

SCALING TRENDS IN P+P COLLISIONS FROM SPS TO LHC

General Characteristics of pp-collisions at LHC

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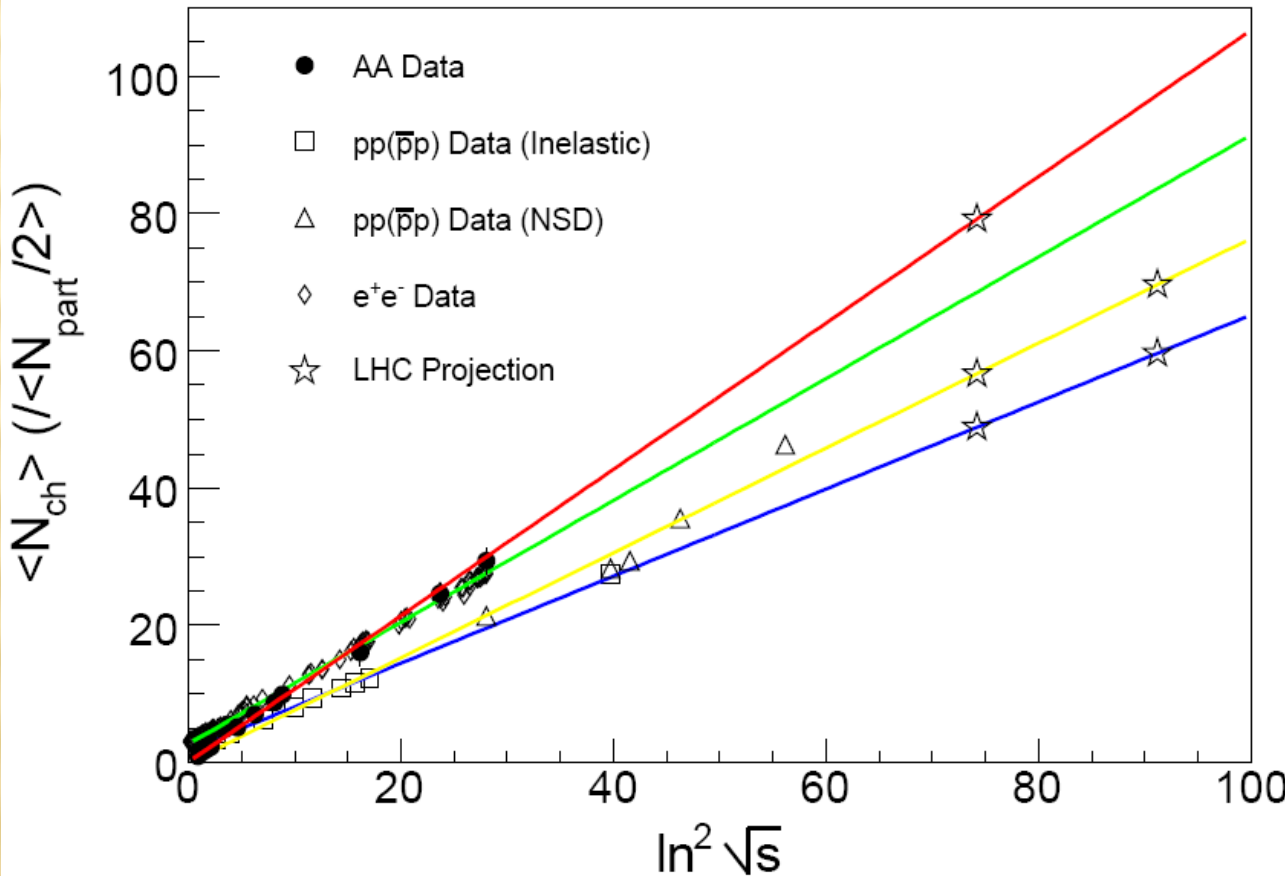
in collaboration with

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R. Kolevator, E. Zabrodin

MOTIVATION: EXPERIMENTAL RESULTS

W. Busza, JPG 35 (2008) 044040



Predictions for LHC

inelastic pp :

$N_{ch} = 60 \pm 10$ (14 TeV)

$N_{ch} = 49 \pm 8$ (5.5 TeV)

NSD pp :

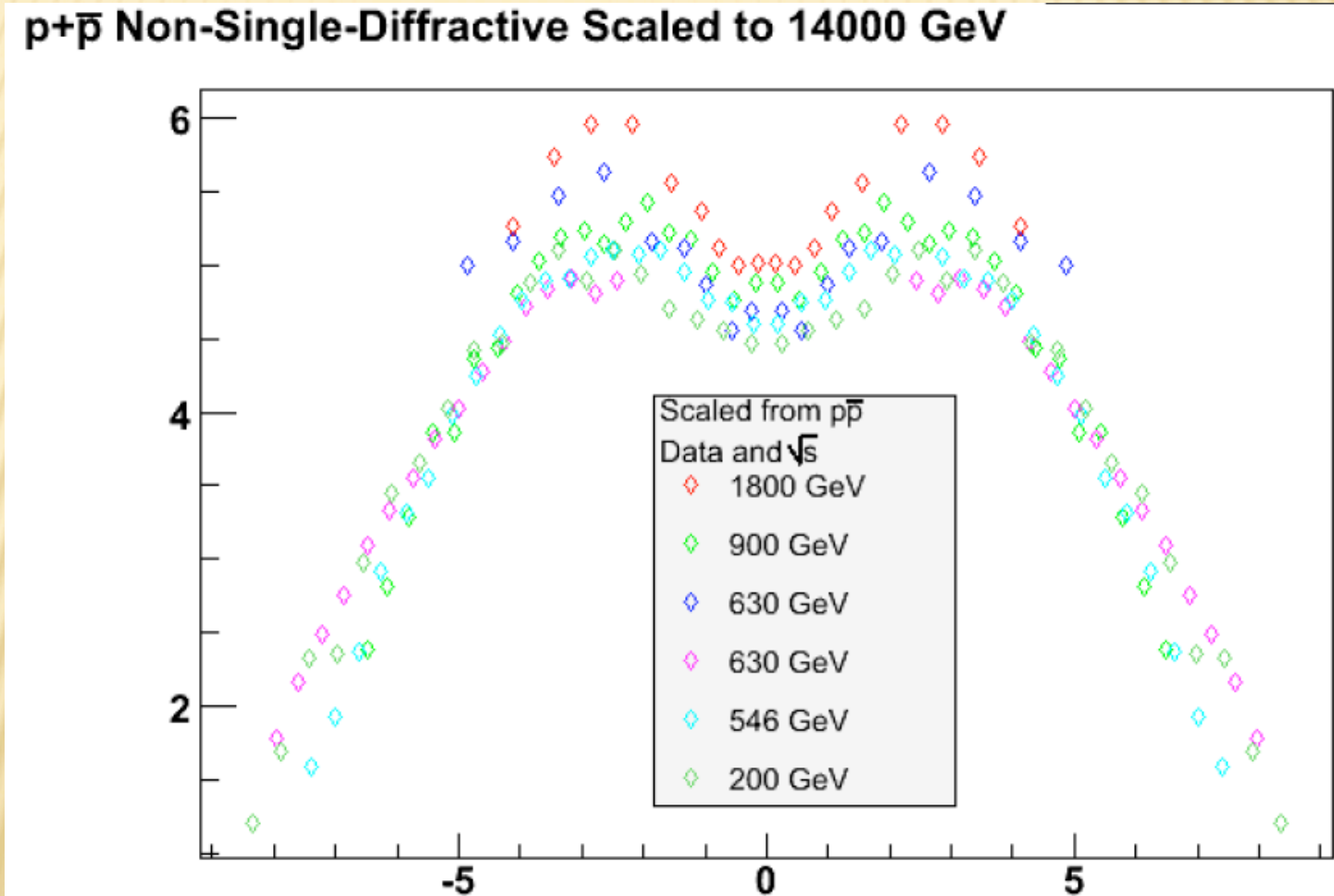
$N_{ch} = 70 \pm 8$ (14 TeV)

$N_{ch} = 57 \pm 7$ (5.5 TeV)

Energy dependence of particle multiplicities

MOTIVATION: EXPERIMENTAL RESULTS

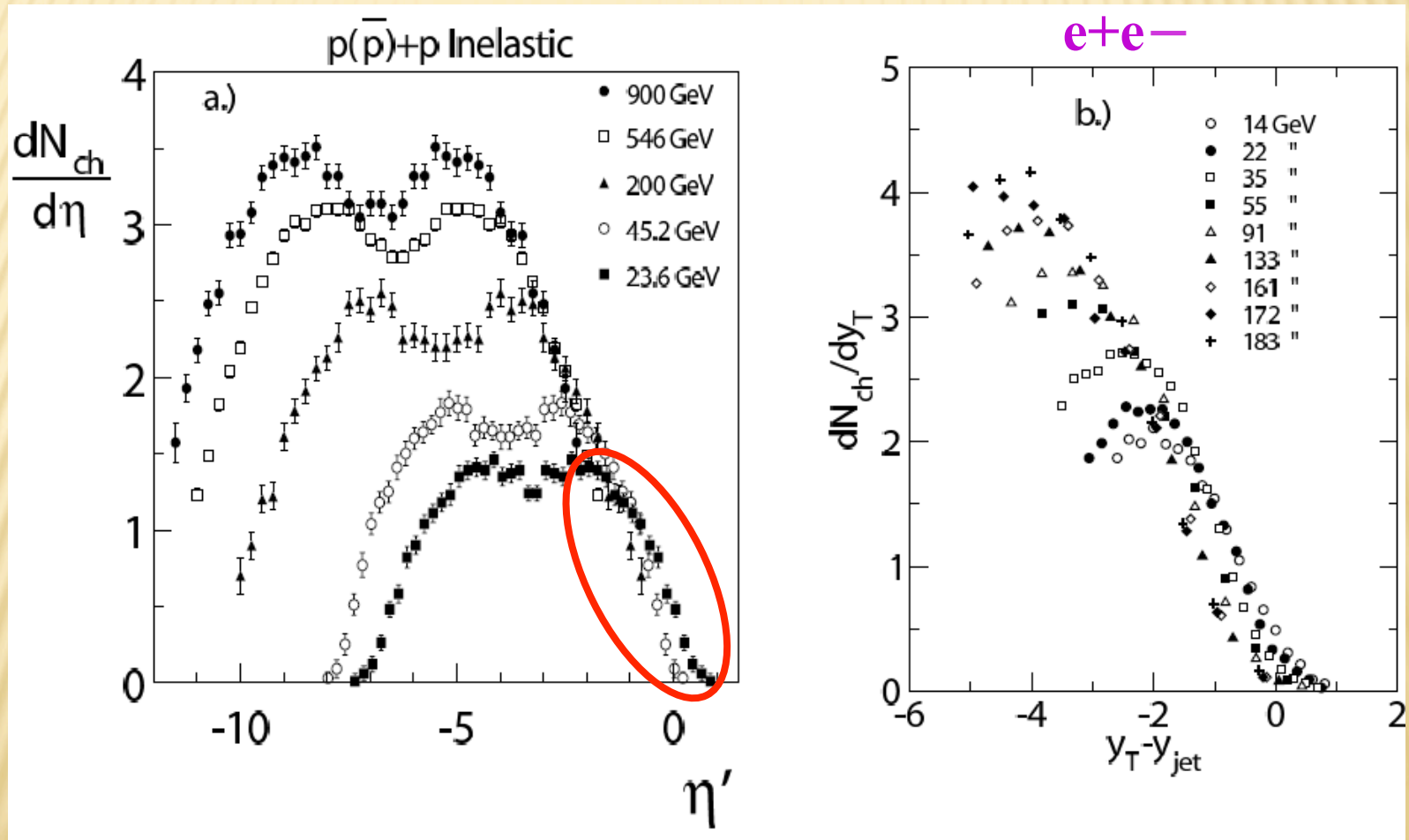
W. Busza, JPG 35 (2008) 044040



Extrapolation of NSD pp data to LHC using $\ln\sqrt{s}$ scaling of the width and height of the distribution

MOTIVATION: EXPERIMENTAL RESULTS

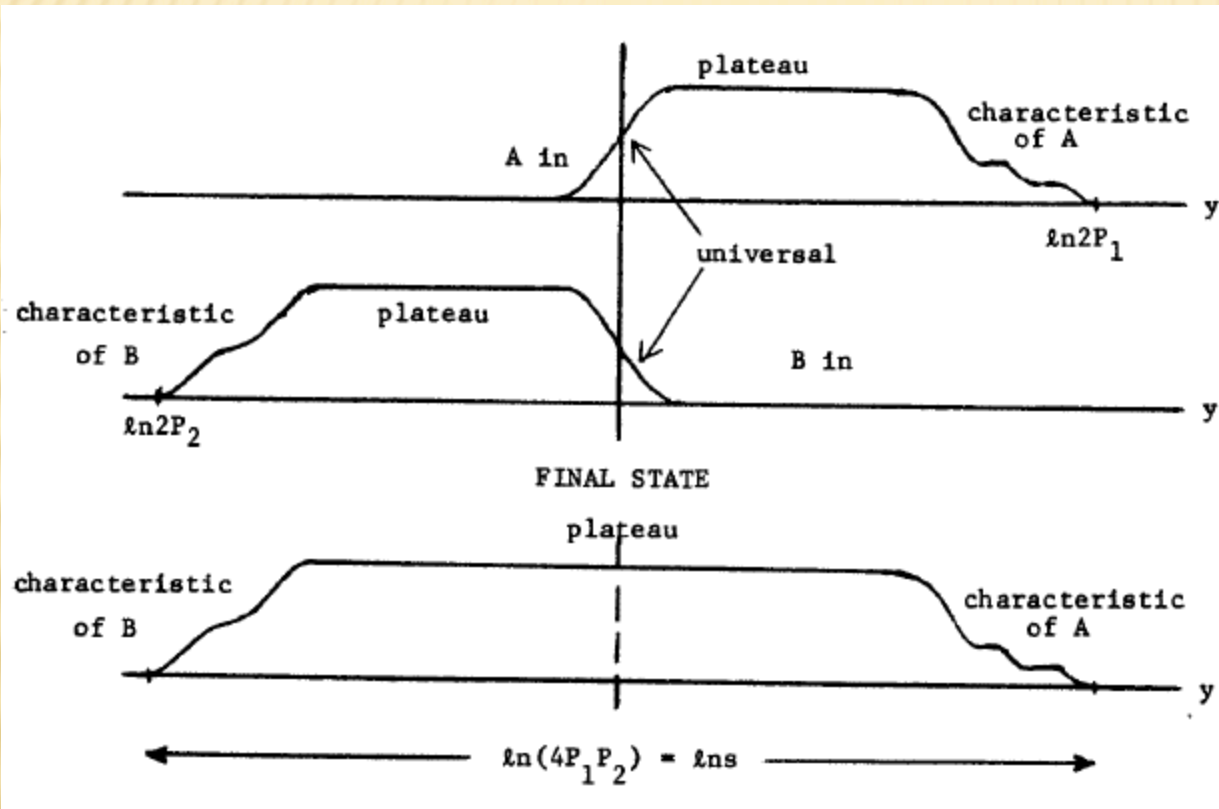
W. Busza, JPG 35 (2008) 044040



Example of extended longitudinal scaling in different reactions

HYPOTHESIS OF FEYNMAN SCALING

R. Feynman, PRL 23 (1969) 1415; also in "Photon-hadron interactions"



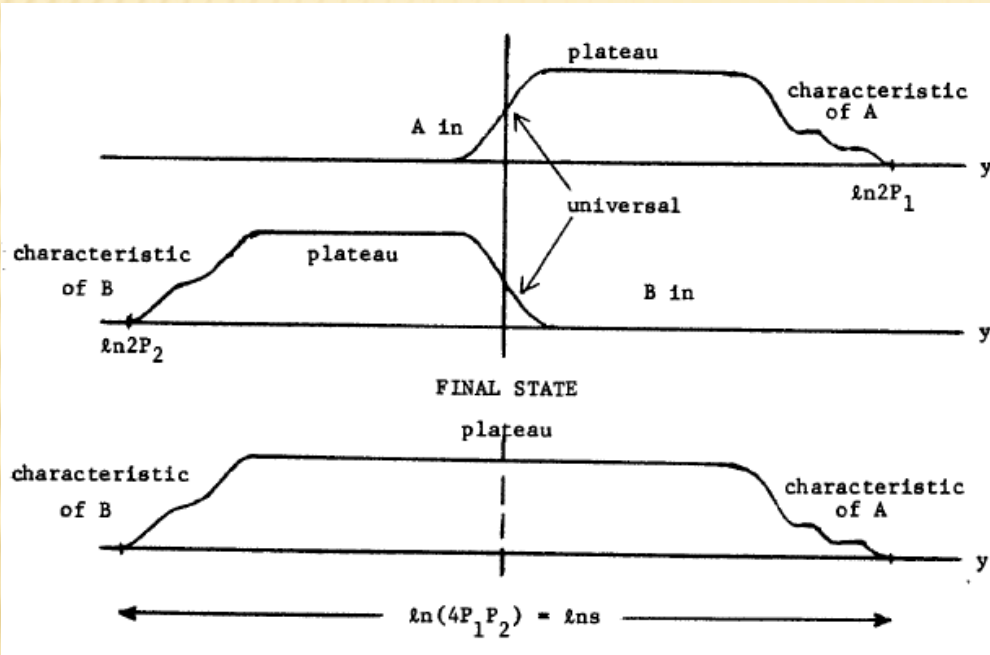
Basic assumption:
scaling of inclusive spectra within the whole kinematically allowed region of x_F (or c.m. y)

In addition:
existence of central area $-x_0 \leq x_F \leq x_0$, where $x_0 \approx (0.1 - 0.2)$ is assumed.

In terms of rapidity

$$-\ln[x_0 \sqrt{s} / m_T] \leq y^* \leq \ln[x_0 \sqrt{s} / m_T]$$

CONSEQUENCES OF FEYNMAN SCALING



- (1) Logarithmic rise of the central rapidity region with energy

$$(\Delta y^*) \approx 2 \ln(x_0 \sqrt{s} / m_T)$$

- (2) Fragmentation regions are fixed

$$(\Delta y^*) \approx \ln(1 / x_0)$$

- (3) Main contribution to mean multiplicity comes from the central area

$$\langle n \rangle : \ln(x_0 \sqrt{s} / m_T)$$

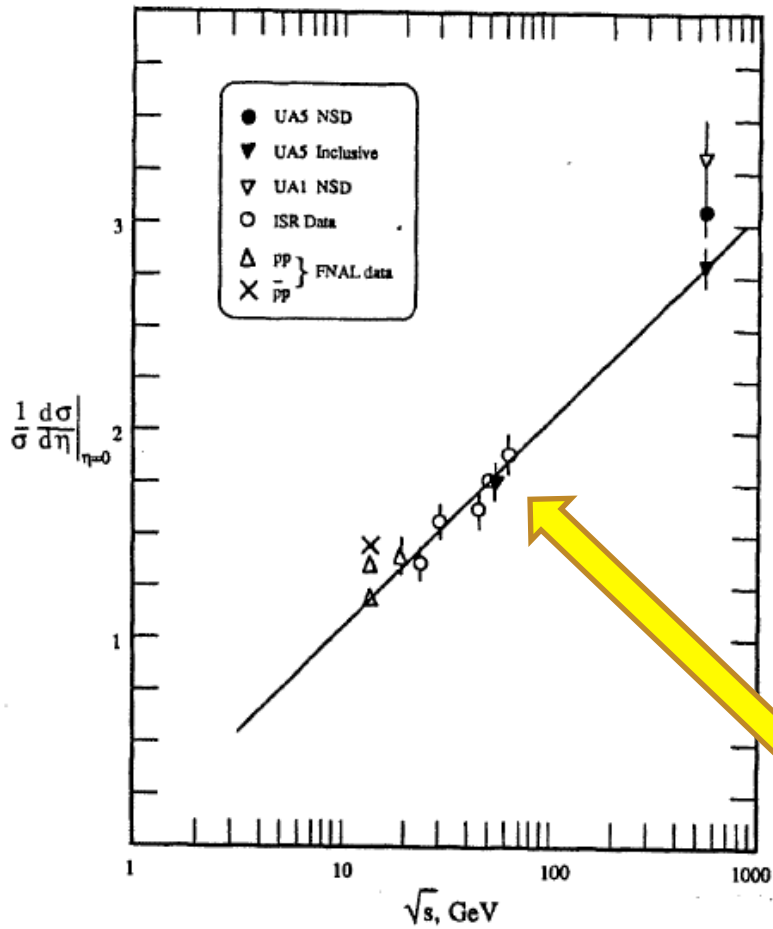
- (4) In the central area particle density does not depend on energy and rapidity

$$\rho(y^*, p_T; \sqrt{s}) = \rho(p_T)$$

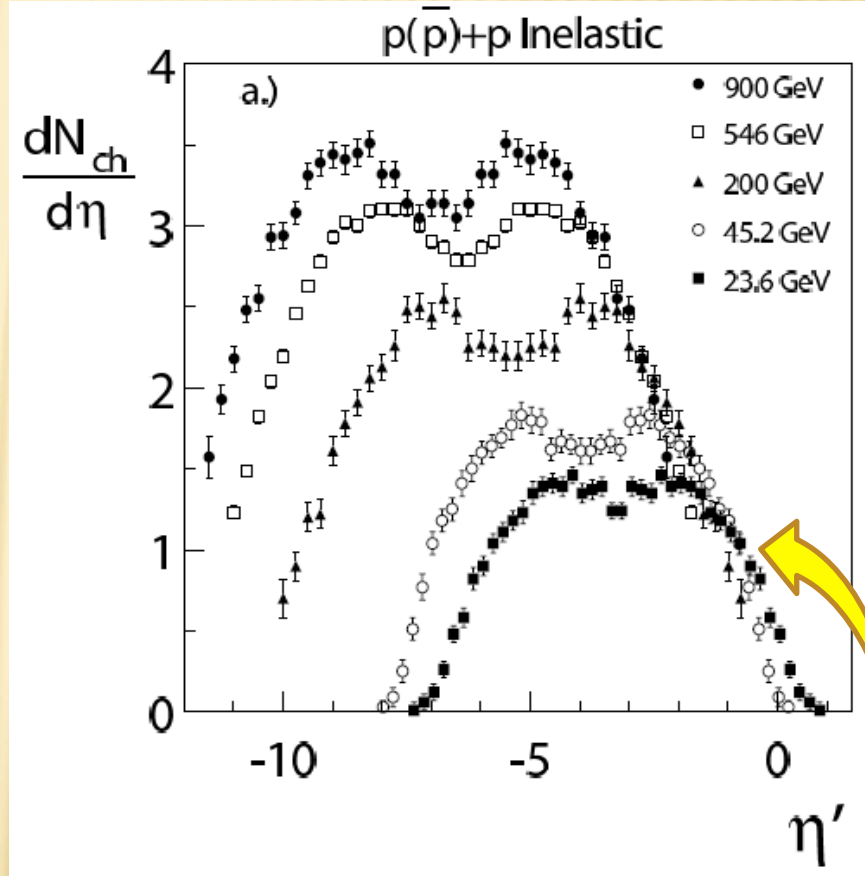
- (5) Contribution from the fragmentation regions is energy independent

VIOLATION OF FEYNMAN SCALING

UA5 Collab., Phys. Rep. 154 (1987) 247



W. Busza, JPG 35 (2008) 044040



Charged particle pseudorapidity density at $\eta = 0$ as a function of \sqrt{s}

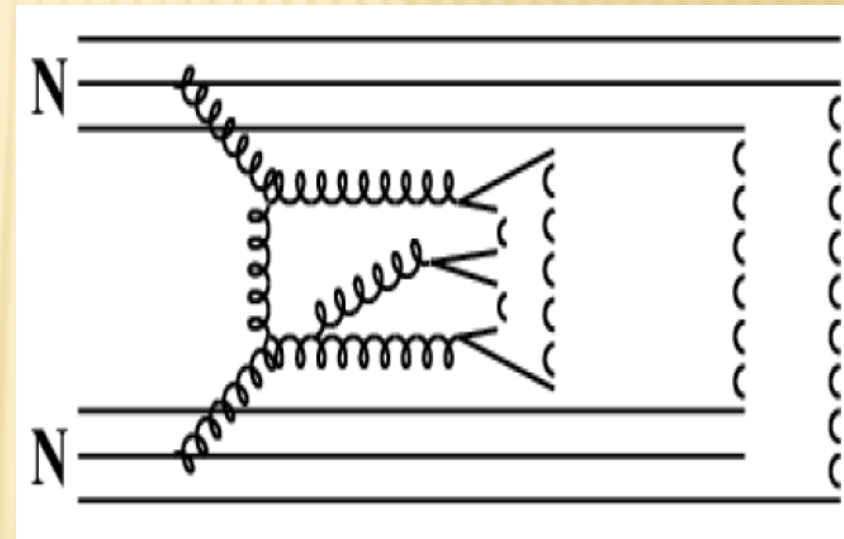
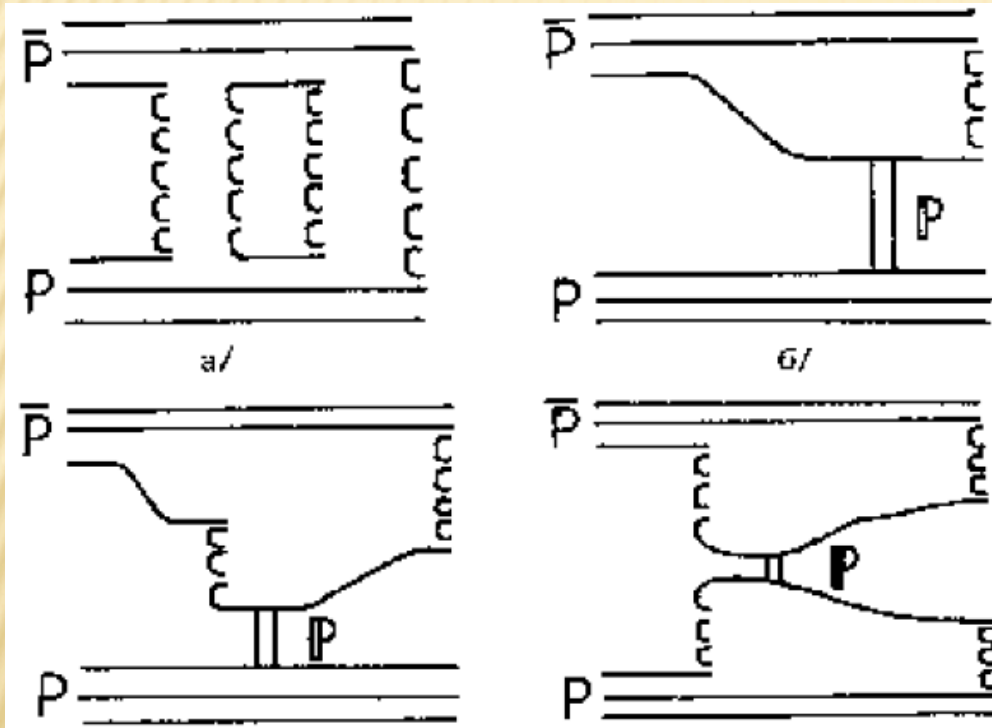
Violation of Feynman scaling, but ext. long. scaling holds?!

QUARK-GLUON STRING MODEL

A.B. Kaidalov, K.A. Ter-Martirosyan, PLB 117 (1982)

N.S. Amelin, L.V. Bravina, Sov. J. Nucl. Phys. 51 (1990) 133

N.S. Amelin, E.F. Staubo, L.P. Csernai, PRD 46 (1992) 4873



At ultra-relativistic energies: multi-Pomeron scattering, single and double diffraction, and jets (hard Pomeron exchange)

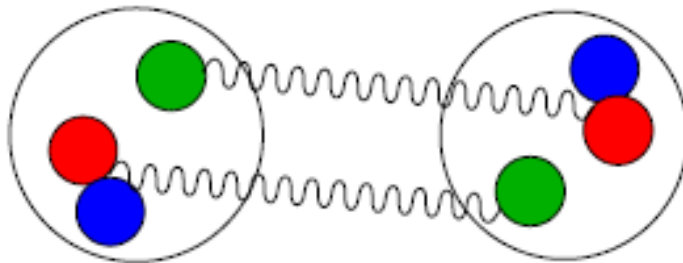
Gribov's Reggeon Calculus + string phenomenology

QUARK-GLUON STRING MODEL

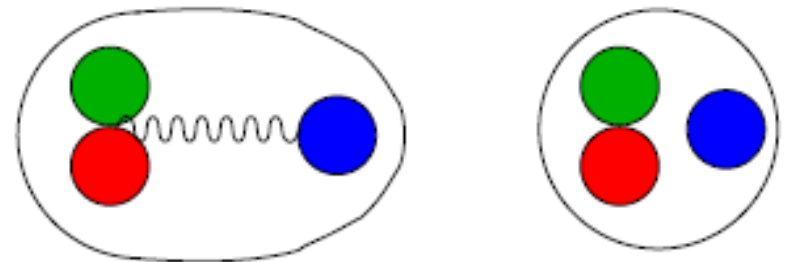
two different mechanisms:

- ▶ excitation due to exchange of pomerons (color exchange)
- ▶ transverse strings

- ▶ excitation due to transfer of momentum to a single parton
- ▶ longitudinal string



- ▶ n cut pomerons give $2n$ strings



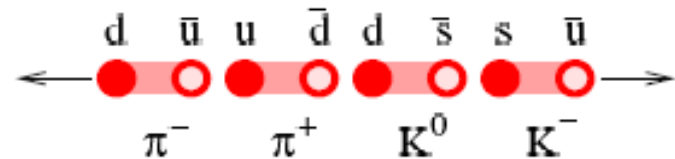
- ▶ purely phenomenological process

Excitation of color neutral strings

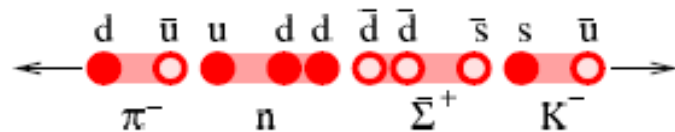
QUARK-GLUON STRING MODEL

Decay of strings - production of mesons and baryons:

- ▶ the colorfield between a quark and a antiquark gets “stretched”
- ▶ a meson (baryon) with some transverse momentum is formed and gets a fraction z of the primordial momentum of the string
- ▶ z is generated from the fragmentation function
- ▶ the rest of the string either decays further or forms a cluster



production of mesons



production of baryons

Decay of strings and particle production

Predictions for LHC.

7. $\langle n_{ch} \rangle$ $80 \div 100$

8. $\left. \frac{dn_{sb}}{dy} \right|_{y=0}$ $5.5 \div 6.0$

9. Structures in σ_n

10. Strong long-range (in y) correlations

11. Large amount of minijets.

7. QGSM: Predictions for LHC.

1. $\sigma^{(tot)}$ 103 mb ($\sigma^{(tot)}_{(s)} \sim \ln^2 \frac{s}{s_0}$)

2. $\sigma^{(el)}$ 26 mb ($\sigma^{(el)}_{(s)} \sim \ln^2 \frac{s}{s_0}$)

3. $B(0)$ 21.5 GeV⁻² ($B(0) \sim \ln^2 \frac{s}{s_0}$)

4. $\rho = \frac{\text{Re}T(0)}{\text{Im}T(0)}$ 0.11

5. σ_{SD} 12 ÷ 13 mb ($\sigma_{SD} \sim \sigma_{DD} \sim \ln \frac{s}{s_0}$)

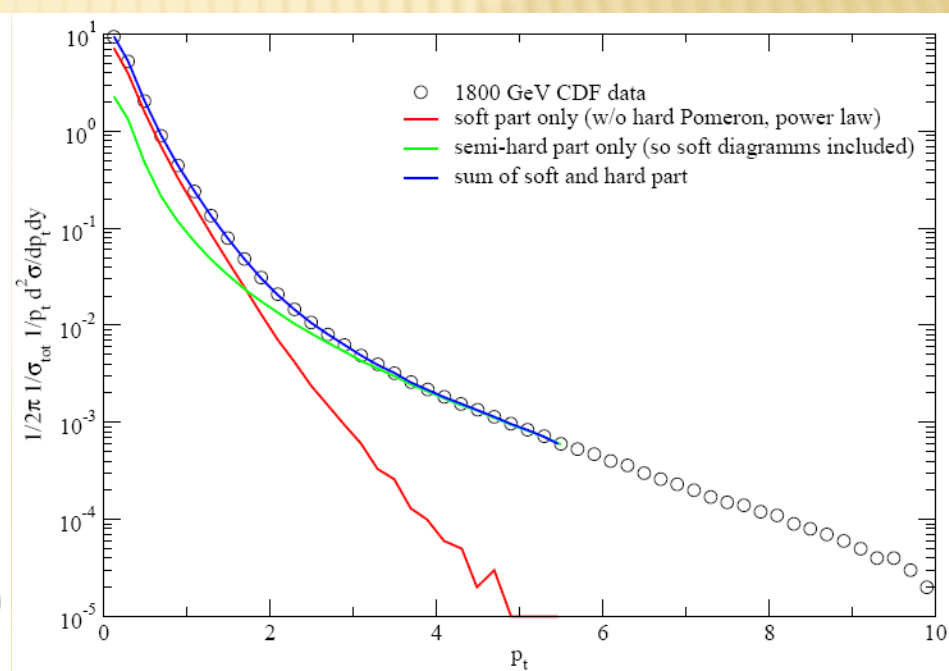
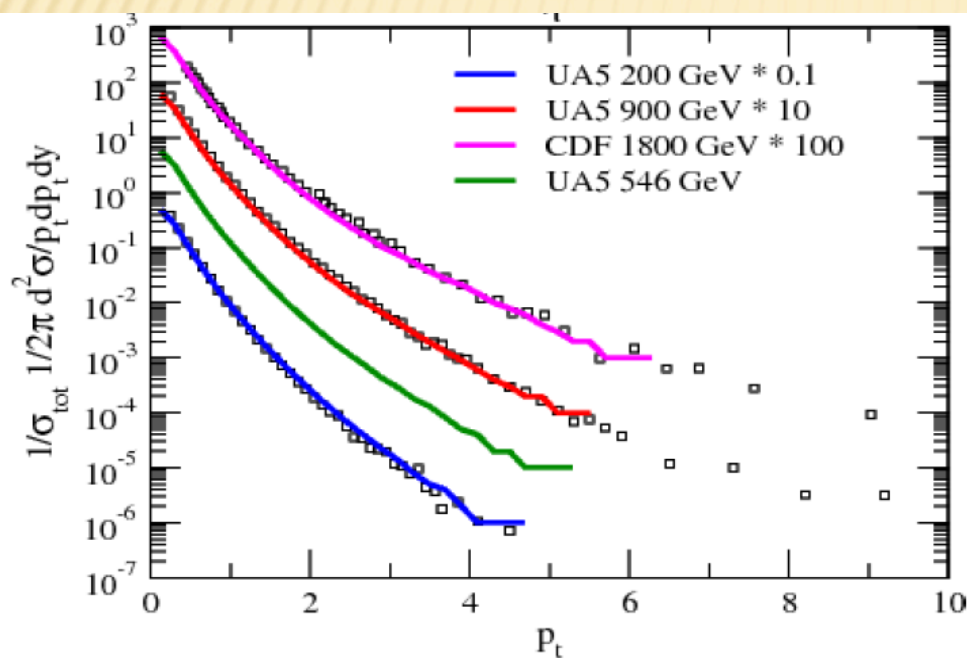
6. σ_{DD} 11 ÷ 13 mb

$$\sigma^{(el)} + \sigma_{SD} + \sigma_{DD} = 51 \text{ mb} \approx \frac{1}{2} \sigma^{(tot)}$$

RAPIDITY AND P_T SPECTRA: MODEL VS. DATA

Transverse momentum spectra

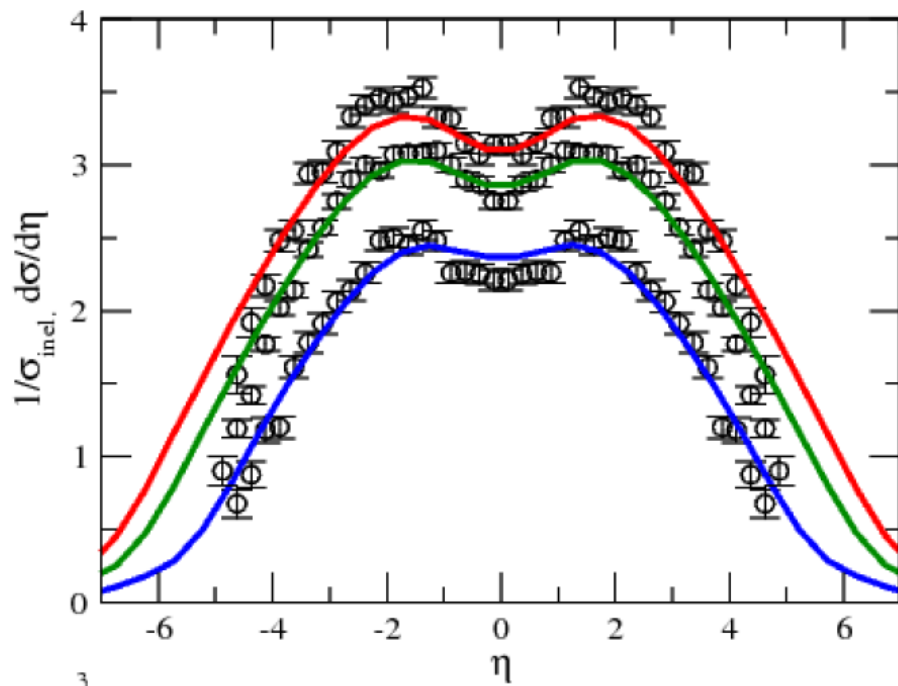
Hard and soft components



Description of P_t distributions seems to be good

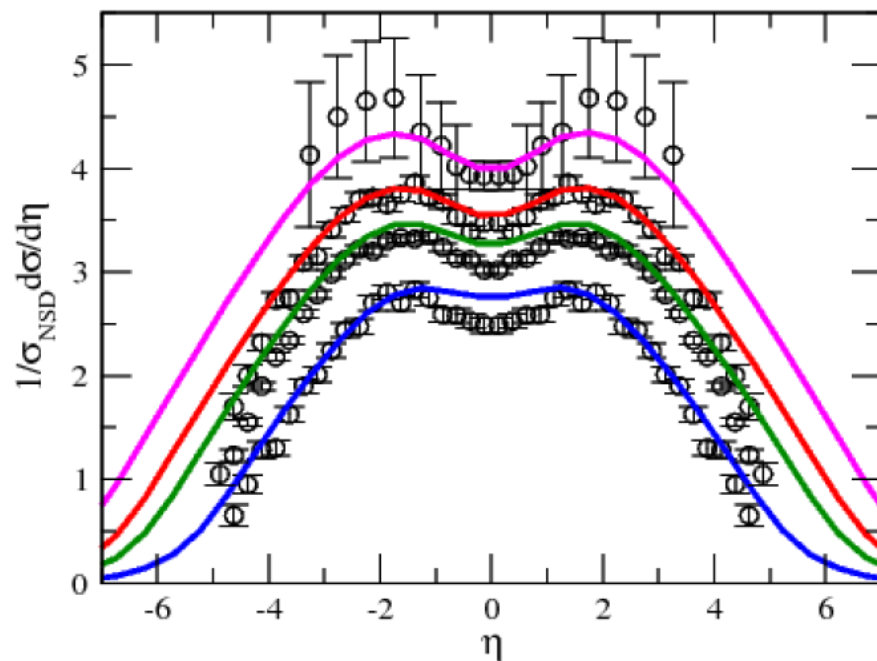
RAPIDITY AND P_T SPECTRA: MODEL VS. DATA

Inelastic collisions



200, 546, 900 GeV

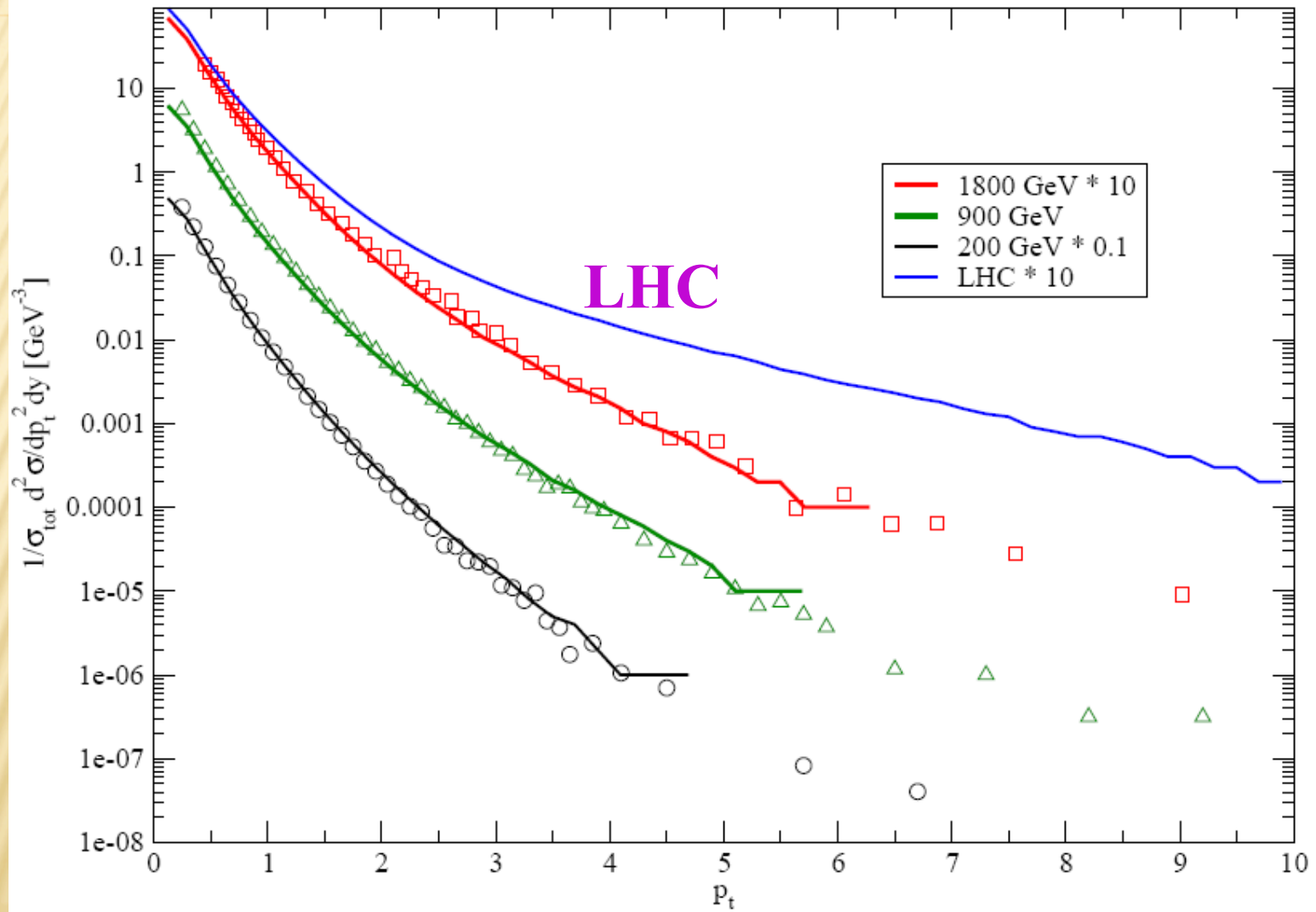
NSD collisions



200, 546, 900, 1800 GeV

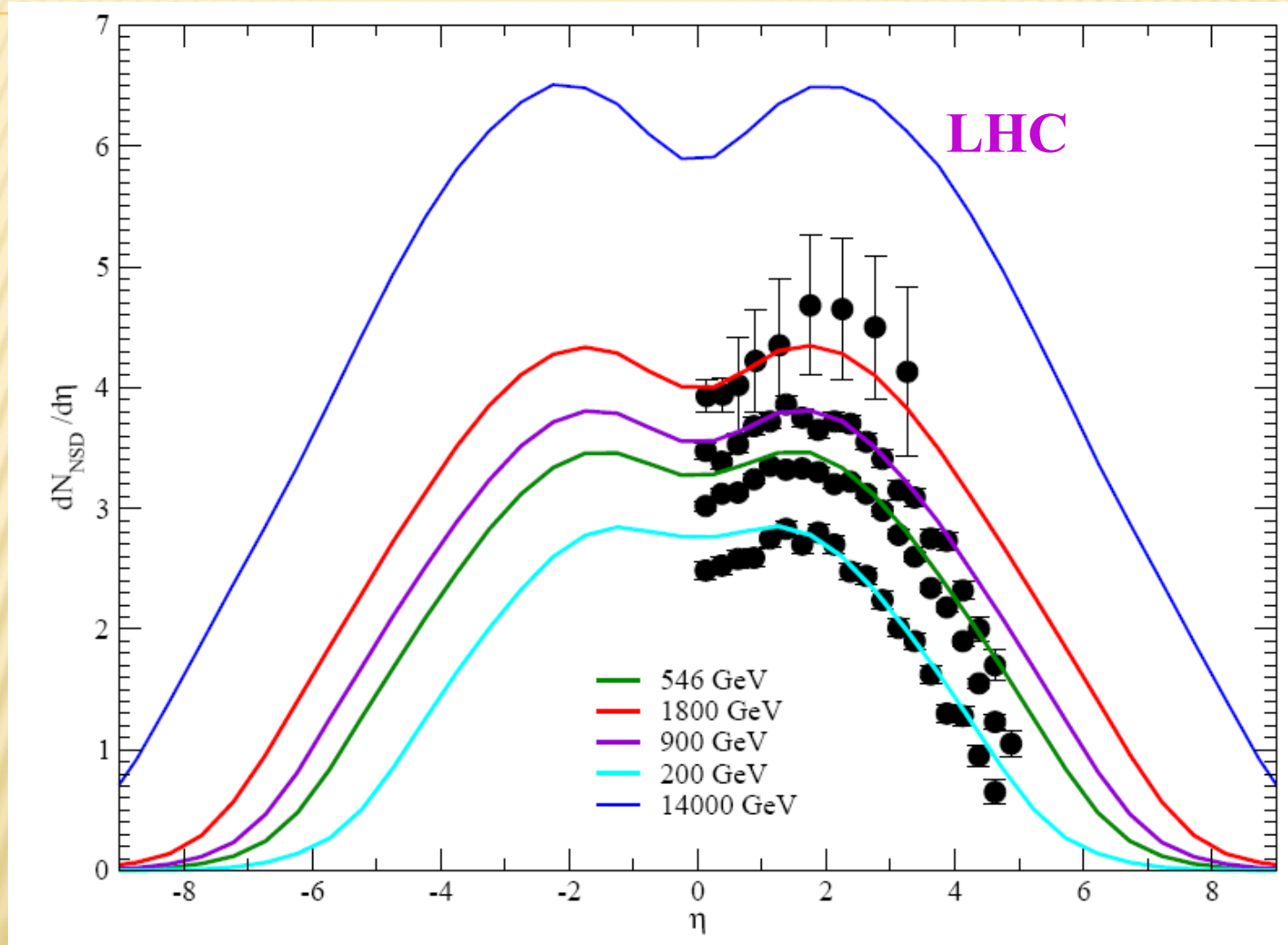
Description of pseudorapidity spectra seems to be good

PREDICTIONS FOR P+P @ LHC



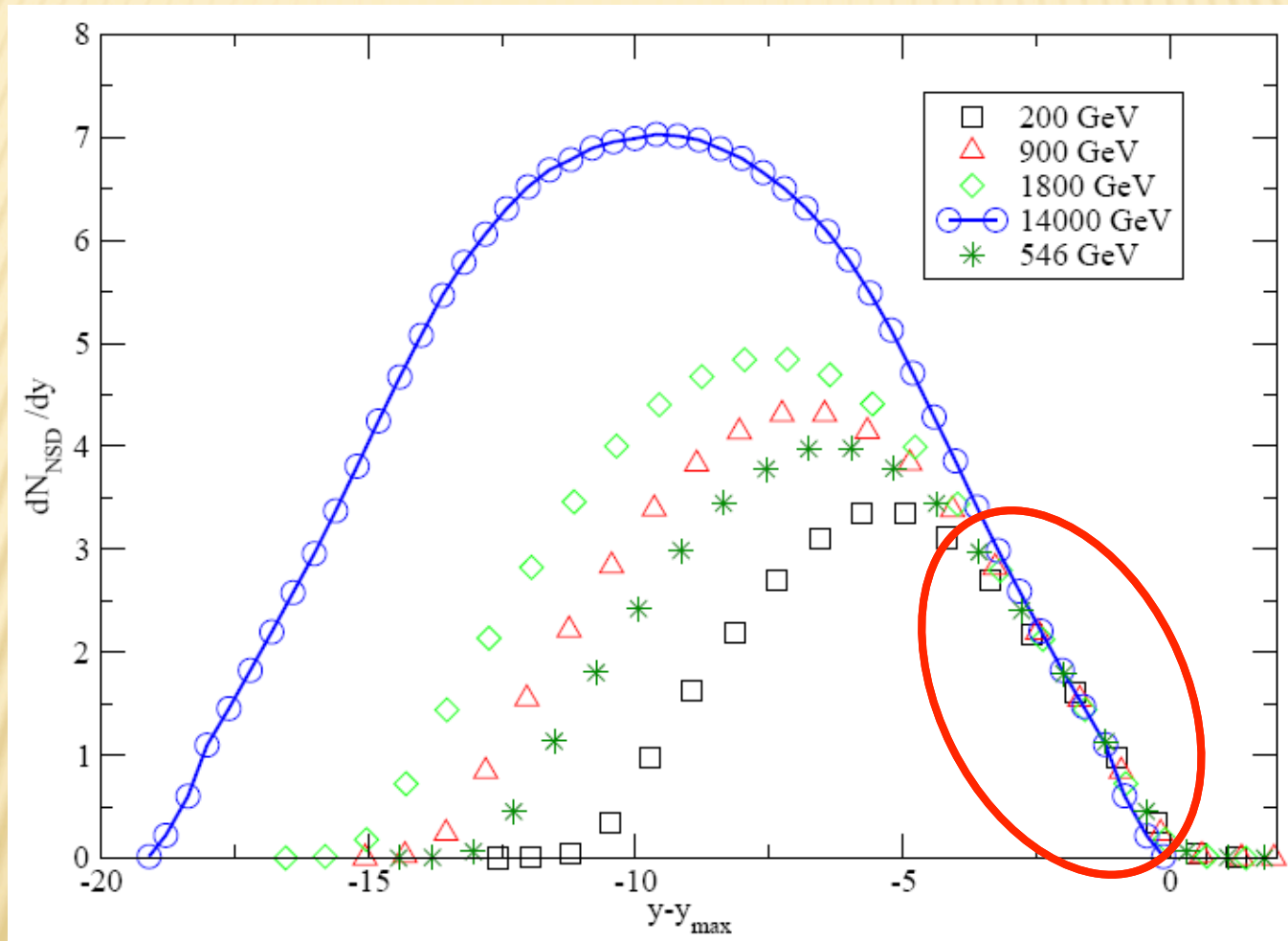
QGSM: transverse momentum distribution of particles

PREDICTIONS FOR P+P @ LHC



QGSM: pseudorapidity distribution of particles

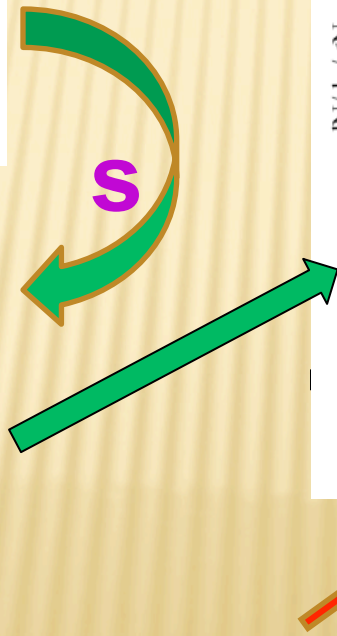
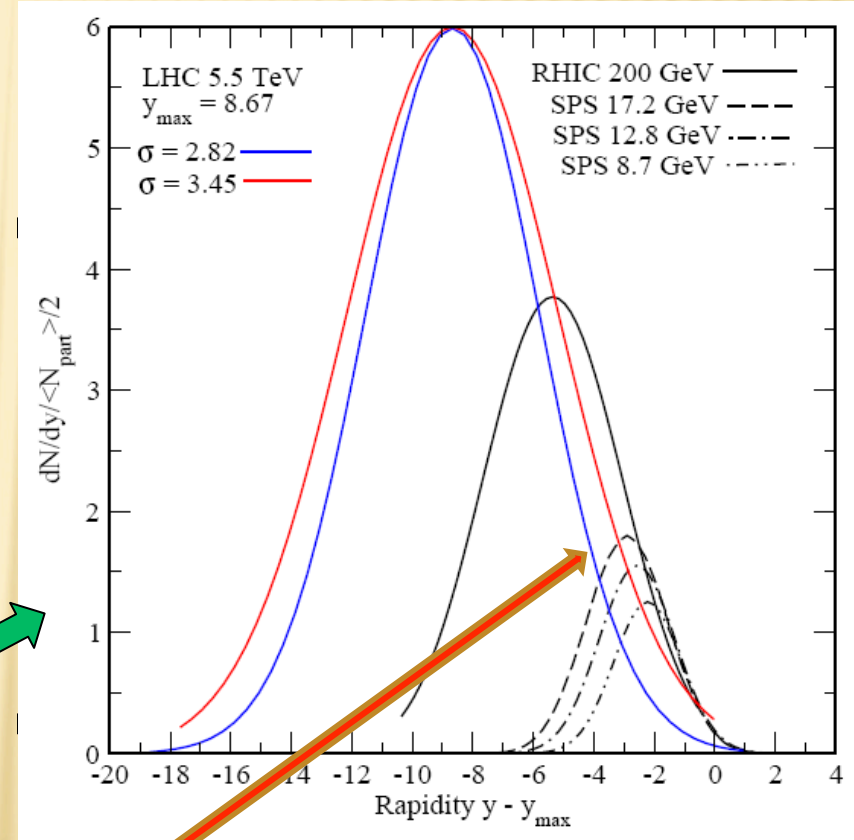
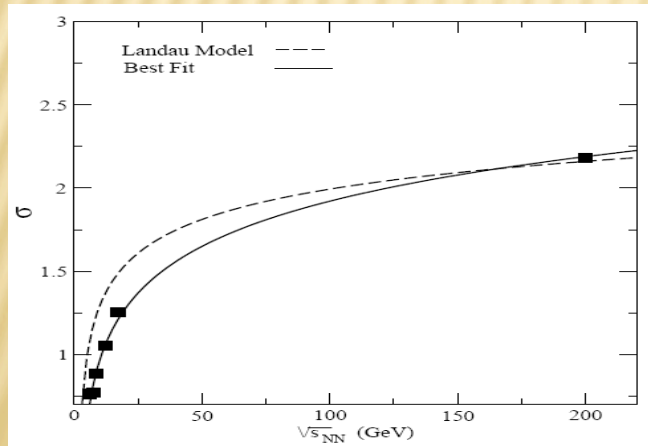
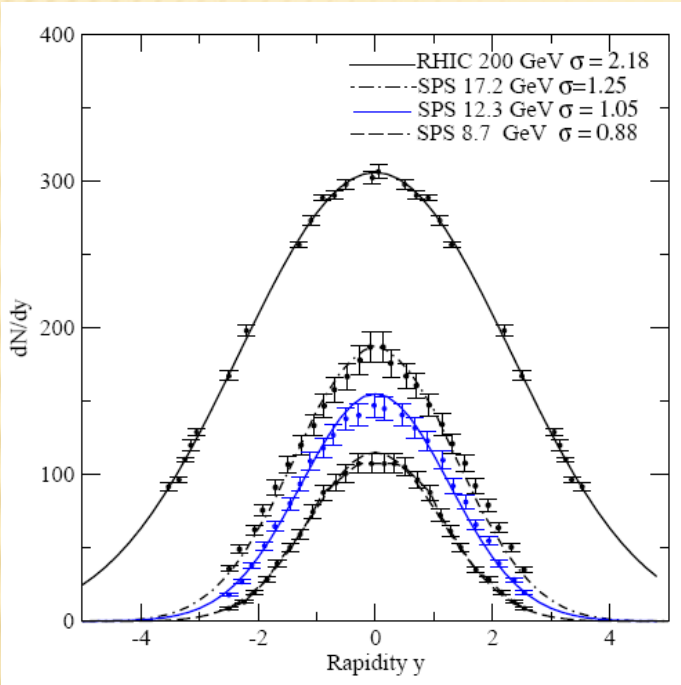
PREDICTIONS FOR P+P @ LHC



QGSM: extended longitudinal scaling in p+p collisions holds

VIOLATION OF ELS IN A+A AT LHC?

J. Cleymans, J.Struempfer, L.Turko, PRC 78 (2008) 017901



Statistical thermal model: ELS will be violated in A+A @ LHC. What about p+p ?

WHY SCALING HOLDS IN THE MODEL?

Correlation function

$$C(y_i, y_j) \propto \exp\{-\lambda(y_i - y_j)\}$$

Particles are uncorrelated if

$$y_i - y_j \equiv \Delta y \gg 1$$

Consider now inclusive process

$$1 + 2 \rightarrow i + X$$

Particle inclusive cross section

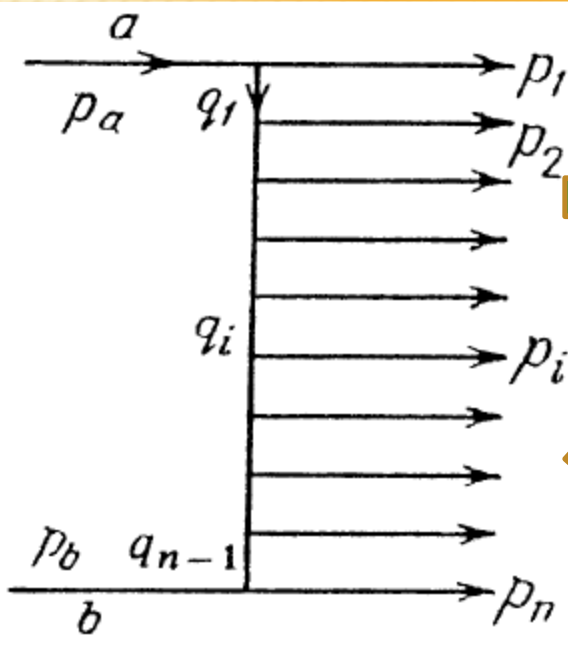
$$f_i = \frac{d^2 \sigma(y_1 - y_i, y_i - y_2, p_{iT}^2)}{dy_i d^2 p_{iT}}$$

In the fragmentation region of particle 1

$$y_1 - y_i \approx 1, y_i - y_2 \approx y_1 - y_2 \gg 1$$

Inclusive density

$$n_i = f_i / \sigma_{inel} = \phi(y_1 - y_i, p_{iT}^2)$$



Short range correlations

$$x_F^{(i)} \equiv \frac{p_{iP}}{p_P^{\max}} \approx \exp\{-(y_1 - y_i)\}$$

therefore

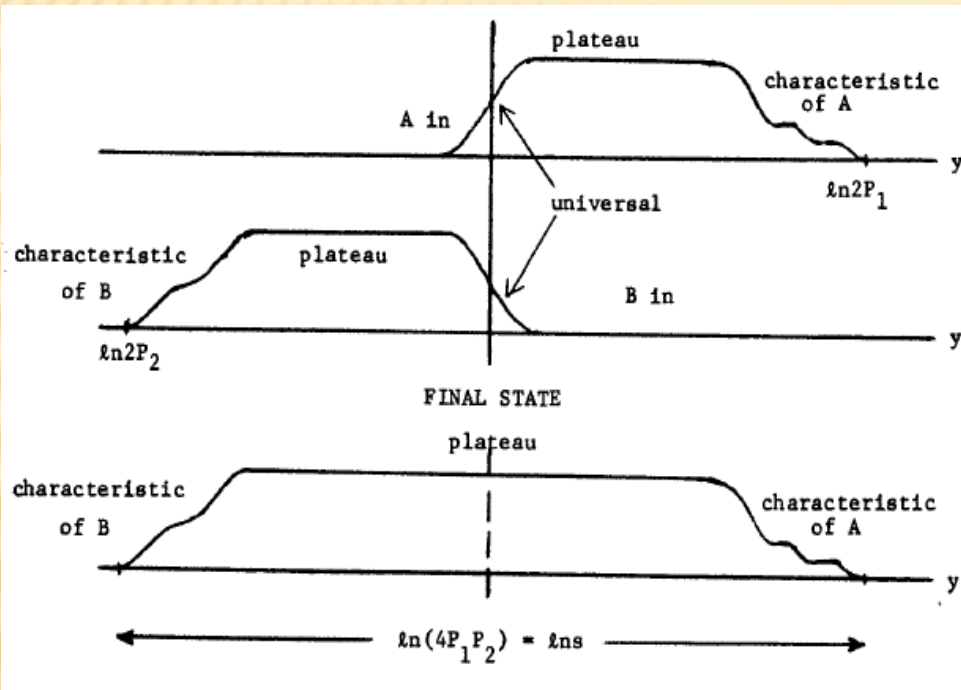
$$n_i = \psi(x_F^{(i)}, p_{iT}^2)$$



In string models both FS and ELS holds in the fragmentation regions

KOBA-NIELSEN-OLESEN (KNO) SCALING

Z.Koba, H.B.Nielsen, P.Olesen, NPB 40 (1972) 317



They claim that if Feynman scaling holds, then the **multiplicity distribution is independent of energy** except through the variable

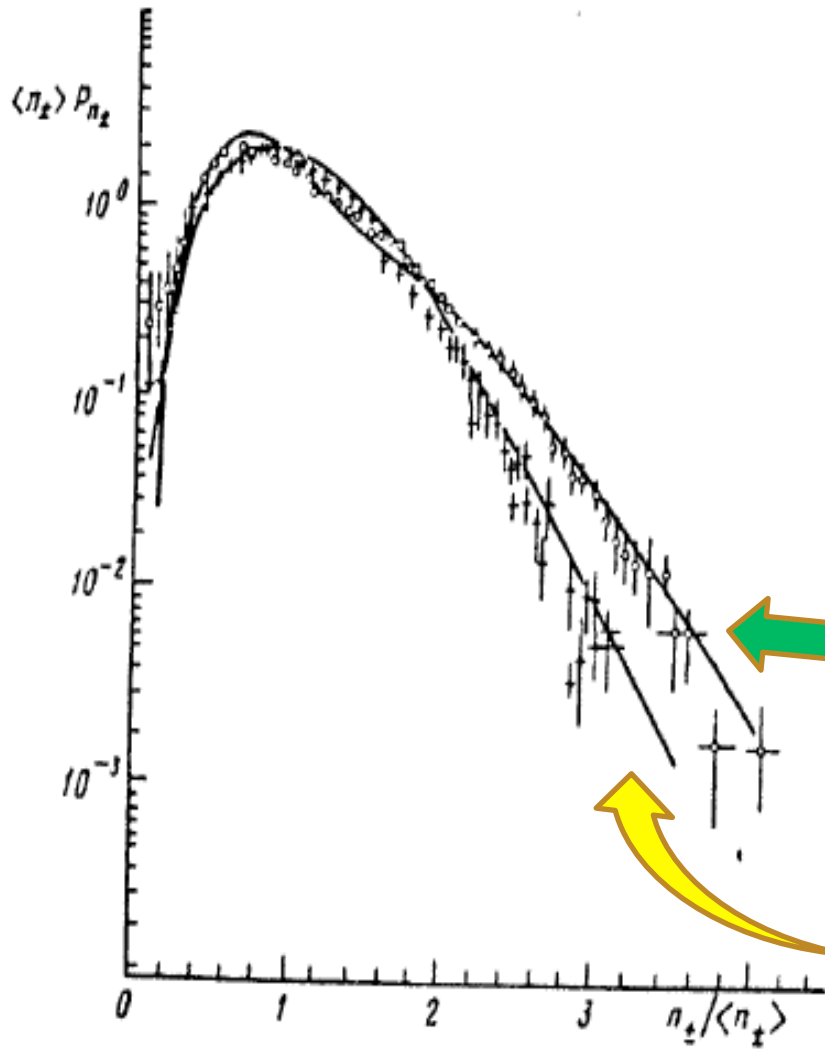
$$z = n / \langle n \rangle$$

$$P_n(s) = \frac{\sigma_n(s)}{\sigma_{tot}(s)} = \frac{1}{\langle n \rangle} \Psi \left(\frac{n}{\langle n \rangle} \right)$$

Experimental data: KNO scaling holds in hh collisions up to $\sqrt{s} = 53$ GeV (ISR)

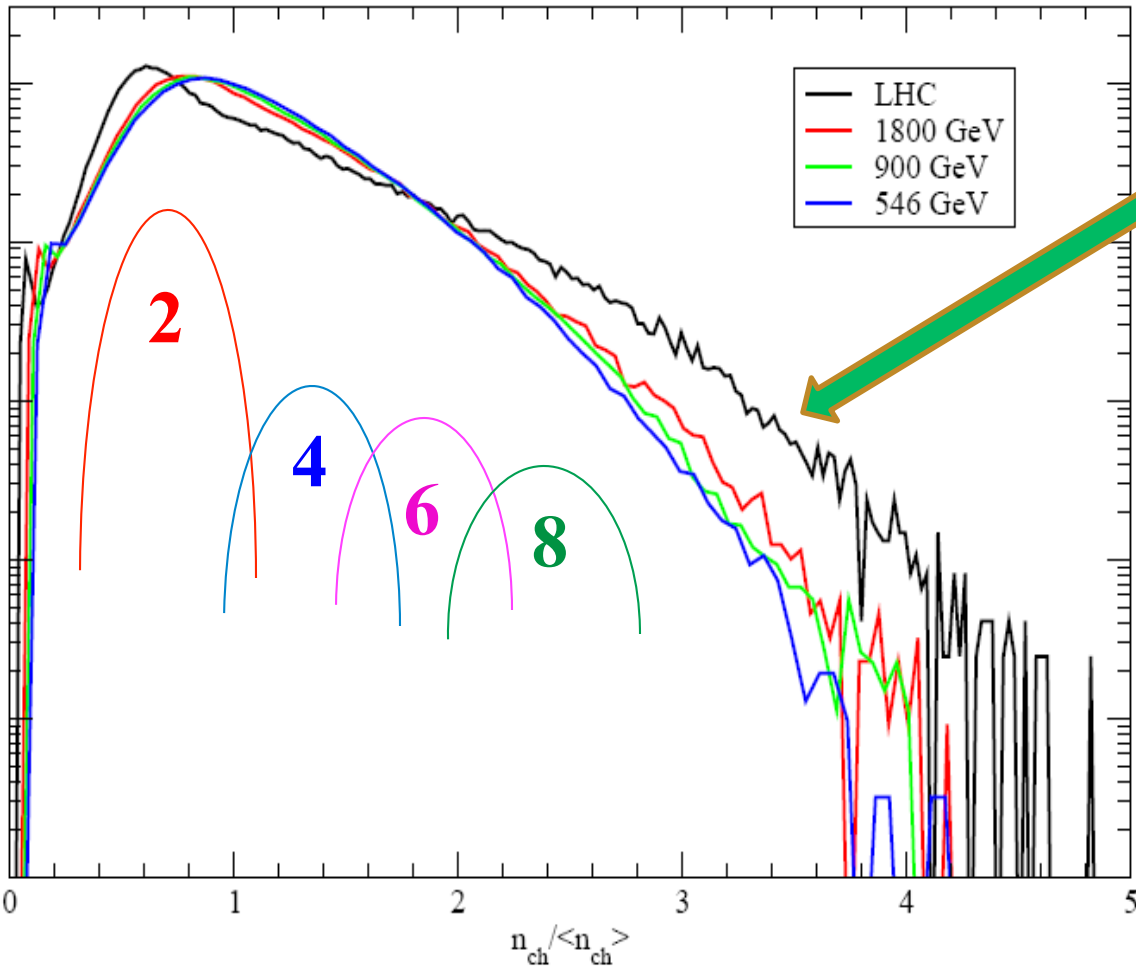
VIOLETION OF KNO SCALING

A.B.Kaidalov, K.A.Ter-Martirosyan, PLB 117 (1982) 247
UA5 Collaboration, Phys. Rep. 154 (1987) 247
N.S.Amelin, L.V.Bravina, Sov.J.Nucl.Phys. 51 (1990) 133



**Charged-particle
multiplicity distributions
in the KNO variables in
nondiffractive
antiproton-proton
collisions at
 $\sqrt{s} = 546 \text{ GeV}$ and
 53 GeV**

VIOATION OF KNO SCALING AT LHC

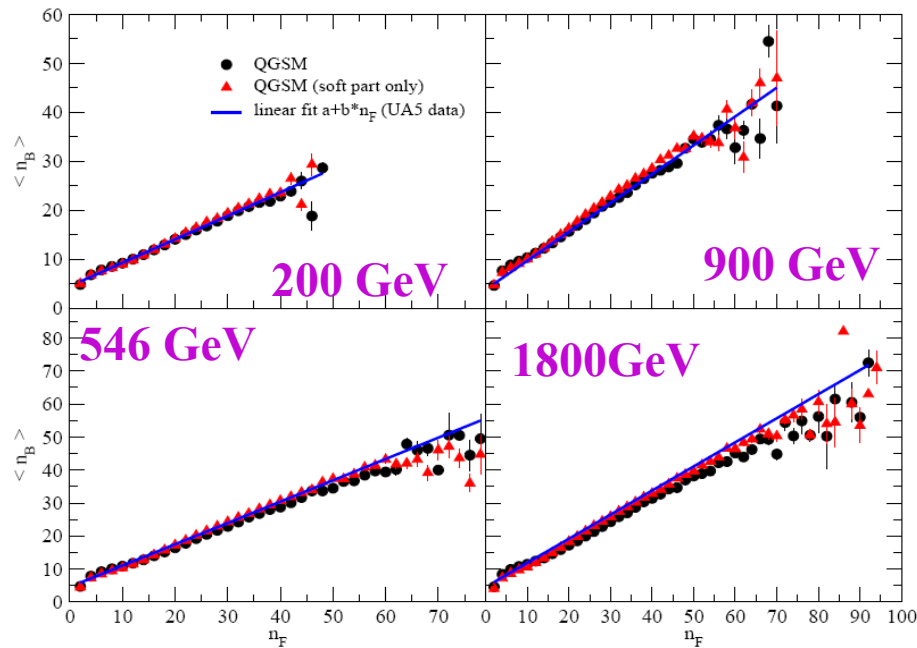


High-multiplicity tail is pushed up, whereas maximum of the distribution is shifted towards small values of z

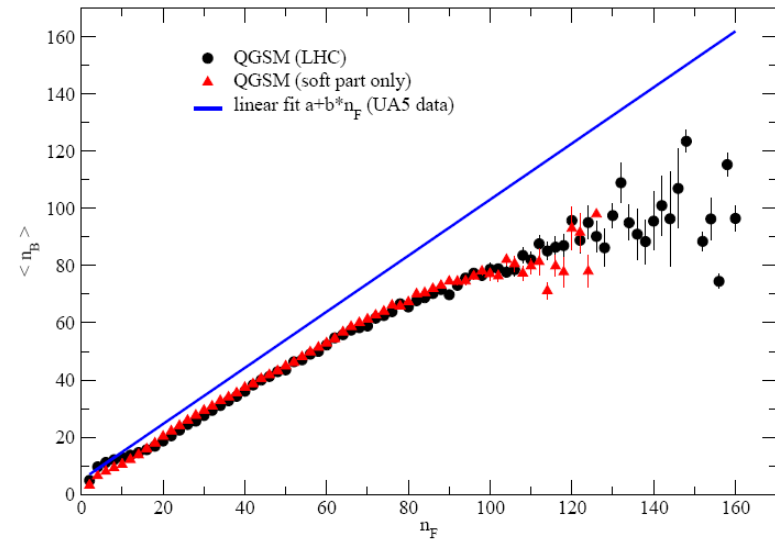
At energies below 100 GeV different contributions overlap strongly, whereas at higher energies – more multi-string processes

\Rightarrow Enhancement of high multiplicities

FORWARD-BACKWARD MULTIPLICITY CORRELATIONS

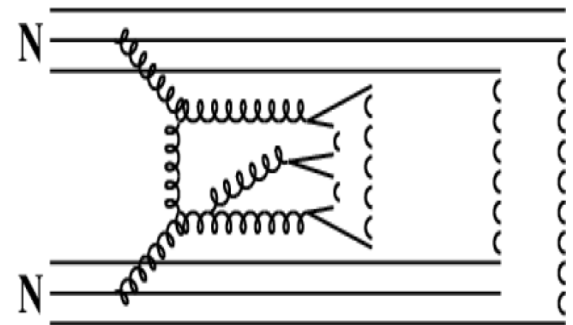


pp 14 TeV

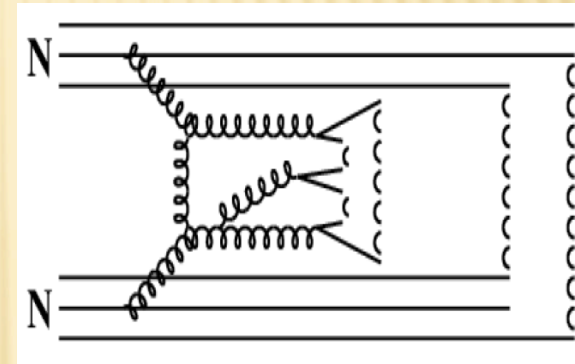
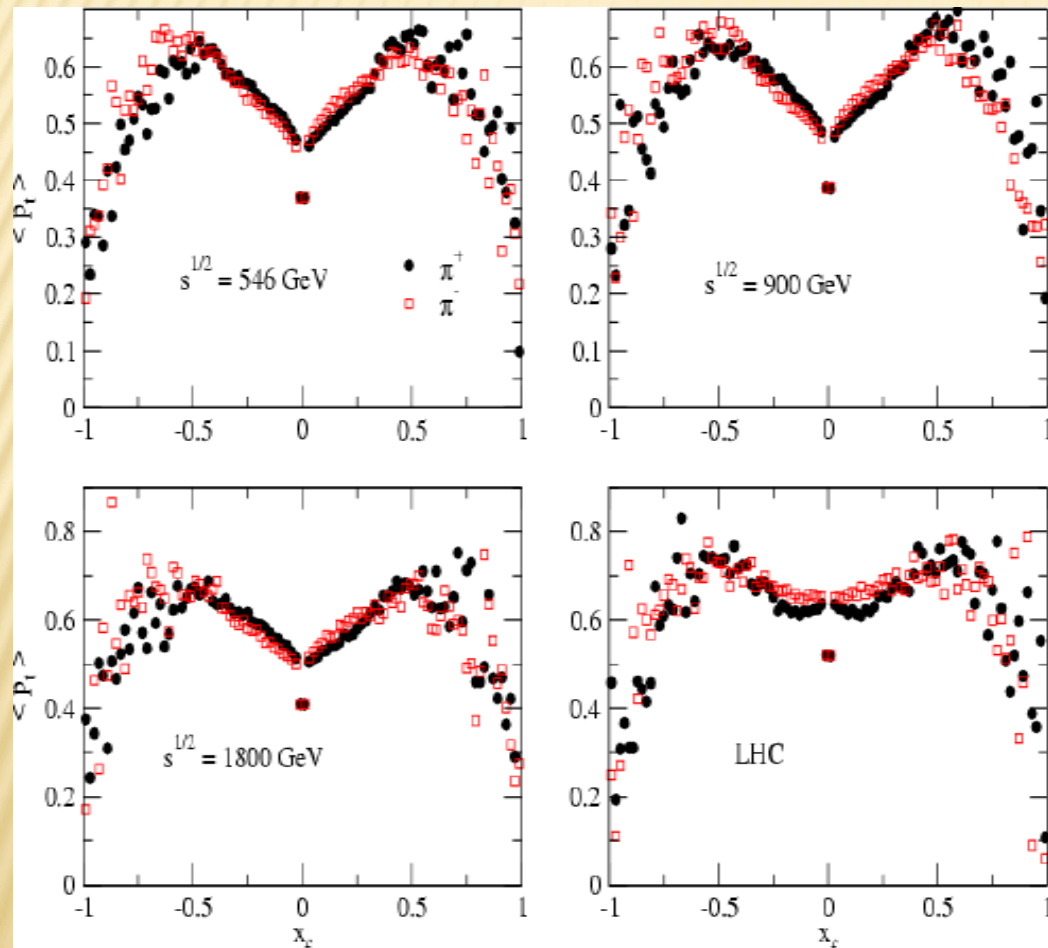


$\langle n_B(n_F) \rangle = a + b n_F$ is linear with increasing of the slope b with energy due to

- 1) Multi-chain diagrams
- 2) Color exchange type of string excitation



STRONG SEA-GULL EFFECT $\langle PT(XF) \rangle$



Sea-gull effect becomes more pronounced with energy

8. Anisotropic flow in pp

Azimuthal anisotropy in relativistic string fragmentation, I

Accepted picture for flow in heavy ion collisions – hydro expansion of QGP. Still, flow in pp and light AA is an open question:

- ? possible reasons for it
- ? magnitude
- ? possibility of observation

All the points are linked with each other

⇒ Importance of models as a test-ground for study of possible mechanisms.

Possibility of flow in DPM

- **DPM**: final particles come as fragments of qg strings, N of strings is defined via **RFT**.
- **RFT** study (K.Boreskov, A.Kaidalov, O.Kancheli) proposes azimuthal anisotropy.
- Model for \mathbb{P} with transverse separation of its ends – qg string → relativistic string with transverse separation of its ends.

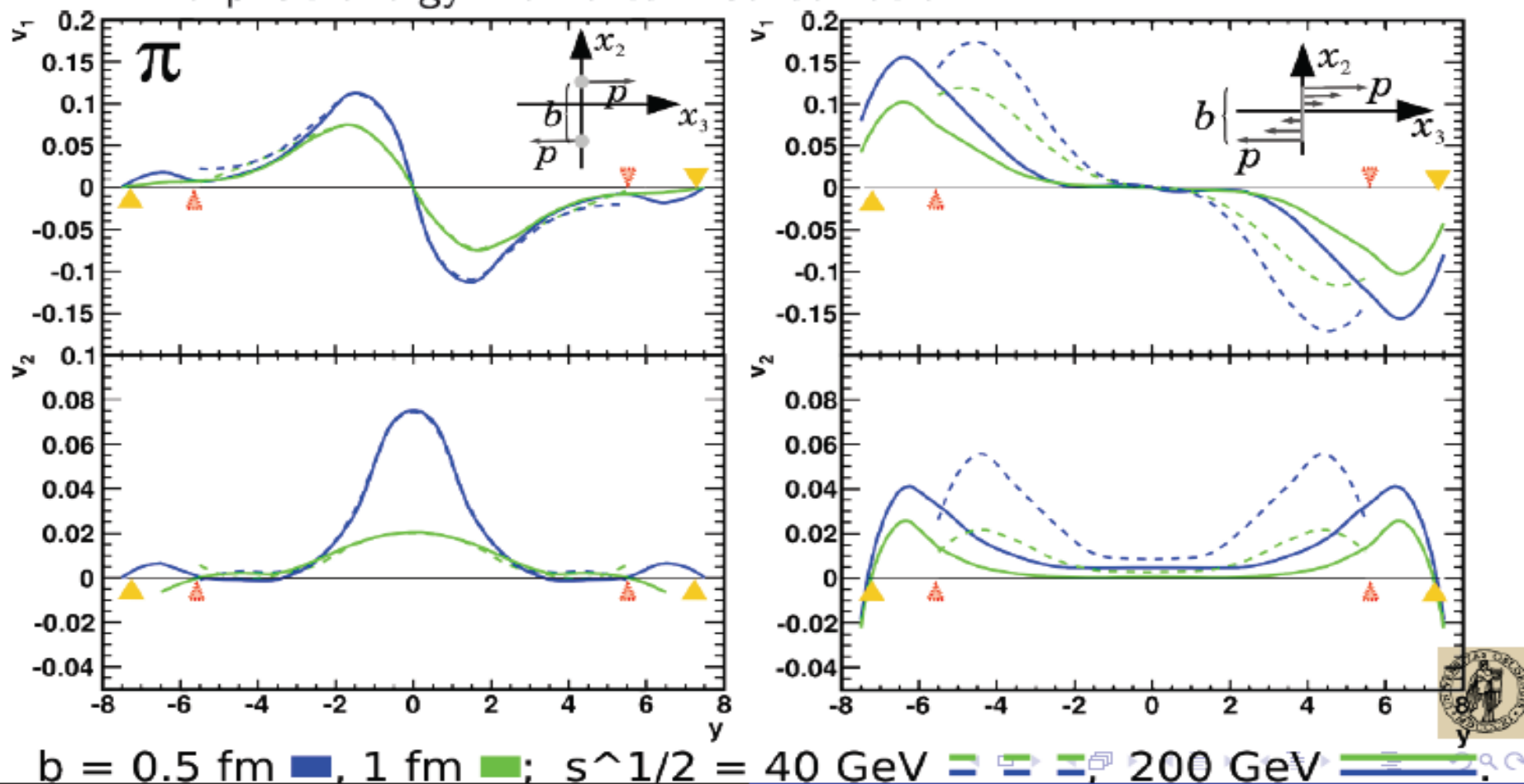


8. Anisotropic flow in pp

Azimuthal anisotropy in relativistic string fragmentation, II

Comparatively simple model, only one sort of particles (“ π -mesons”).

But: explicitly observed string dynamics;
explicit energy-momentum conservation.



8. Anisotropic flow in pp

Azimuthal anisotropy in relativistic string fragmentation, III

RESULTS:

- 1 Both v_1 and v_2 present; positive v_2 , v_1 comes with the same sign as v_1 in AuAu experiment.
- 2 Extreme sensitivity to the internal momentum distribution.

Paper R.Kolevator "On azimuthal anisotropy in fragmentation of classical relativistic string " (arXiv:0912.5377v1 [hep-ph]); submitted to EPJC.

OUTLOOK:

- 1 Application to pp involve $2 \times n$ strings asymmetric in rapidity,
- 2 Need much deeper understanding of string formation within RFT (see p.2 of the results)



Summary and perspectives

- *Feynman scaling should hold in pp collisions at LHC in the fragmentation regions only*
- *=> Extended longitudinal scaling holds there as well*
- *It would be interesting to check the ELS for pp collisions within the statistical thermal model*
- *KNO scaling is strongly violated at LHC. The origin of the violation is traced to multi-string processes*