



The Tsallis Distribution at Large Transverse Momenta

J. Cleymans

University of Cape Town, South Africa

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Work done in collaboration with:
M.D. Azmi

C.Y. Wong, G. Wilk, J.L. Cirto, C. Tsallis
Phys. Rev. D91 (2015) 114027

C.Y. Wong, G. Wilk,
Acta Physica Polonica, B43 (2012) 2047-2054



Outline

Tsallis Distribution

Transverse Momentum Distributions

Strange Particles in ALICE

Conclusion

Tsallis Thermodynamics

The Tsallis distribution is given by

$$f(E) = \left[1 + (q - 1) \frac{E - \mu}{T} \right]^{-\frac{1}{q-1}},$$

and the thermodynamic quantities N , E , P , S , ... are integrals over this distribution.

Asymptotically

$$\lim_{E \rightarrow \infty} f(E) = \left(\frac{E}{T} \right)^{-\frac{1}{q-1}}$$

scale is set by T

asymptotic behaviour is set by q .



For high energy physics a consistent form of Tsallis thermodynamics for the particle number, energy density and pressure is given by

$$\begin{aligned} N &= gV \int \frac{d^3p}{(2\pi)^3} \left[1 + (q-1) \frac{E - \mu}{T} \right]^{-\frac{q}{q-1}}, \\ \epsilon &= g \int \frac{d^3p}{(2\pi)^3} E \left[1 + (q-1) \frac{E - \mu}{T} \right]^{-\frac{q}{q-1}}, \\ P &= g \int \frac{d^3p}{(2\pi)^3} \frac{p^2}{3E} \left[1 + (q-1) \frac{E - \mu}{T} \right]^{-\frac{q}{q-1}}. \end{aligned}$$

where T and μ are the temperature and the chemical potential, V is the volume and g is the degeneracy factor. This introduces only one new parameter q which for transverse momentum spectra is always close to 1.



Thermodynamic consistency

$$dE = -pdV + TdS + \mu dN$$

Inserting $E = \epsilon V$, $S = sV$ and $N = nV$ leads to

$$d\epsilon = Tds + \mu dn$$

$$dP = nd\mu + sdT$$

In particular

$$n = \left. \frac{\partial P}{\partial \mu} \right|_T, \quad s = \left. \frac{\partial P}{\partial T} \right|_{\mu}, \quad T = \left. \frac{\partial \epsilon}{\partial s} \right|_n, \quad \mu = \left. \frac{\partial \epsilon}{\partial n} \right|_s.$$

are satisfied.



Thermodynamic consistency: an example

$$\begin{aligned}
 \frac{\partial P}{\partial \mu} &= gV \int \frac{d^3 p}{(2\pi^3)} \frac{p^2}{3E} \frac{\partial}{\partial \mu} f^q \\
 &= -gV \int \frac{d^3 p}{(2\pi^3)} \frac{p^2}{3E} \frac{d}{dE} f^q \\
 &= -gV \frac{4\pi}{(2\pi^3)} \int_0^\infty dp \frac{p^4}{3E} \frac{d}{dE} f^q \\
 &= -gV \frac{4\pi}{(2\pi^3)} \int_0^\infty dp \frac{p^3}{3} \frac{d}{dp} f^q \quad \text{using } EdE = pdp \\
 &= gV \frac{4\pi}{(2\pi^3)} \int_0^\infty dp p^2 f^q \\
 &= n
 \end{aligned}$$



In the Tsallis distribution the total number of particles is given by:

$$N = gV \int \frac{d^3p}{(2\pi)^3} \left[1 + (q-1) \frac{E - \mu}{T} \right]^{-\frac{q}{q-1}}.$$

The corresponding momentum distribution is given by

$$E \frac{dN}{d^3p} = gVE \frac{1}{(2\pi)^3} \left[1 + (q-1) \frac{E - \mu}{T} \right]^{-\frac{q}{q-1}},$$

which, in terms of the rapidity and transverse mass variables, $E = m_T \cosh y$, becomes (at mid-rapidity $y = 0$ and for $\mu = 0$)

$$\left. \frac{d^2N}{dp_T dy} \right|_{y=0} = gV \frac{p_T m_T}{(2\pi)^2} \left[1 + (q-1) \frac{m_T}{T} \right]^{-\frac{q}{q-1}},$$

J.C. and D. Worku, J. Phys. **G39** (2012) 025006;
arXiv:1203.4343[hep-ph].



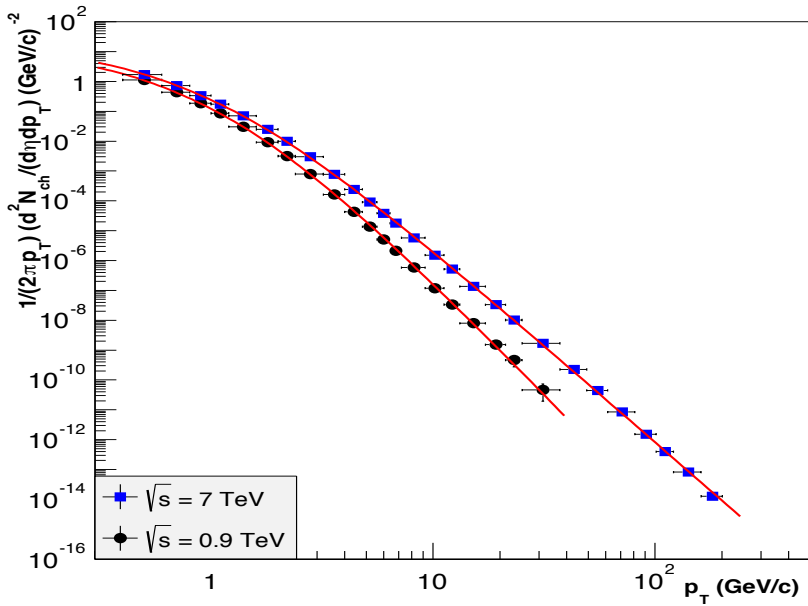
For charged particles use:

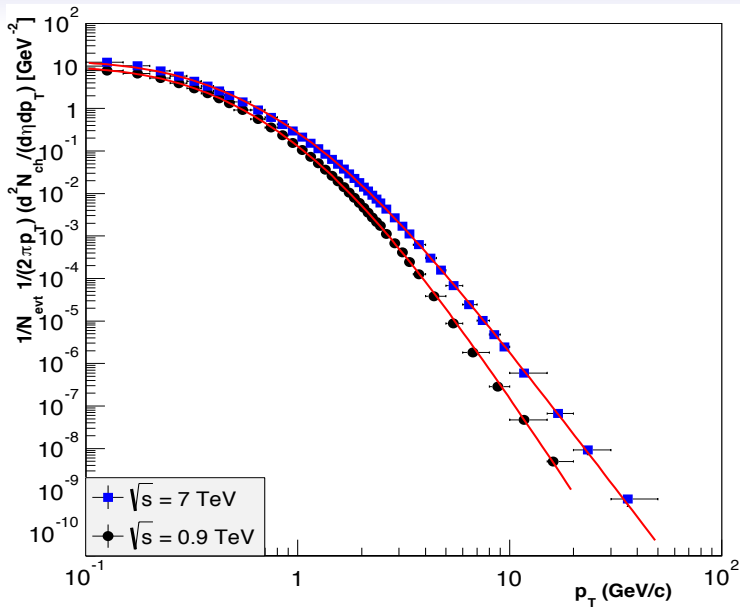
$$\left. \frac{d^2 N(\text{charged})}{dp_T dy} \right|_{y=0} = \sum_{i=\pi, K, p, \dots} g_i V \frac{p_T m_T}{(2\pi)^2} \left[1 + (q-1) \frac{m_T}{T} \right]^{-\frac{q}{q-1}},$$

M.D. Azmi and J.C. , arXiv:1501.07217v3[hep-ph].



Tsallis Distribution p-p CMS





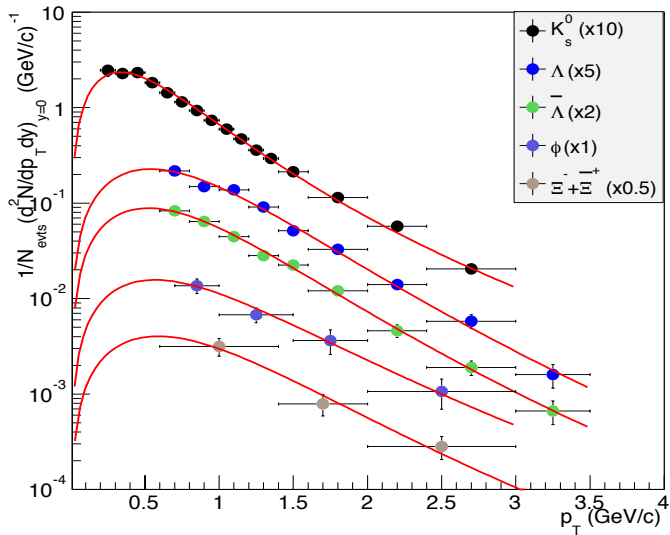
Tsallis Distribution p-p

| Experiment | \sqrt{s} (TeV) | q | T (MeV) |
|-------------------|------------------|-------------------|------------------|
| ATLAS | 0.9 | 1.129 ± 0.005 | 74.21 ± 3.55 |
| ATLAS | 7 | 1.150 ± 0.002 | 75.00 ± 3.21 |
| CMS | 0.9 | 1.129 ± 0.003 | 76.00 ± 0.17 |
| CMS | 7 | 1.153 ± 0.002 | 73.00 ± 1.42 |

Values of the q and T parameters to fit the p_T spectra measured by the ATLAS and CMS collaborations.



Tsallis fits to strange particles



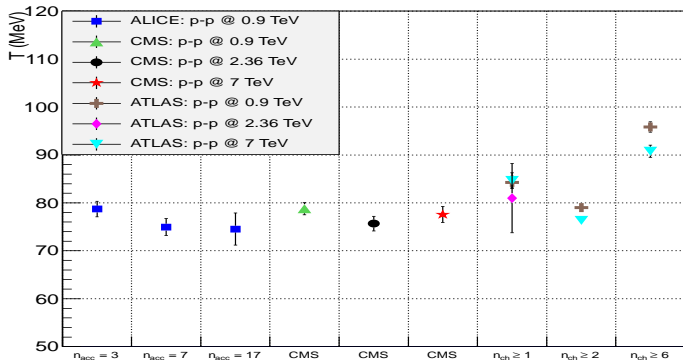
Fit parameters vs ALICE measurements

| Particles | q | T (MeV) Tsallis | T (MeV) ALICE | dN/dy Tsallis | dN/dy ALICE | χ^2/NDF Tsallis | χ^2/NDF ALICE |
|-----------------------------|------------------|--------------------|------------------|------------------|----------------|--------------------------------|------------------------------|
| K_s^0 | 1.15 ± 0.03 | 73.67 ± 3.85 | 168 ± 5 | 0.182 | 0.184 | 2.01/13 | 10.8/13 |
| Φ | 1.14 ± 0.03 | 79.99 ± 6.12 | 164 ± 91 | 0.019 | 0.021 | 0.12/1 | 0.6/1 |
| Λ | 1.11 ± 0.008 | 79.99 ± 5.63 | 229 ± 15 | 0.049 | 0.048 | 1.38/6 | 9.6/6 |
| $\Lambda(\text{bar})$ | 1.11 ± 0.008 | 70.00 ± 9.8 | 210 ± 15 | 0.047 | 0.047 | 0.42/6 | 3.7/6 |
| $\Xi^+ + \Xi(\text{bar})^+$ | 1.11 ± 0.03 | 75.00 ± 7.5 | 175 ± 50 | 0.0096 | 0.0101 | 0.189/0* | - # |

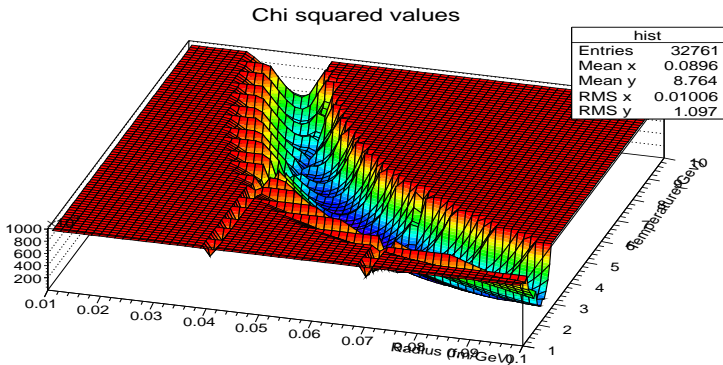
*The number of data sets is 3 and the number of fit parameters are also 3.
So, the NDF = 0



p-p collisions: Summary of results for parameter T



Tsallis: problem in determining parameters T and V



The Tsallis distribution provides an **excellent description of the transverse momentum spectra over 14 orders of magnitude**

up to 200 GeV. **Use**

$$\frac{d^2N}{dp_T dy} = gV \frac{\rho_T m_T}{(2\pi)^2} \left[1 + (q-1) \frac{m_T}{T} \right]^{-\frac{q}{q-1}},$$

Advantages : thermodynamic consistency:

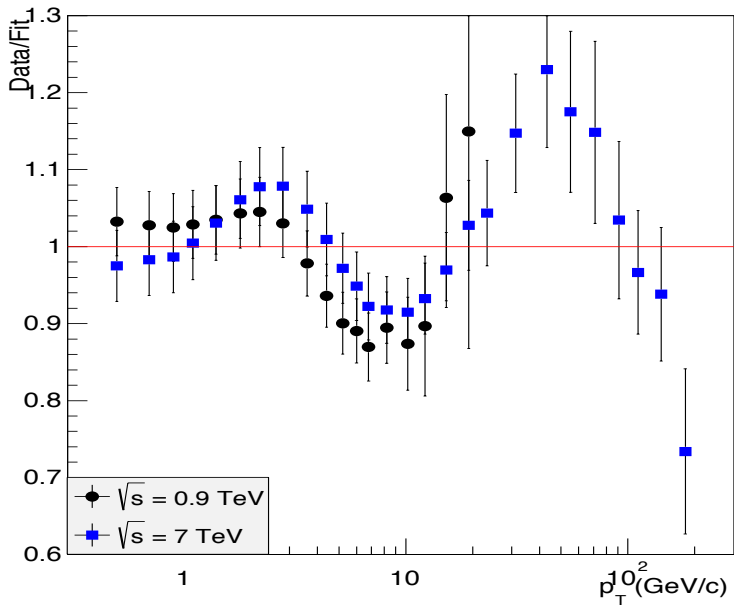
$$n = \frac{\partial P}{\partial \mu} \quad \text{etc...},$$

and the parameter T deserves its name since

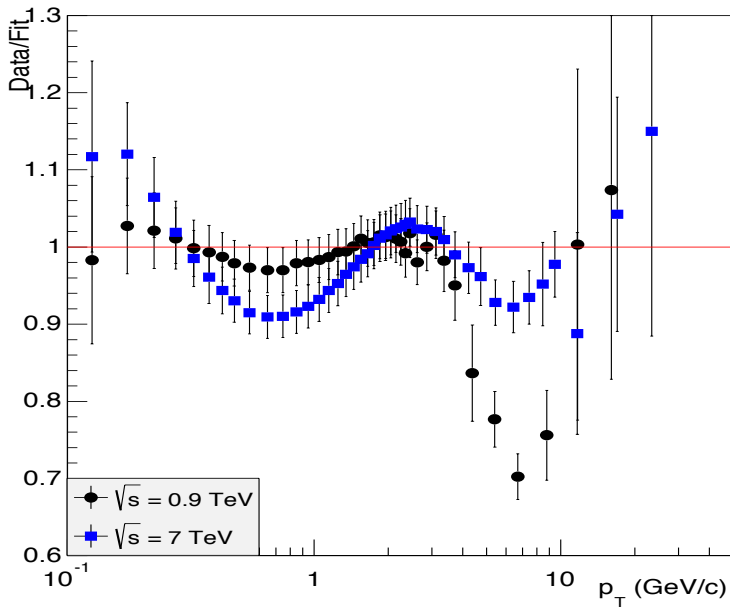
$$T = \frac{\partial S}{\partial E} \quad \dots$$



Tsallis Distribution p-p



Tsallis Distribution p-p



Tsallis Distribution p-p

| Experiment | \sqrt{s} (TeV) | R (fm) | χ^2/NDF |
|-------------------|------------------|-----------------|--------------|
| ATLAS | 3.55 | 4.62 ± 0.29 | 0.657503/36 |
| ATLAS | 3.21 | 5.05 ± 0.07 | 4.35145/41 |
| CMS | 0.9 | 4.32 ± 0.29 | 0.648806/17 |
| CMS | 7 | 5.04 ± 0.27 | 0.521746/24 |

Values of R and χ^2/NDF parameters to fit the p_T spectra measured by the ATLAS and CMS collaborations.

