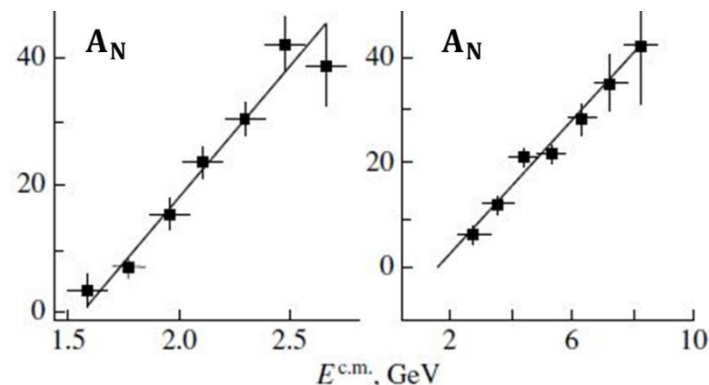
# Polar angle distribution



# Approximation from energy at the center of mass

A.N. Vasiliev and V. V. Mochalov,  
*Physics of Atomic Nuclei, Vol. 67, p. 2169 (2004)*

$$A_N = \begin{cases} = 0, E_{\text{cm}} < E_0 (\approx 1.6 \text{ GeV}); \\ = k(E_{\text{cm}} - E_0 (\approx 1.6 \text{ GeV})) \end{cases}$$



**NICA energy (<http://nica.jinr.ru>):**

- ❖ Collider mode (ions  $A > 1$ )  $\sqrt{S_{\text{NN}}} \leq 11 \text{ GeV}$   $\longleftrightarrow$  Fix.targ.  $E_b \leq 5.5 \text{ GeV}/N$
- ❖ Collider mode (protons)  $\sqrt{S_{\text{NN}}} \leq 22 \text{ GeV}$   $\longleftrightarrow$  Fix.targ.  $E_b \leq 11 \text{ GeV}$

**Fix.targ**

$$E_b \leq 5.5 \text{ GeV}/N \longleftrightarrow \sqrt{S_{\text{NN}}} \leq 3.5 \text{ GeV} \longrightarrow E_{\text{cm},\pi} \leq 1.25 \text{ GeV}; A_N = 0$$

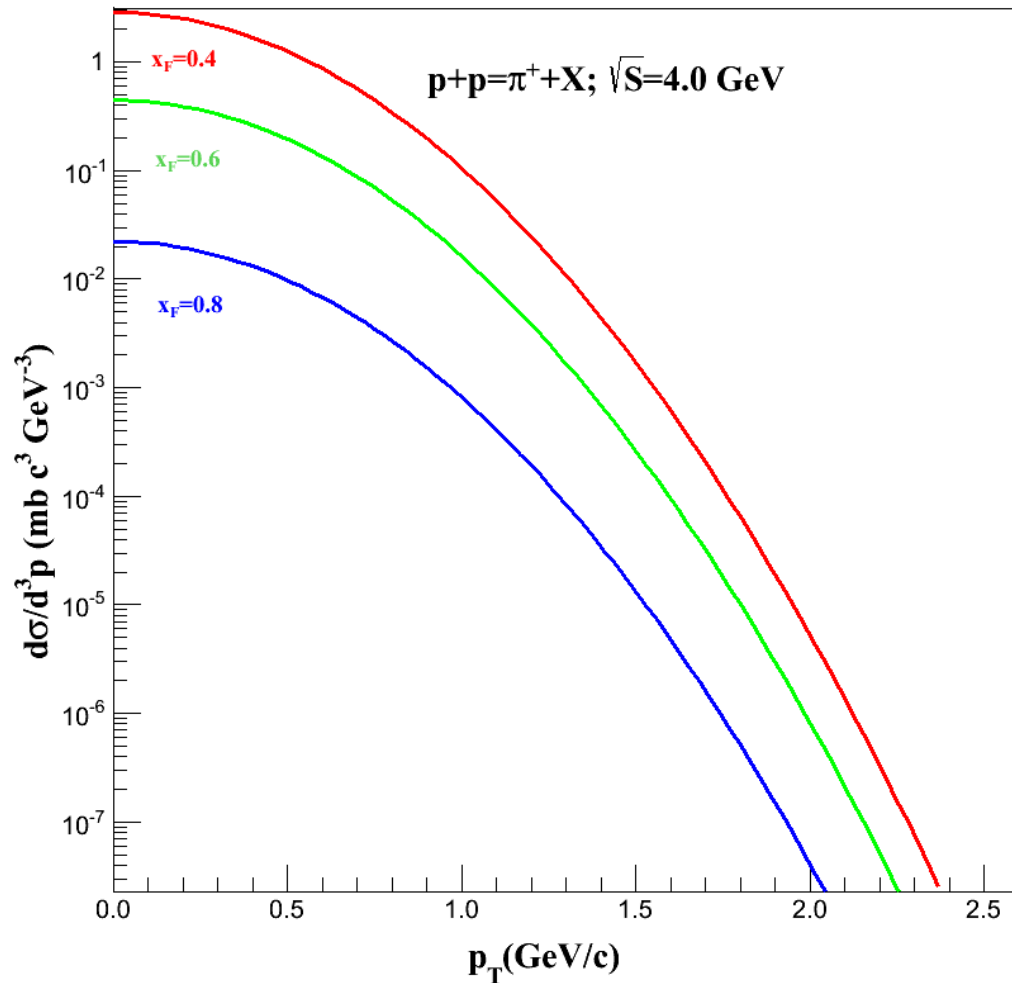
$$E_b \leq 11 \text{ GeV} \longleftrightarrow \sqrt{S_{\text{NN}}} \leq 4.7 \text{ GeV} \longrightarrow E_{\text{cm},\pi} \leq 2.0 \text{ GeV}; A_N \leq 0.13$$

## Conclusion II

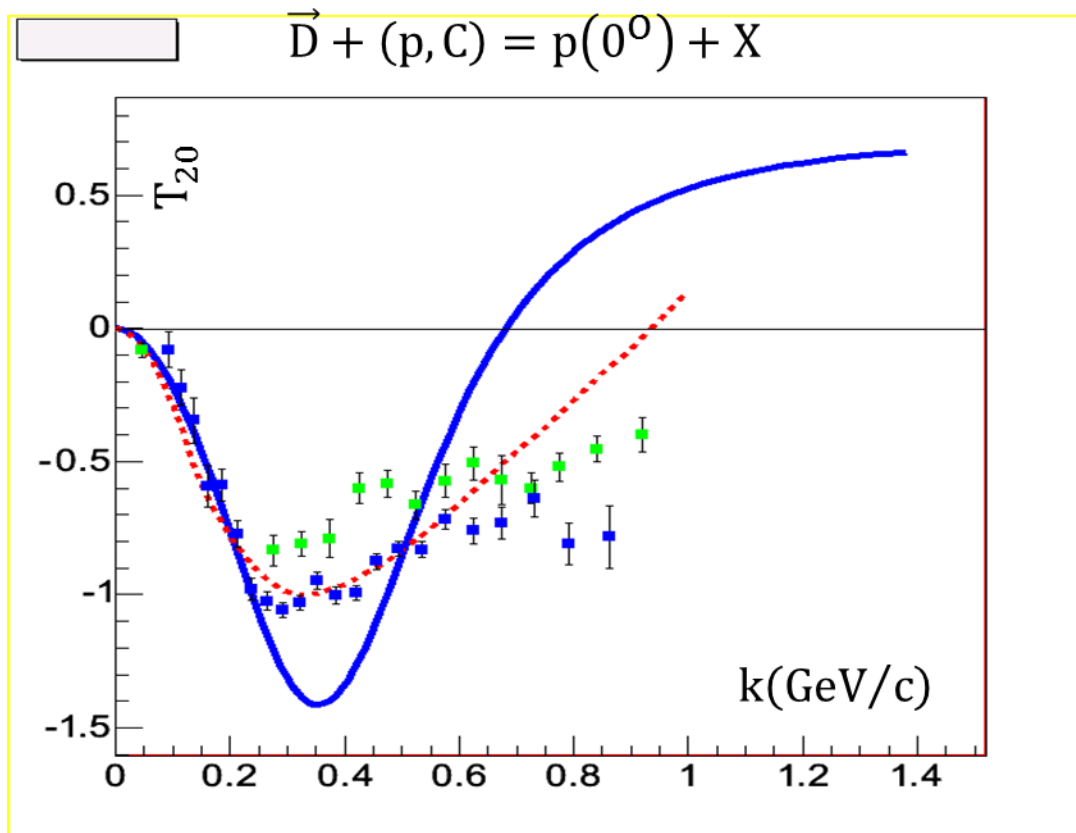
There is no obvious reason why the study of transverse asymmetry in the fixed target experiment is interesting both for a beam of transversely polarized protons and deuterons. (proton beam  $A_N \equiv A_y \leq 0.13$ , deuteron beam  $A_N \equiv A_y \equiv 0$ ).

# Cross section

V.S.Barashenkov, N. V.Slavin, PEPAN, v.15, p.997, (1984)



# Deuteron fragmentation into cumulative particle Polarization observables



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Deuteron Beams»