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Thermal production of charms and charmonia in quark-gluon plasma

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Outline

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- Thermal production of charms
- Thermal production of charmonia
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- Summary

Introduction

Matsui and Satz(1986): J/Psi suppression as a *probe* of QGP in HIC

Hot Nuclear Matter effects:

1) **suppression in QGP and HG** (Masui, Satz et al.)

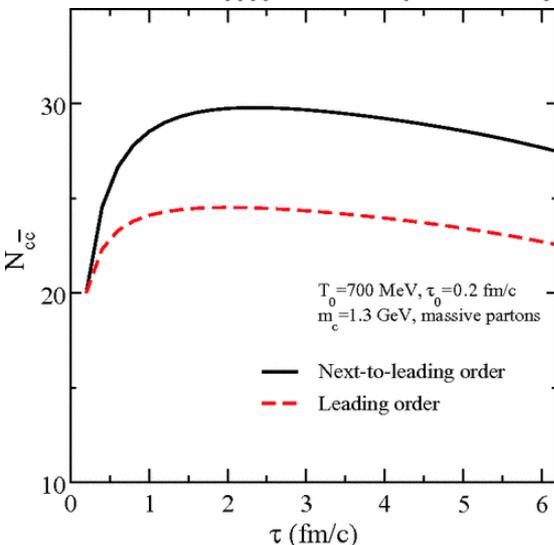
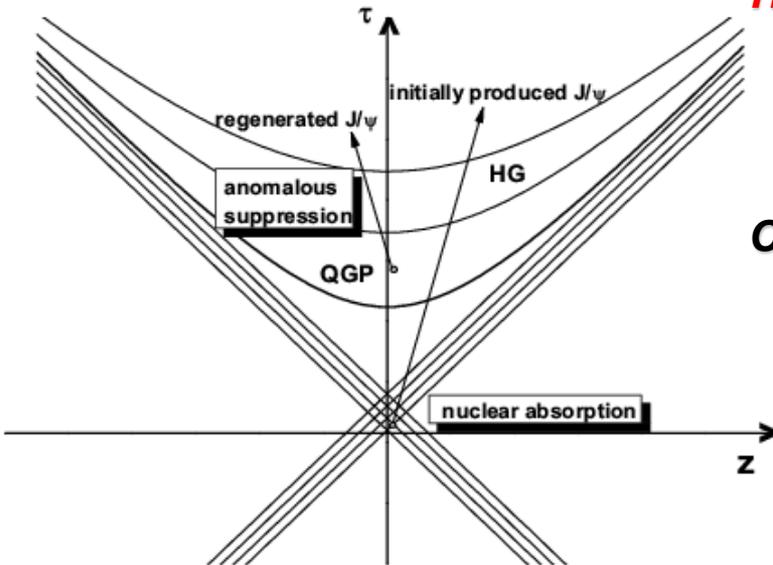
2) **regeneration in QGP and HG** (PBM, Thews, R.Rapp et al.)

Cold Nuclear Matter effects:

1) **nuclear absorption** (J.Huefner, A.Capella et al.)

2) **Cronin effect** (J.W.Cronin, J.Huefner et al.)

3) **shadowing effect** (A.H.Mueller, M.Gyulassy, R.Vogt, E.Ferreiro et al.)



All the heavy quarks are initially produced!

$$\partial_{\mu} N^{\mu} = 0$$

How about the thermal charm production and its contribution to charmonia ?

Thermal production of charms

● charm production processes

Leading Order:

$$q + \bar{q} \rightarrow c + \bar{c} \quad g + g \rightarrow c + \bar{c}$$

Next-to-leading Order:

$$q + \bar{q} \rightarrow c + \bar{c} + g \quad g + g \rightarrow c + \bar{c} + g$$

● rate equation

$$\partial_\mu N^\mu = R_{\text{gain}} - R_{\text{loss}}$$

$$R_{\text{gain}} = \frac{dN_{\text{reaction}}}{d^4x} = \frac{1}{v} \int \frac{d^3p_1}{(2\pi)^3} \frac{d^3p_2}{2E_1} \frac{d^3p_2}{(2\pi)^3} \frac{d^3p_2}{2E_2} 4F_{12} \sigma_{12} f_1 f_2$$

σ_{12} is the charm pair production cross section (P. Nason et al. 1988)

R_{loss} can be obtained by using detailed balance

● local temperature and fluid velocity from hydrodynamics

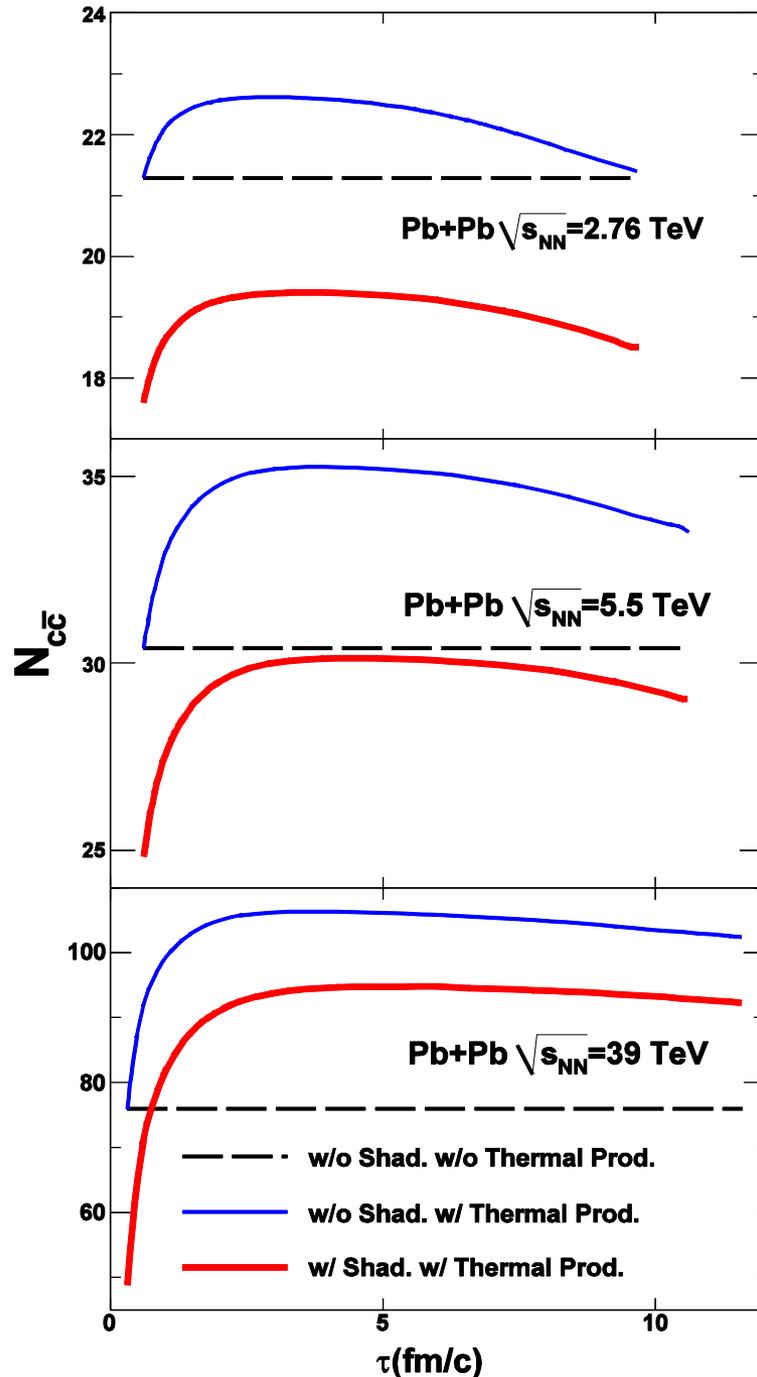
$$\partial_\mu T^{\mu\nu} = 0$$

$$\partial_\mu j^\mu = 0$$

$$T^{\mu\nu} = (\varepsilon + p)u^\mu u^\nu - g^{\mu\nu} p$$

$$j^\mu = nu^\mu$$

charm number evolution



@ 2.76 TeV, the thermal production is cancelled by shadowing.

@ 5.5 TeV, thermal charm production comparable with shadowing.

@ 39 TeV, thermal charm production overcomes shadowing.

Thermal production of charmonia

● *transport approach*

L. Yan et al. (2006), Y. Liu et al. (2009), K. Zhou et al. (2014)

● *quarkonium transport equations*

$$\partial_\tau f_\Psi + \mathbf{V}_\Psi \cdot \nabla f_\Psi = -\alpha_\Psi f_\Psi + \beta_\Psi$$

$$(\Psi = J/\psi, \chi_c, \psi')$$

$$\alpha_\Psi(\mathbf{p}_t, \mathbf{x}_t, \tau|\mathbf{b}) = \frac{1}{2E_\Psi} \int \frac{d^3\mathbf{p}_g}{(2\pi)^3 2E_g} W_{g\Psi}^{c\bar{c}}(s) f_g(\mathbf{p}_g, \mathbf{x}_t, \tau) \Theta(T(\mathbf{x}_t, \tau|\mathbf{b}) - T_c),$$

$$\beta_\Psi(\mathbf{p}_t, \mathbf{x}_t, \tau|\mathbf{b}) = \frac{1}{2E_\Psi} \int \frac{d^3\mathbf{p}_g}{(2\pi)^3 2E_g} \frac{d^3\mathbf{p}_c}{(2\pi)^3 2E_c} \frac{d^3\mathbf{p}_{\bar{c}}}{(2\pi)^3 2E_{\bar{c}}} W_{c\bar{c}}^{g\Psi}(s) \underline{f_c(\mathbf{p}_c, \mathbf{x}_t, \tau|\mathbf{b}) f_{\bar{c}}(\mathbf{p}_{\bar{c}}, \mathbf{x}_t, \tau|\mathbf{b})} \\ \times (2\pi)^4 \delta^{(4)}(p + p_g - p_c - p_{\bar{c}}) \Theta(T(\mathbf{x}_t, \tau|\mathbf{b}) - T_c),$$

***Including thermal
charm production***

Numerical results

@2.76TeV

Thermal production increase R_{AA} , 0.48--
→0.55(with shadowing);

Thermal production increase R_{AA} , 0.48--
→0.77(without shadowing).

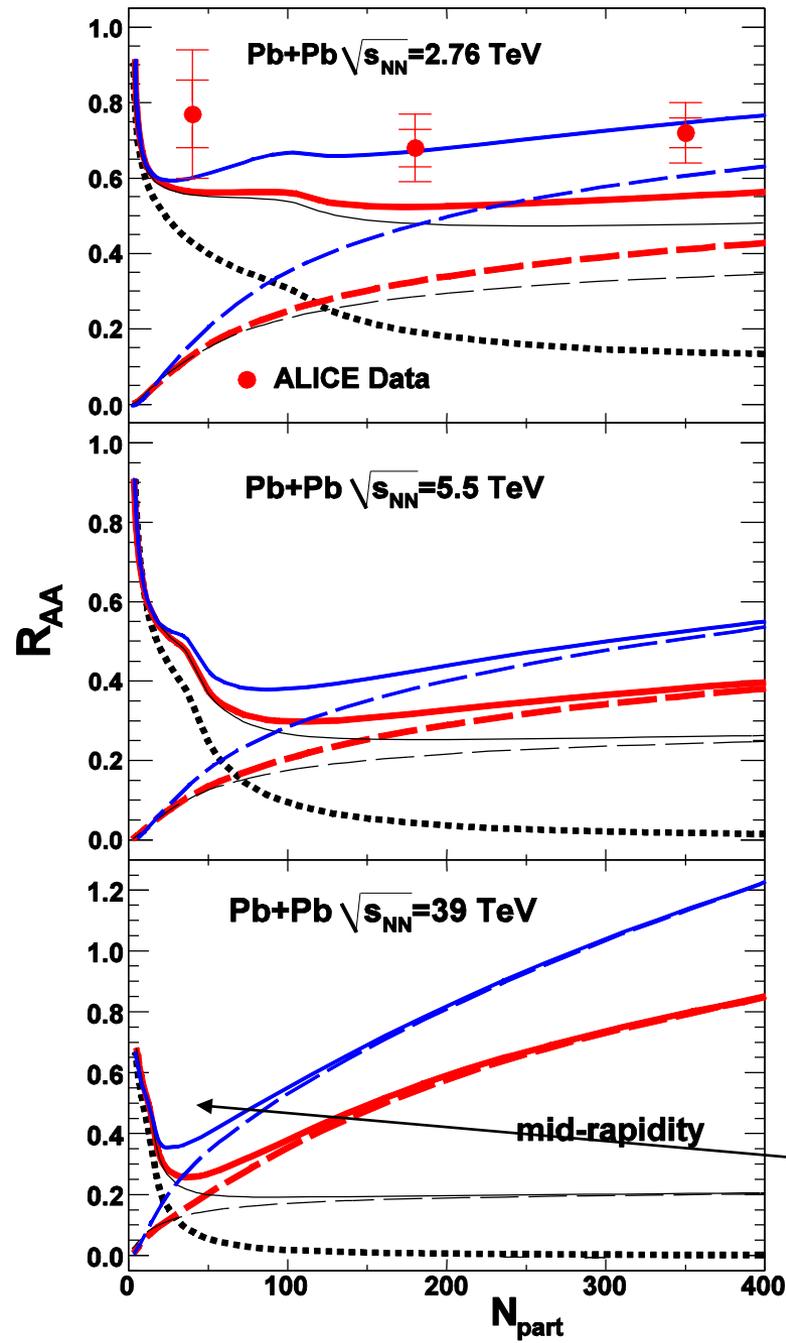
@5.5TeV

With thermal production, R_{AA} appears a flat structure;
Without thermal production, R_{AA} , shows a slightly increasing trend.

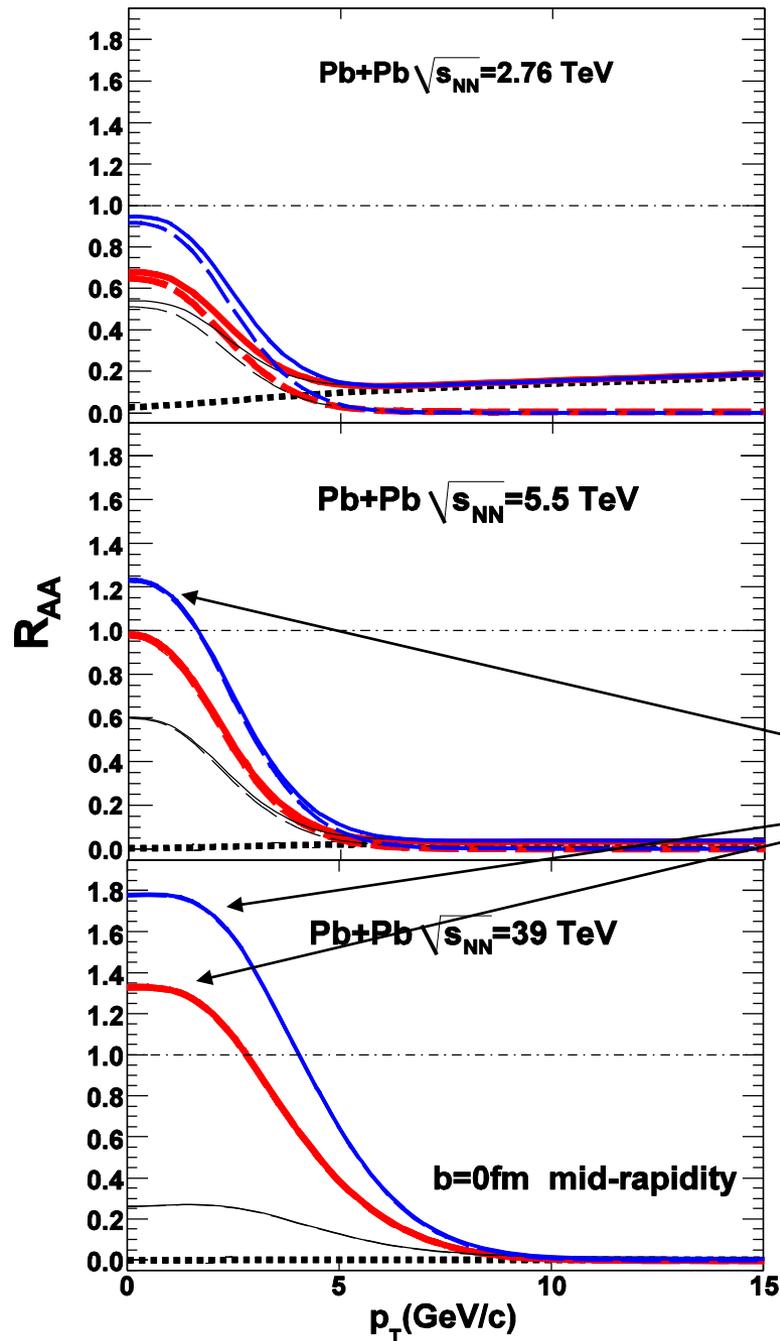
@39TeV

Thermal production increase R_{AA} , 0.2--
→0.85(with shadowing);

Thermal production increase R_{AA} , 0.2--
→1.25(without shadowing).



valley structure



@LHC

Thermal charm production enhances the charmonium regeneration further, which strengthens the decreasing trend.

@FCC

The QGP becomes a rich source for charm pair thermal production, and this overcome the stronger suppression through regeneration.

exceed unit

Summary

- *At FCC , the thermal charm production dramatically enhances the charm quark pairs yield.*
- *At FCC , RAA presents a clearly valley structure due to the competition between the strong suppression and strong regeneration with centrality.*
- *At FCC , differential RAA as a function of p_T can even exceed 1 with thermal charm production in central collisions.*

Thank You!