
Transport Study on Heavy Quarkonium production in Heavy Ion Collisions

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Pengfei Zhuang (Tsinghua U.)

Based on : PRC89,054911(arXiv:1401.5845)

JPG: Nucl.Part.Phys.41,124006(arXiv:1409.5559)

NPA931,654(arXiv:1408.3900)



Outline

- Introduction
- Transport Model
- Numerical Results
- Summary

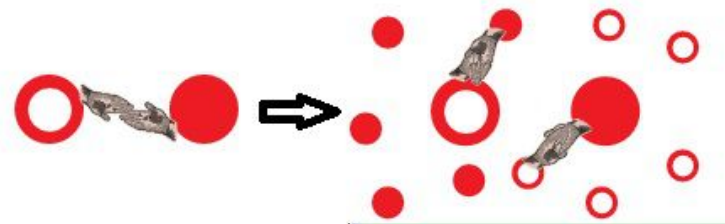
Introduction

large mass scale $m_Q \gg \Lambda_{QCD}, T$

- produced via **Hard Processes** from early stage
- "calibrated" QCD Force---**Heavy quark interaction**

- In vacuum **NR potential (or NRQCD)** e.g. $V(r) = -\alpha_c / r + kr$
---spectroscopy well described

- In medium **Color screening**



Satz and Matsui: PLB178, 416(1986):

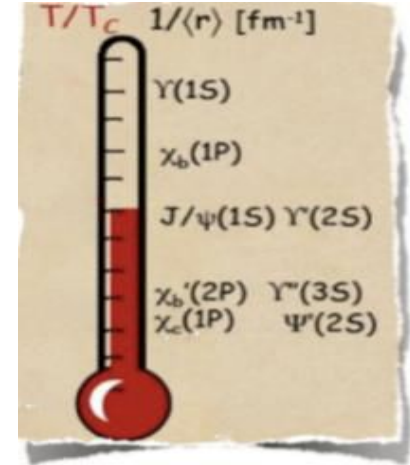
J/Psi suppression as a probe of QGP in HIC

Introduction

Thermometer

e.g for $V=U=F+TS$ (H.Satz et al, 06) F from IQCD :

state	J/ψ(1S)	χ _c (1P)	ψ'(2S)	Υ(1S)	χ _b (1P)	Υ(2S)	χ _b (2P)	Υ(3S)
T _d /T _c	2.10	1.16	1.12	> 4.0	1.76	1.60	1.19	1.17



Not so simple, many other effects affecting...

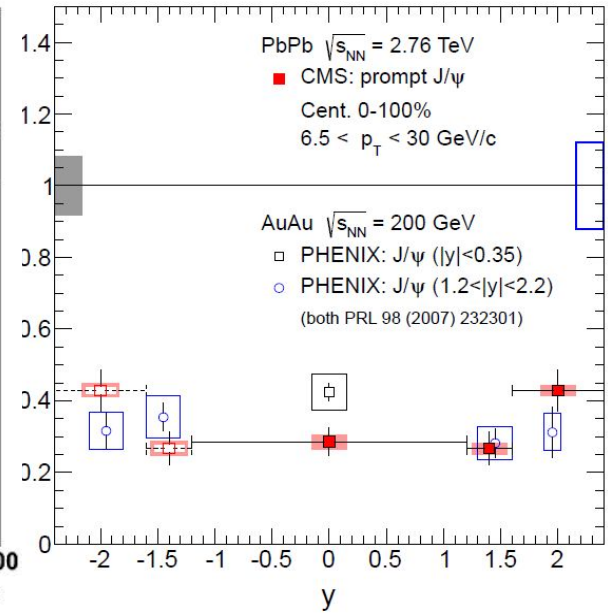
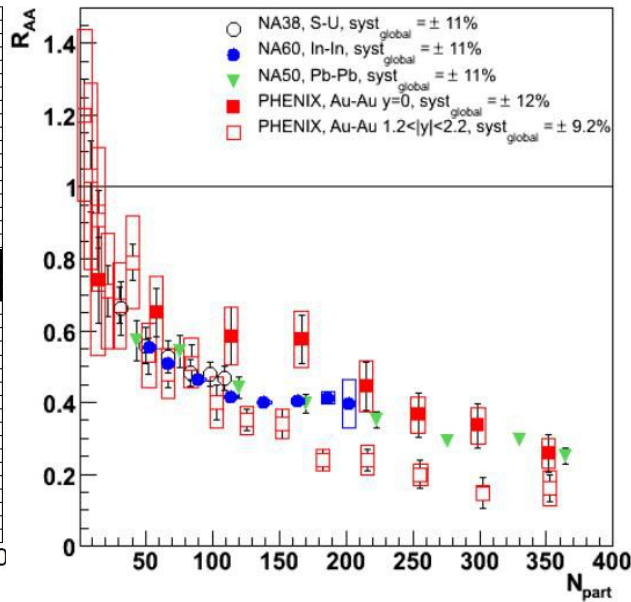
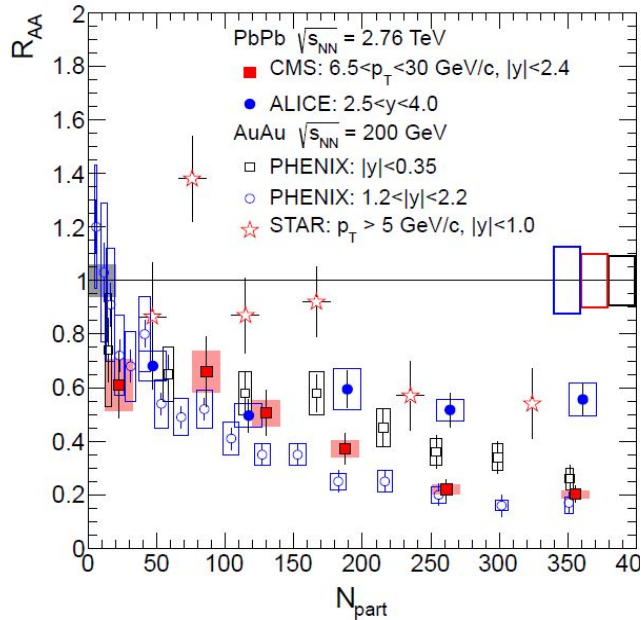
(A.Capella et al) (J.W.Cronin et al) (A.H.Mueller, R.Vogt, et al)

- Cold matter effects, nuclear absorption, Cronin, Shadowing
- collisional break-up, gluo-diss.(G.Bhanot and M.H.Peskin) quasi-free diss.(R.Rapp)
- Regeneration/Recombination (PBM, Thews, R.Rapp,et al)

● Observation $R_{AA} = \frac{N_{J/\psi}^{AA}}{N_{coll} N_{J/\psi}^{pp}} \sim \frac{"QCD_{medium}"}{"QCD_{vacuum}"}$ $\left\{ \begin{array}{l} = 1 \text{ No effect} \\ < 1 \text{ Suppression} \\ > 1 \text{ Enhancement} \end{array} \right.$

Introduction

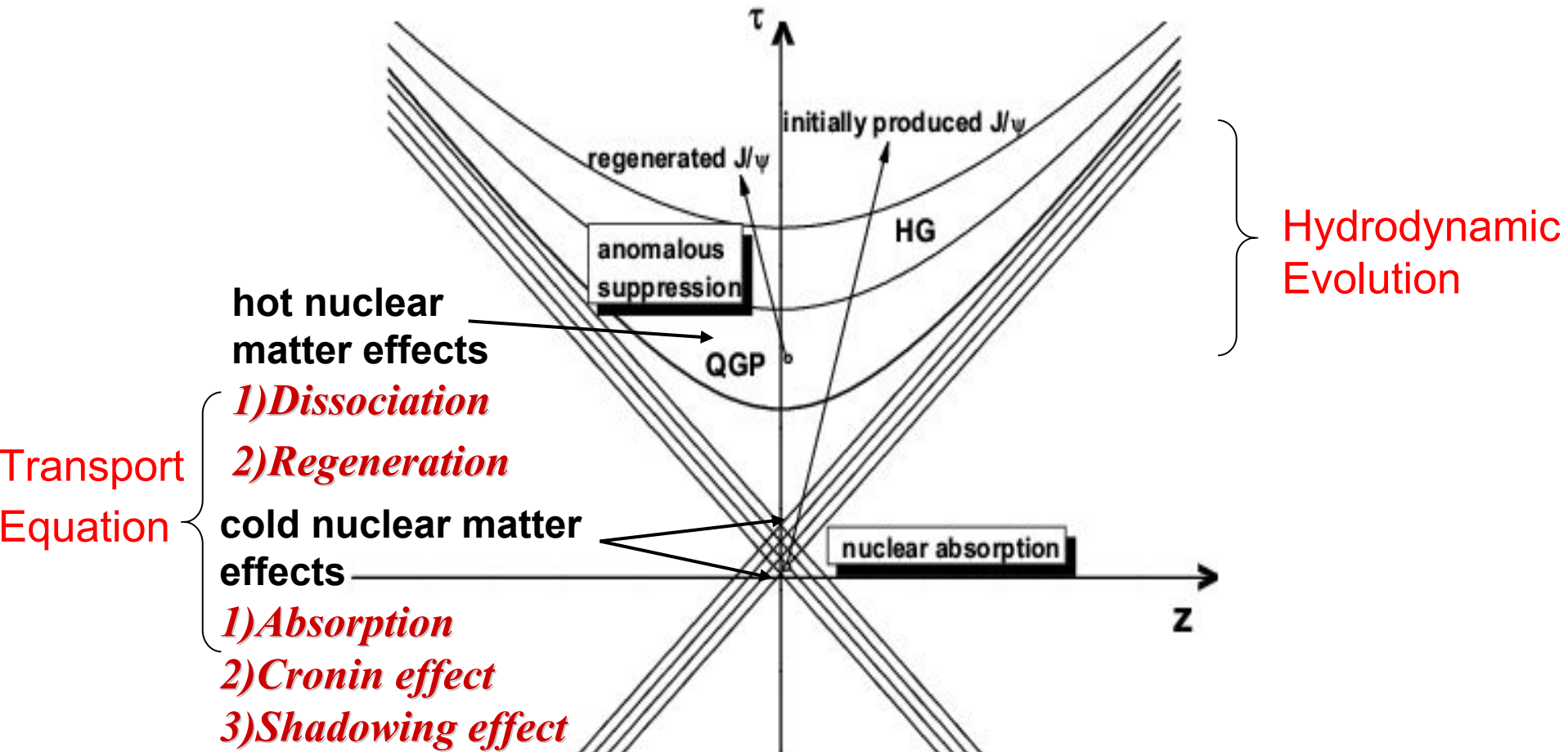
from **SPS**, to **RHIC**, Now, we are at **LHC** era



- ✓ Unified model including interplay of Cold and Hot matter effects
- ✓ With increasing coll.energy, Hot medium effects increase? where?

Transport Model

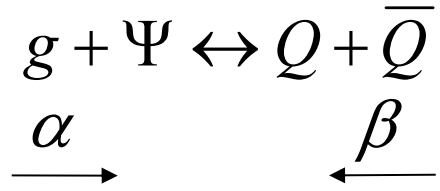
Transport(cold&hot) + Hydrodynamic



Transport Model- transport equation & hot effects

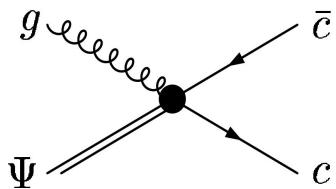
- quarkonium distribution function in phase space $f_\psi(\vec{p}, \vec{x}, t)$

$$\partial_t f + \vec{v}_T \cdot \nabla_T f + v_z \partial_z f = -\alpha f + \beta$$



1) Gluon dissociation :

$$\alpha = \frac{1}{2E_T} \int \frac{d^3 \vec{k}}{(2\pi)^3 2E_g} \sigma_{g\Psi} \cdot 4F_{g\Psi} \underline{f_g(k, x)} \leftarrow \frac{N_g}{(e^{p_g^\mu u_\mu / T} - 1)}$$



in Vacuum

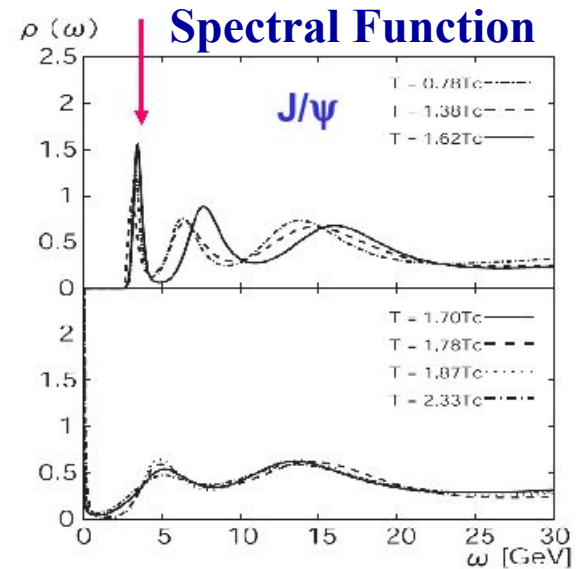
OPE (Peskin, 1979)

$$\sigma_g(\omega) = A_0 \cdot \frac{(\omega/\epsilon_\psi - 1)^{3/2}}{(\omega/\epsilon_\psi)^5}$$

$$\epsilon_\psi = \text{const, for } T_c < T < T_d,$$

in Medium

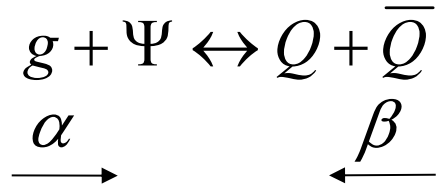
spectral peak disappear above some tem. T_d



Transport Model- transport equation & hot effects

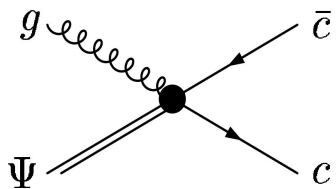
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in Vacuum

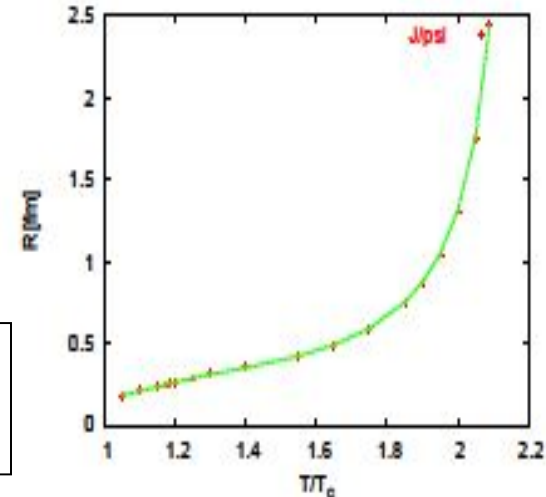
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in Medium

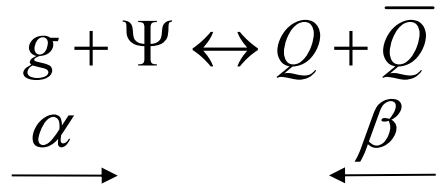
$$\sigma_{g\Psi}(T) = \sigma_{g\Psi}(T=0) \frac{\langle r_\Psi^2 \rangle(T)}{\langle r_\Psi^2 \rangle(T=0)}$$



Transport Model- hot nuclear matter effects

- quarkonium distribution function in phase space $f_\Psi(\vec{p}, \vec{x}, t)$

$$\partial_t f + \vec{v}_T \cdot \nabla_T f + v_z \partial_z f = -\alpha f + \beta$$



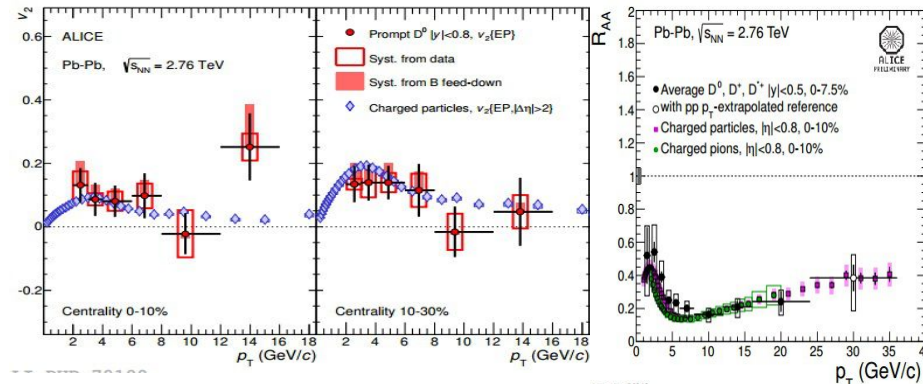
2) in-Medium Regeneration :

$$\beta = \frac{1}{2m_t} \int \frac{d^3\vec{k}}{(2\pi)^3 2E_g} \frac{d^3\vec{q}_1}{(2\pi)^3 2E_Q} \frac{d^3\vec{q}_2}{(2\pi)^3 2E_{\bar{Q}}} (2\pi)^4 \delta^4(p+k-q_1-q_2) W_{pro}(s) f_Q(k, x) f_{\bar{Q}}(k, x)$$

- Detailed balance : $\sigma_{reg.}(s) = \frac{4}{3} \frac{(s - m_\Psi^2)^2}{s(s - 4m_Q^2)} \sigma_{diss.}(s)$

- Distribution for heavy quarks is assumed (strong interaction limit) as:

$$f_Q(q, x) = N(x) \rho_Q(x) / (e^{u \cdot q/T} + 1)$$



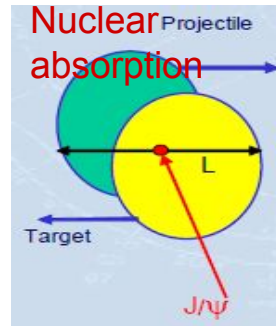
Transport Model- cold nuclear matter effects

- Initial condition $f(\vec{p}, \vec{x}, t_0)$ for transport

Glauber superposition from pp collisions along with modification from CNMs:

Cold Effects

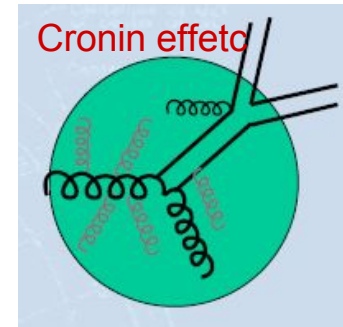
Absorption



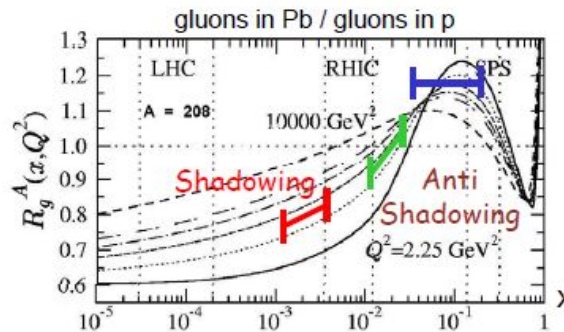
$t_{coll} \ll t_{\Psi}$ so it's neglected at LHC

Cronin

Gaussian smearing treatment



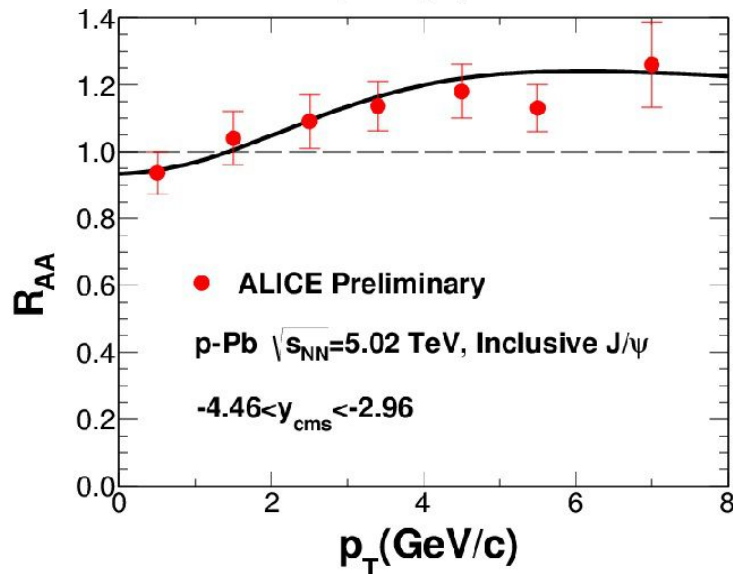
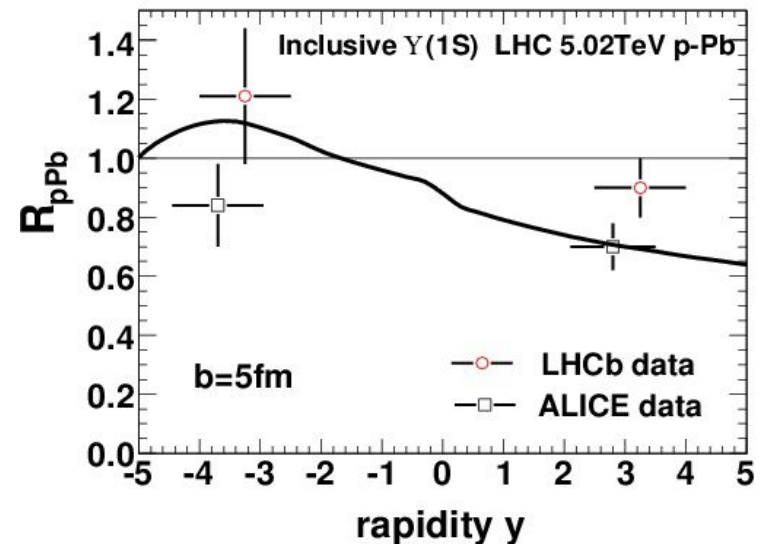
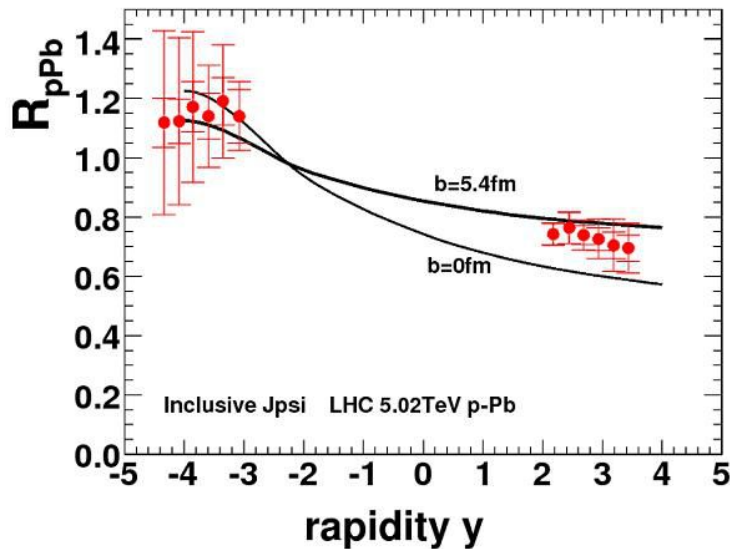
Shadowing



nPDF vs. free PDF

R.Vogt et al. PRL91 (2003)
142301.PRC71(2005) 05490

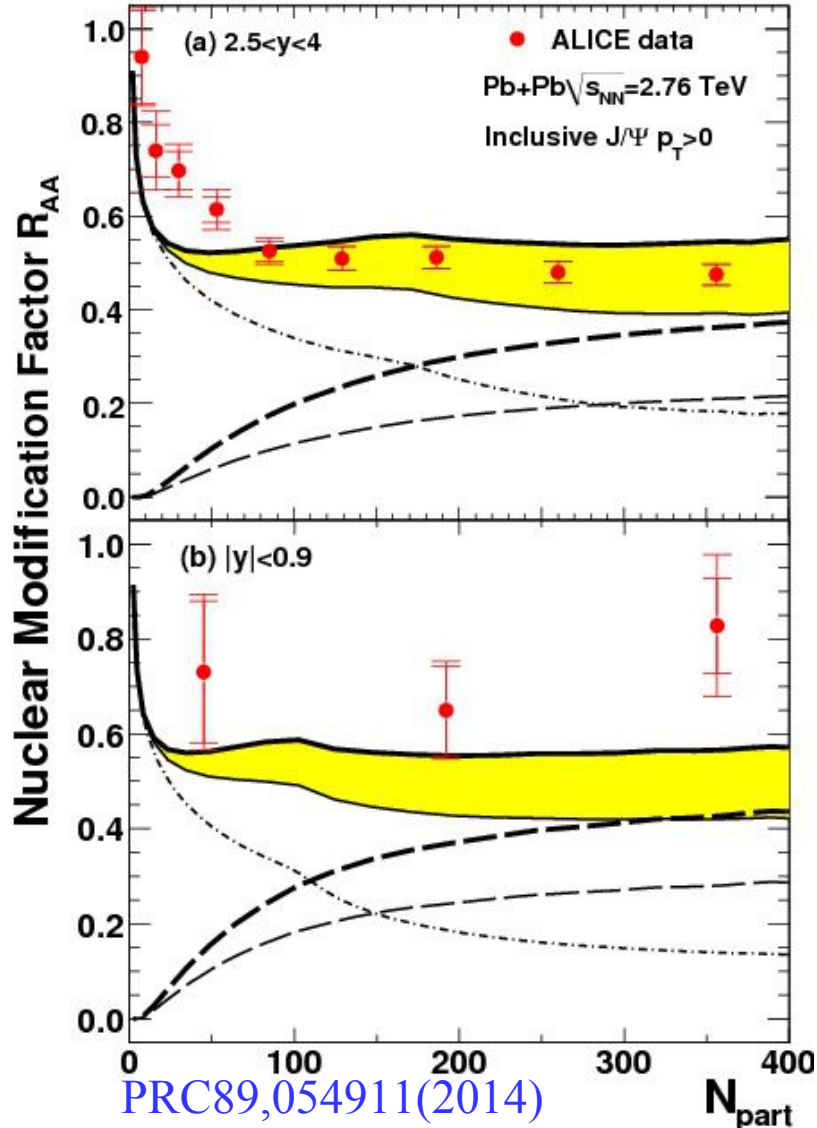
Transport Model- test of cold matter in p-Pb



p-Pb 5.02 TeV

Cronin + Shadowing(EKS98)
can describe the p-Pb(5.02 TeV) data well !

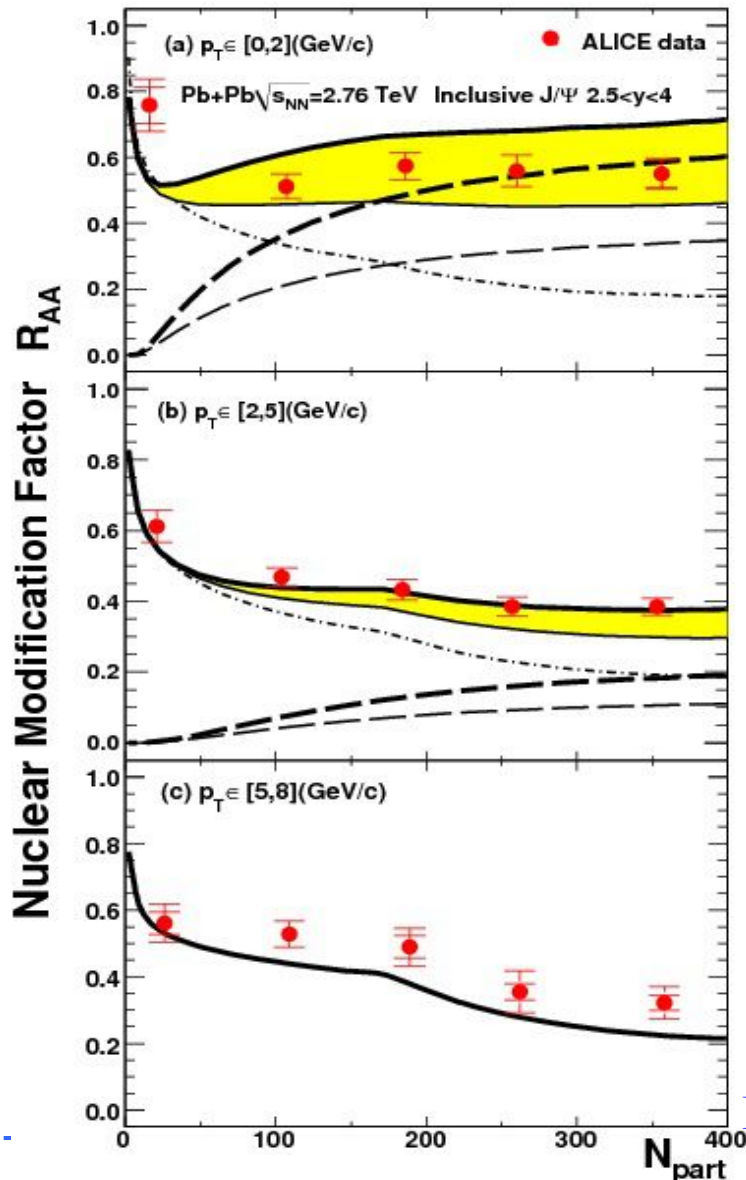
Results—Yield's Centrality depen. (all pT)



1、 **Regeneration** plays an important roll in most of centralities, and can be dominant.

2、 Competition leads to **platform** structure in most centralities.

Results—Yield's Centrality depen. (pT bin)



Forward Rapidity

1、flat structure gradually disappears with pT.---->

Regeneration is mostly contributed in **low pT** part.

2、Jpsi naturally provide two probes:

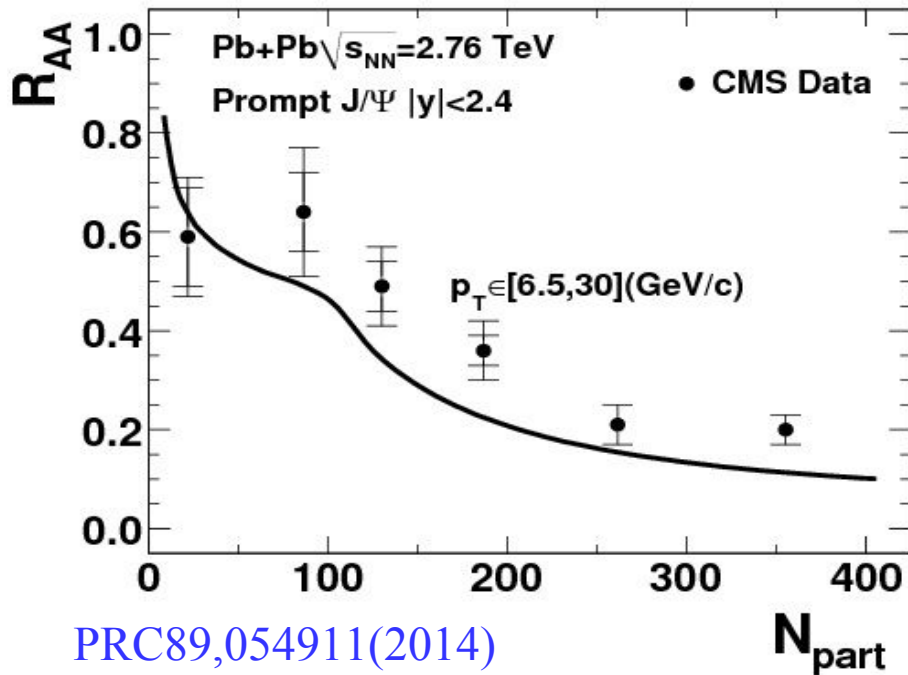
a) **Hard Probe**: high pT, Color Screening

b) **Soft Probe**: low pT, Thermalization

[PRC89,054911\(2014\)](#)

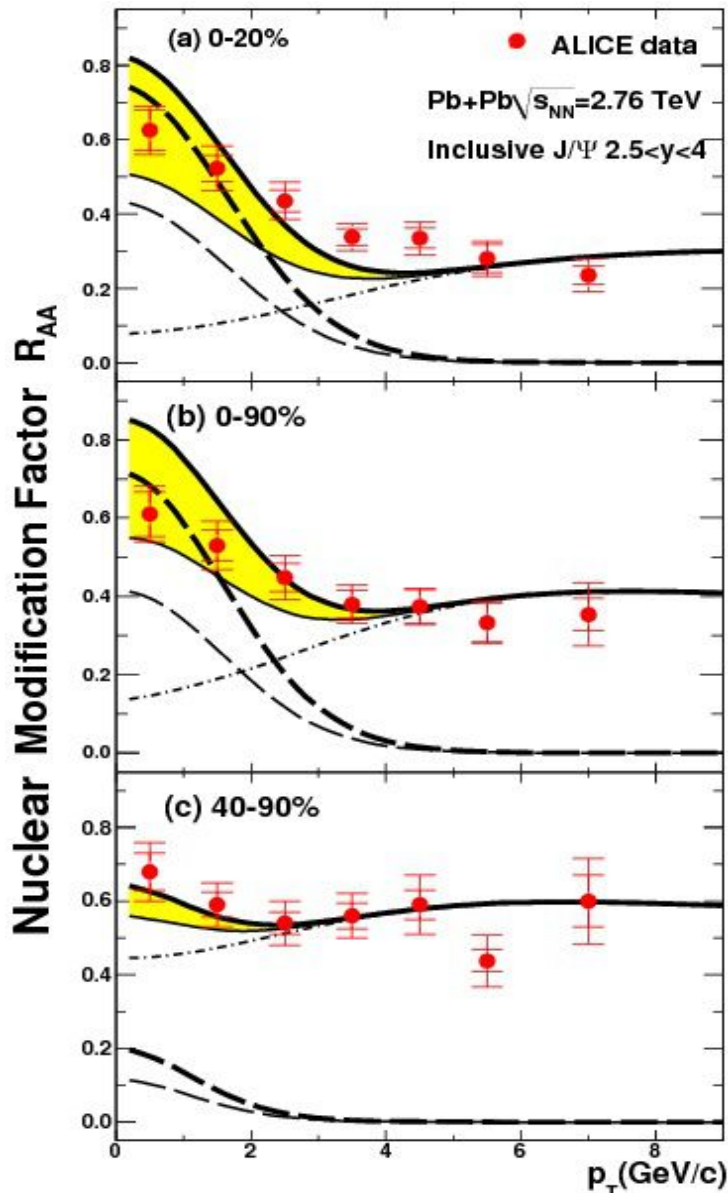
Results—Yield's Centrality depen. (pT bin)

Mid-Rapidity



Note the **"kink"**-----
Melting Temperature from
Color Screening

Results—RAA(pT)



Forward Rapidity

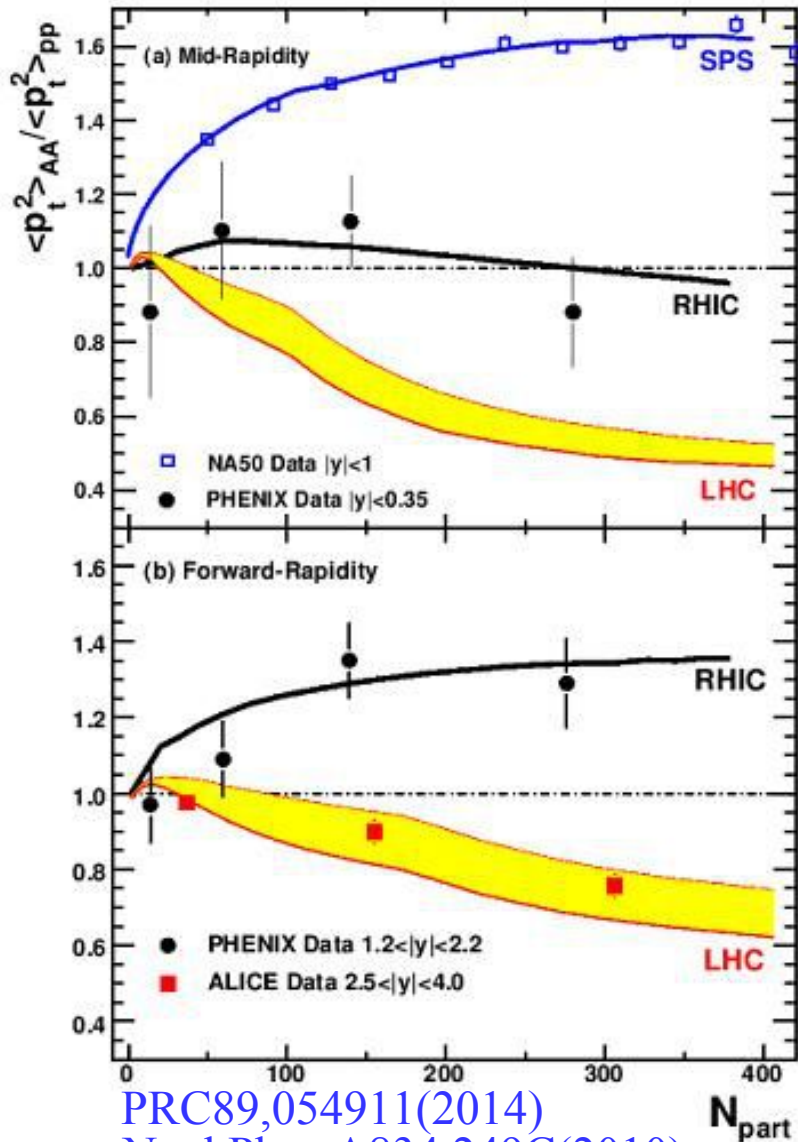
1、Regeneration -- low p_T
Initial -- high p_T .

2、Competition leads to slightly
"Valley" structure

3、The decreasing behavior
indicates regeneration

PRC89,054911(2014)

Results—Modification for Trans. pT: rAA



SPS: Cronin effect for *initial*

RHIC: competition betw. *initial & regeneration*

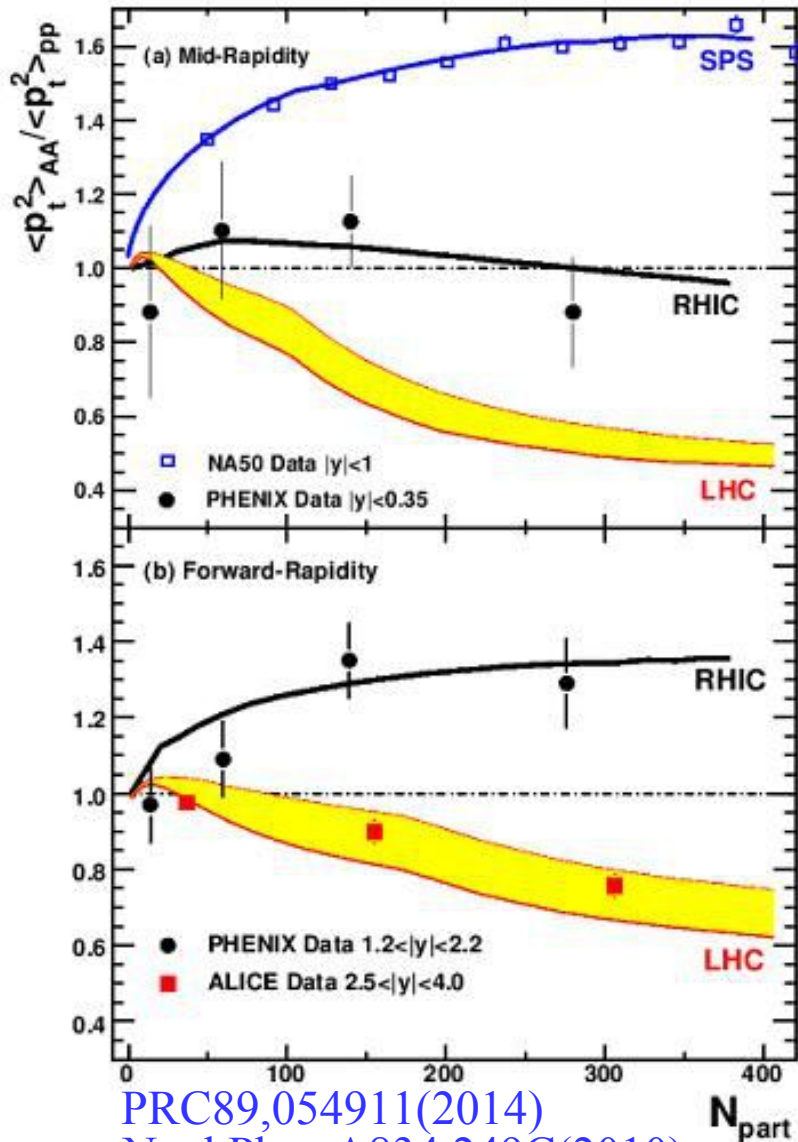
LHC: dominant *regeneration*

$$r_{AA} = \frac{\langle p_T^2 \rangle_{AA}}{\langle p_T^2 \rangle_{pp}}$$

PRC89,054911(2014)

Nucl.Phys.A834,249C(2010)

Results—Modification for Trans. pT: rAA

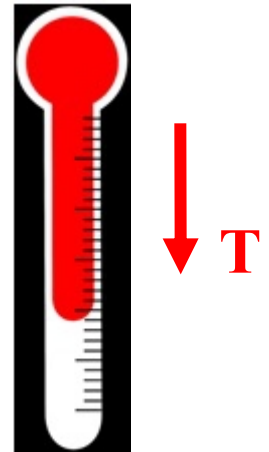


$\sqrt{s_{NN}} \uparrow$ $QGP \uparrow$

hotter

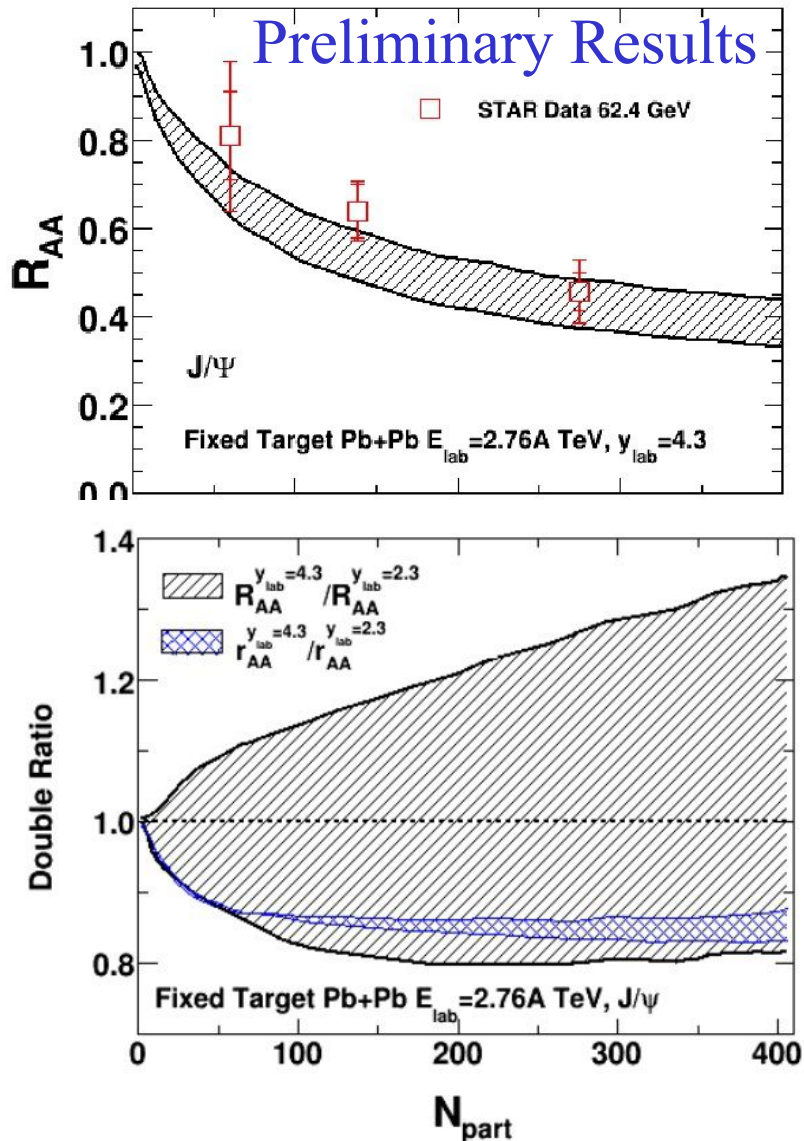
$$r_{AA} = \frac{\langle p_T^2 \rangle_{AA}}{\langle p_T^2 \rangle_{pp}}$$

Clearly shows
a hotter medium
been created at
LHC !



PRC89,054911(2014)
Nucl.Phys.A834,249C(2010)

Fixed Target Pb+Pb 2.76A TeV (AFTER) $\sim \sqrt{s_{NN}} = 72 GeV$



lower border : w/o Shadowing
 upper border : with Shadowing

$$\Delta y = \tanh^{-1} \beta_{cms} = 4.3$$

mid-y (lab-y=4.3) : Anti-shadowing
 for-y (lab-y=2.3) : Shadowing

Sensitive probe to gluon distribution

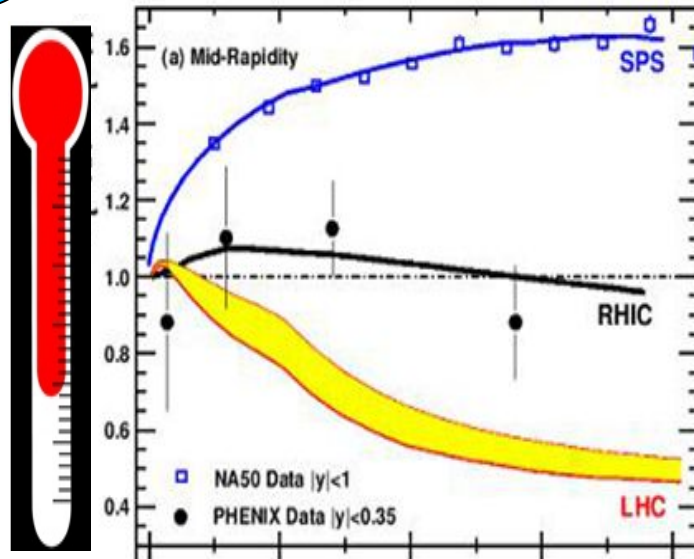
Summary

$$r_{AA} = \langle p_T^2 \rangle_{AA} / \langle p_T^2 \rangle_{pp}$$

cold? hot?



"heavy quarkonia cat"



not that hot

a little hot

very hot !

- ❑ Finite chemical potential 's effects, which is relevant for lower energy case : FAIR
- ❑ Event-by-event simulation for both open and hidden heavy mesons

Thank You !

Transport Model- ideal Hydro dynamics

● 2+1D hydrodynamics

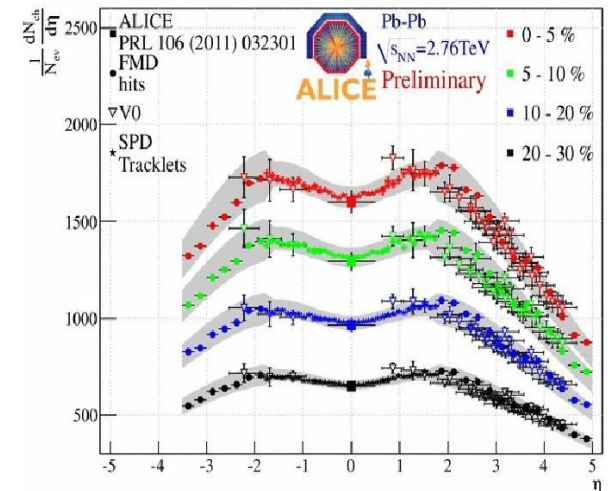
$$\begin{cases} \partial_{\mu} T^{\mu\nu} = 0 \\ \partial_{\mu} j^{\mu} = 0 \end{cases} \quad \begin{cases} T^{\mu\nu} = (\varepsilon + p)u^{\mu}u^{\nu} - g^{\mu\nu} p \\ j^{\mu} = nu^{\mu} \end{cases}$$

● Equation Of State:

Ideal Gas with quarks and gluons for QGP
& **HRG** for Hadronic phase

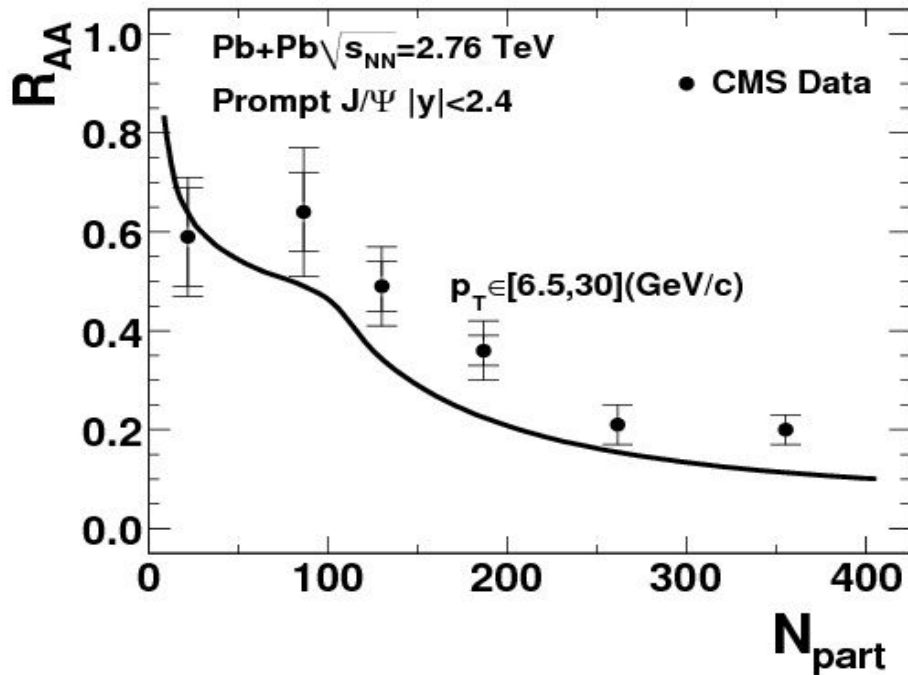
● Initialization:

Glauber model & constrained by
fitting **Charged Multiplicities**



Backup—Yield's Centrality depen. (pT bin)

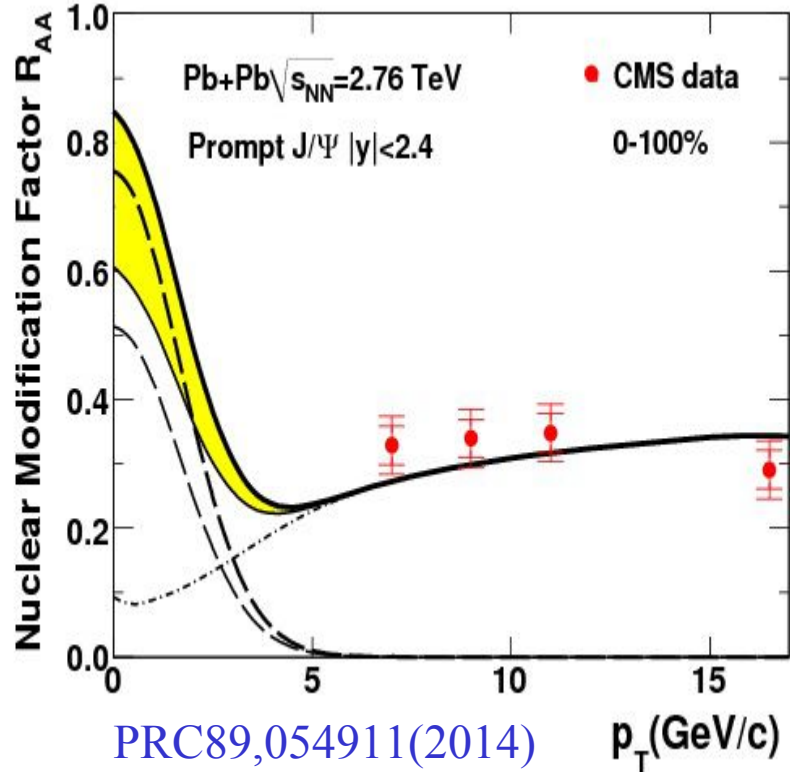
Mid-Rapidity



Note the "kink"-----
Melting Temperature from
Color Screening

PRC89,054911(2014)

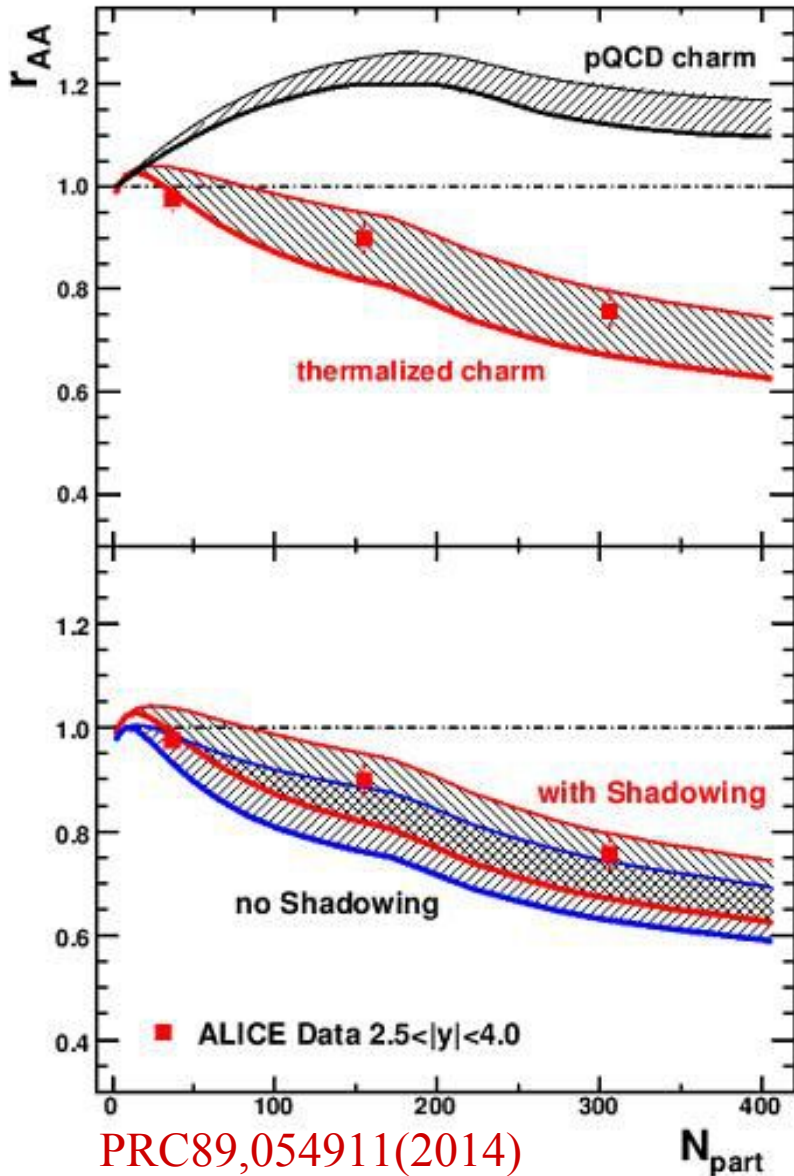
Backup—RAA(pT)



Mid-Rapidity

"Valley" structure more clearly

Backup-- Trans. pT: rAA



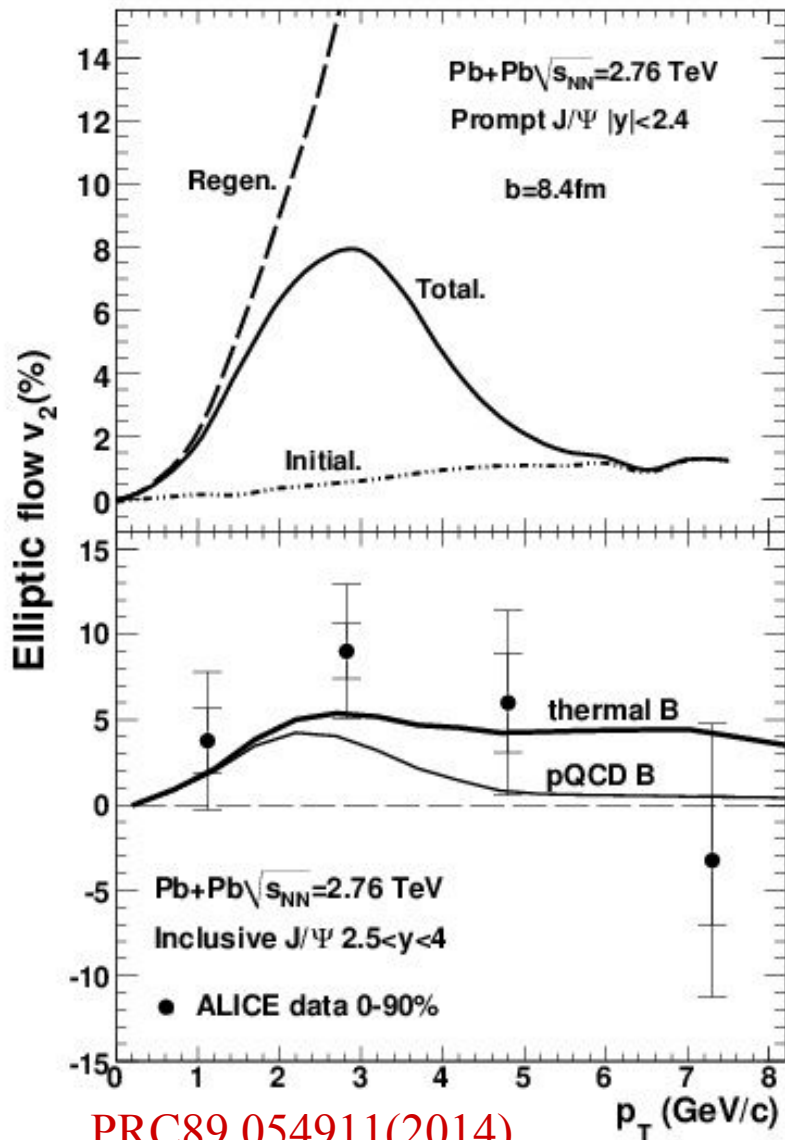
$$r_{AA} = \frac{\langle p_T^2 \rangle_{AA}}{\langle p_T^2 \rangle_{pp}}$$

1, sensitive to the degree of heavy quark thermalization --energy loss.

2, not sensitive to the cold nuclear matter effect----- Shadowing effect.

clearly indicates QGP's medium effects

Results—Elliptic flow v_2



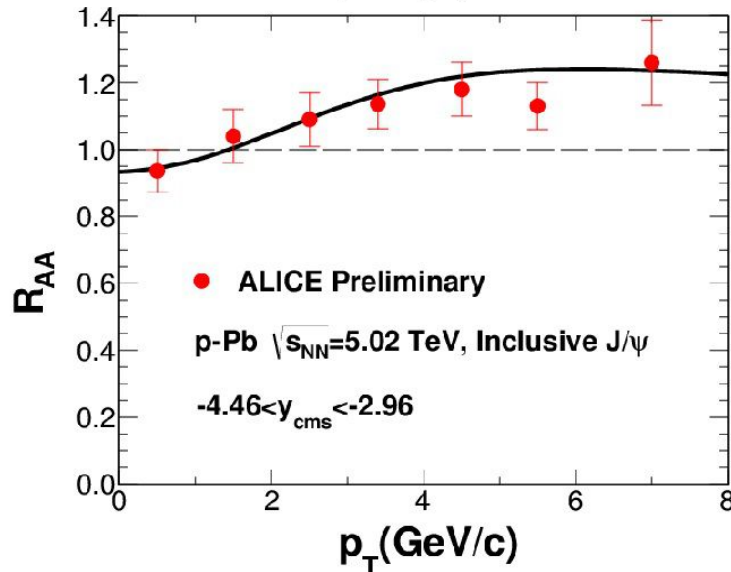
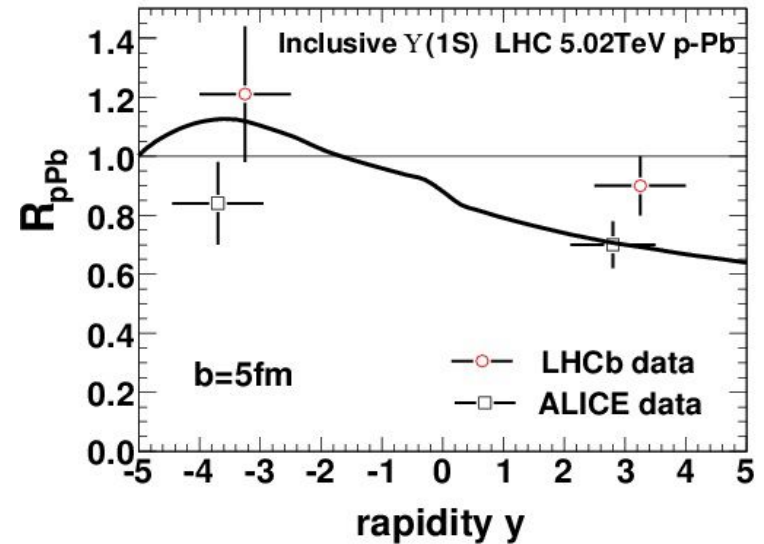
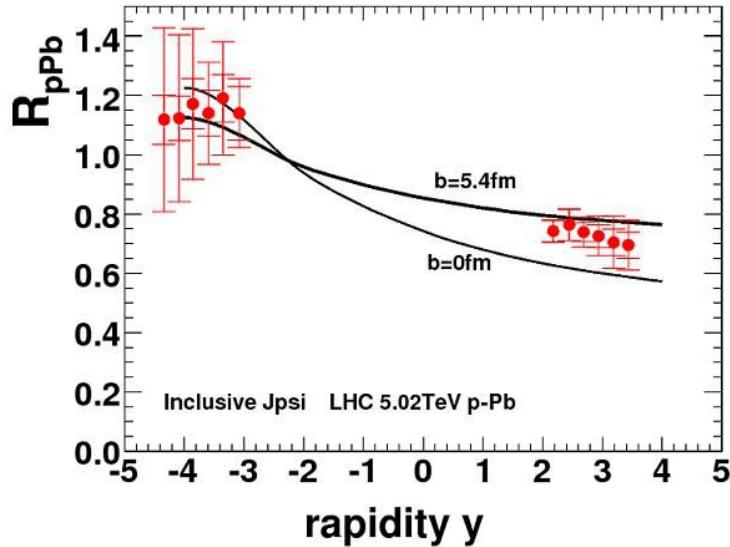
PRC89,054911(2014)

1, remarkable v_2 due to important regeneration !
-->Reflects energy loss of heavy quark.

2, "Ridge" structure from two-component:
{ hard (initial, jet)
{ soft (regeneration, hydro)

3, v_2 for high p_T indicates B meson's thermalization

Transport Model- test of cold matter in p-Pb



p-Pb 5.02 TeV

Cronin + Shadowing(EKS98)
can describe the p-Pb(5.02
TeV) data well !

Transport Model- solution of transport equation

$$\left[\cosh(y - \eta) \frac{\partial}{\partial \tau} + \frac{1}{\tau} \sinh(y - \eta) \frac{\partial}{\partial \eta} + \vec{v}_t \cdot \vec{\nabla}_t \right] f = -\alpha f + \beta$$

$$\begin{aligned} & f(\vec{p}_t, y, \vec{x}_t, \eta, \tau) \\ = & f(\vec{p}_t, y, \vec{r}_t(\tau_0), Y(\tau_0), \tau_0) e^{-\int_{\tau_0}^{\tau} d\tau' A(\vec{p}_t, y, \vec{r}_t(\tau'), Y(\tau'), \tau')} \\ & + \int_{\tau_0}^{\tau} d\tau' B(\vec{p}_t, y, \vec{r}_t(\tau'), Y(\tau'), \tau') e^{-\int_{\tau'}^{\tau} d\tau'' A(\vec{p}_t, y, \vec{r}_t(\tau''), Y(\tau''), \tau'')} \end{aligned}$$

$$\vec{v}_t = \frac{\vec{p}_t}{E_t}$$

$$\vec{r}_t(\tau') = \vec{x}_t - \vec{v}_t [\tau \cosh(y - \eta) - \tau' \cosh(\Delta(y - \eta))]$$

$$Y(\tau') = y - \Delta(y - \eta)$$

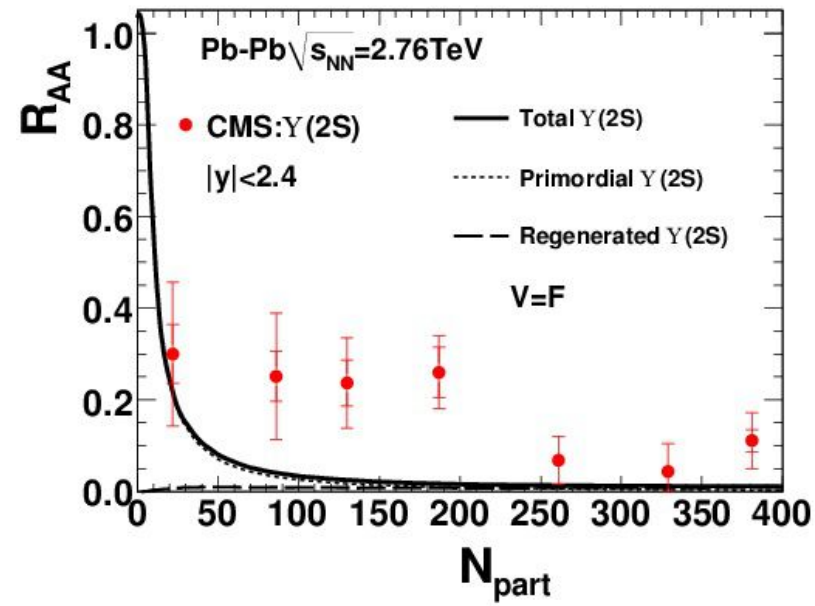
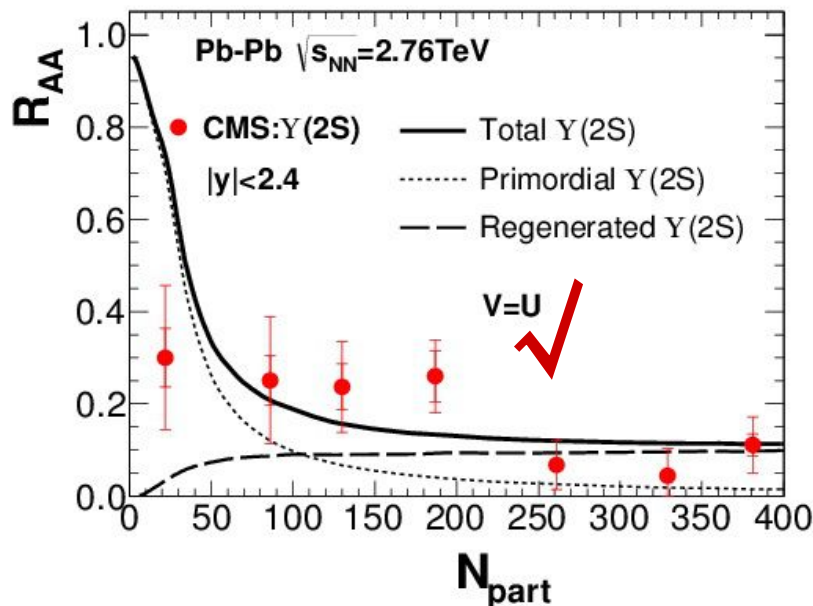
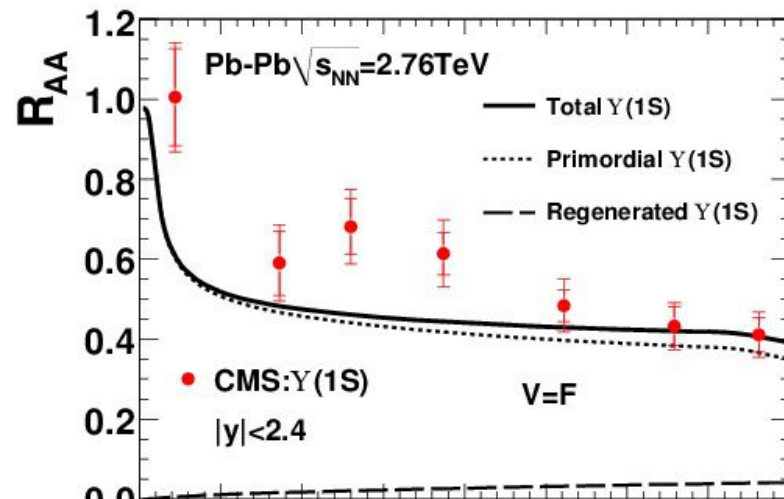
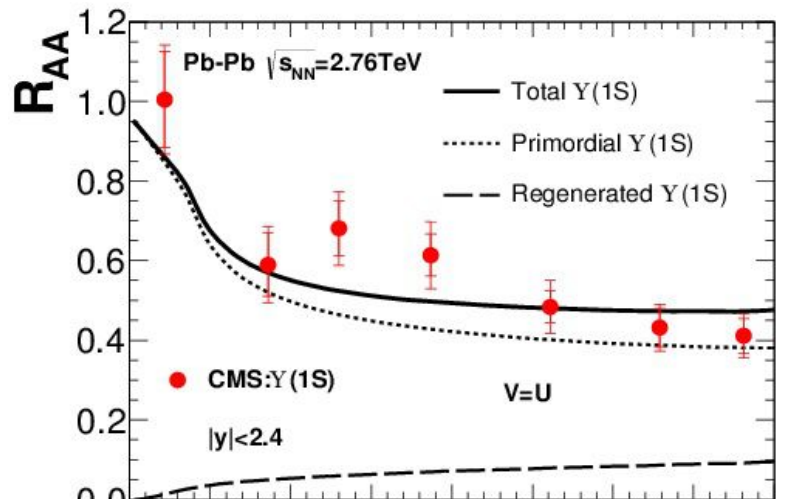
$$A(\vec{p}_t, y, \vec{r}_t(\tau'), Y(\tau'), \tau') = \frac{\alpha(\vec{p}_t, y, \vec{r}_t(\tau'), Y(\tau'), \tau')}{\cosh(\Delta(y - \eta))}$$

$$B(\vec{p}_t, y, \vec{r}_t(\tau'), Y(\tau'), \tau') = \frac{\beta(\vec{p}_t, y, \vec{r}_t(\tau'), Y(\tau'), \tau')}{\cosh(\Delta(y - \eta))}$$

$$\Delta(y - \eta) \equiv \operatorname{arcsinh}\left(\frac{\tau}{\tau'} \sinh(y - \eta)\right)$$

Both Initial production and Regeneration suffers **Suppression**

Results—Bottomonium differs $V=U$ or $V=F$



Transport Model- cold nuclear matter effects

Absorption $\times e^{-\sigma_{abs}(T_A(\vec{x}_T, z_A, +\infty) + T_B(\vec{x}_T - \vec{b}, -\infty, z_B))}$

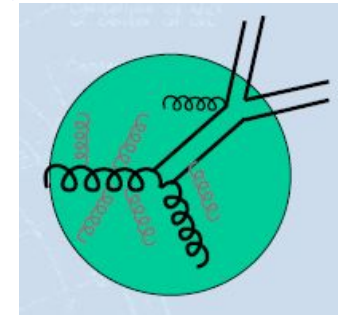
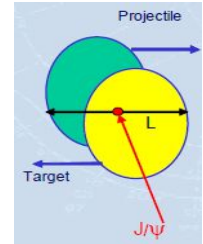
$t_{coll} \ll t_{\Psi}$ (so at LHC can safely be neglected)

Cronin p_T broadening (use Gaussian smearing)

$$\bar{f}_{pp}(\vec{p}_T, \vec{x}_T, z_A, z_B) = \frac{1}{\pi a_{gN} \cdot l(\vec{x}_T, z_A, z_B)} \int d^2 \vec{p}'_T e^{-\frac{p_T'^2}{a_{gN} \cdot l(\vec{x}_T, z_A, z_B)}} f_{pp}(|\vec{p}_T - \vec{p}'_T|)$$

$$a_{gN} = \Delta^2(\mu) \sigma_{pp}^{inelastic} \rho_0$$

$$a_{gN} = 0.15 \text{GeV}^2 / c^2 \quad \underline{\text{@ LHC Pb-Pb 2.76TeV}}$$



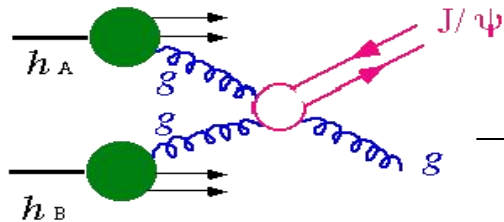
Init.J.Mod.Phys.E.12,211(2003)

Phys.Rev. C 73, 014904(2006)

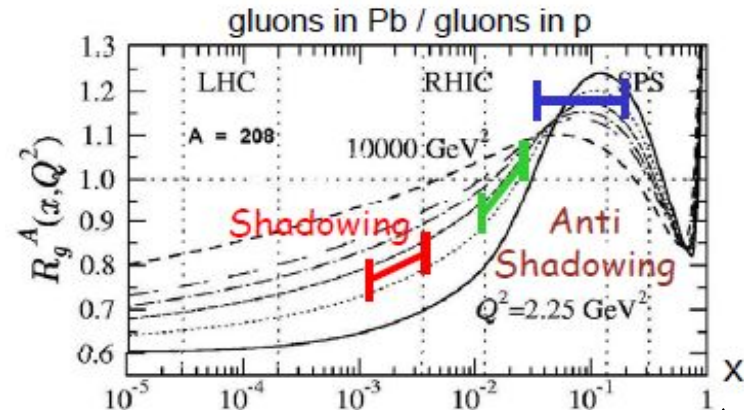
Transport Model- cold nuclear matter effects

Shadowing $R_g^A(x, \mu_F) = \frac{f_g^A(x, \mu_F)}{A f_g^{\text{Nucleon}}(x, \mu_F)}$

for open & hidden heavy mesons



(2->1)process
Color Evaporation Model



$$x_{1,2}^g = \frac{\sqrt{m_{c\bar{c}}^2 + p_T^2}}{\sqrt{s_{NN}}} e^{\pm y}$$

pp $\frac{d\sigma_{pp}^\Psi}{dp_T^\Psi dy_\Psi} = \int dy_g x_1 x_2 \cdot f_g(x_1, \mu_F) f_g(x_2, \mu_F) \frac{d\sigma_{gg \rightarrow \psi g}}{dt}$

AA $f_0(\vec{p}, \vec{x}_T) = \frac{(2\pi)^3}{E_T^\Psi \cosh y_\Psi} \frac{d\sigma_{pp}^\Psi}{dy} \int dz_A dz_B \rho_A(\vec{x}_T, z_A) \cdot$

$$\rho_B(\vec{x}_T - \vec{b}, z_b) \mathcal{R}_g(\vec{x}_T, x_1, \mu_f) \cdot$$

$$\mathcal{R}_g(\vec{x}_T - \vec{b}, x_2, \mu_f) \bar{f}_{pp}(\vec{p}_T, \vec{x}_T, z_A, z_B)$$

$$\mathcal{R}_g(\vec{x}_T, x, \mu_f) = 1 + N_{A,\rho} [R_g^A(x, \mu_f) - 1] \frac{T_A(\vec{x}_T)}{T_A(0)}$$

R.Vogt et al. PRL91 (2003) 142301.

PRC71(2005) 054902

Transport Model- ideal Hydro dynamics

● 2+1D hydrodynamics($\mu_B = 0$)

$$\left\{ \begin{array}{l} \partial_\tau \rho_T + \nabla_T \cdot (\rho_T \vec{v}_T) = 0 \quad (\rho_T(\vec{x}_T, \tau) = \tau \cdot n_{c\bar{c}}^{Lab}) \leftarrow \text{kinetic thermalization for HQ} \\ \partial_\tau E + \nabla_T \cdot \vec{M}_T = -(E + p) / \tau \\ \partial_\tau M_x + \nabla_T \cdot (M_x \vec{v}_T) = -M_x / \tau - \partial_x p \\ \partial_\tau M_y + \nabla_T \cdot (M_y \vec{v}_T) = -M_y / \tau - \partial_y p \end{array} \right. \leftarrow \left\{ \begin{array}{l} \partial_\mu T^{\mu\nu} = 0 \\ \text{Boost Invariance in z-direction} \end{array} \right.$$

$$E = (\varepsilon + p)\gamma^2 - p \quad \vec{M} = (\varepsilon + p)\gamma^2 \vec{v}$$

● Equation Of State:

Ideal Gas with quarks and gluons for QGP
& **HRG** for Hadronic phase

● Initialization:

Glauber model & constrained by fitting **Charged Multiplicities**

