
Lattice-QCD Studies of the H-Dibaryon

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Strangeness in Quark Matter

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Introduction - The H-Dibaryon

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PHYSICAL REVIEW LETTERS

31 JANUARY 1977

Perhaps a Stable Dihyperon*

R. L. Jaffe†

Stanford Linear Accelerator Center, Stanford University, Stanford, California 94305, and Department of Physics and Laboratory of Nuclear Science, ‡ Massachusetts Institute of Technology, Cambridge, Massachusetts 02139

(Received 1 November 1976)

In the quark bag model, the same gluon-exchange forces which make the proton lighter than the $\Delta(1236)$ bind six quarks to form a stable, flavor-singlet (with strangeness of -2) $J^P = 0^+$ dihyperon (H) at 2150 MeV. Another isosinglet dihyperon (H^*) with $J^P = 1^+$ at 2335 MeV should appear as a bump in $\Lambda\Lambda$ invariant-mass plots. Production and decay systematics of the H are discussed.

- ❖ MIT bag model predicts di-hyperon state (H) with

$$I = 0, S = -2, J^P = 0^+$$

and a mass of $m_H = 2150 \text{ MeV}$

- ❖ H-Dibaryon must decay weakly

Experimental Searches

VOLUME 87, NUMBER 21

PHYSICAL REVIEW LETTERS

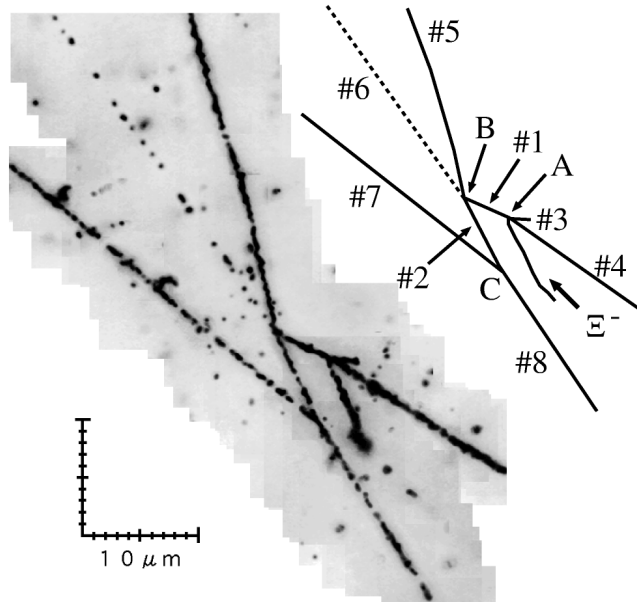
19 NOVEMBER 2001

Observation of a ${}_{\Lambda\Lambda}^6\text{He}$ Double Hypernucleus

(E373@KEK):

A double-hyperfragment event has been found in a hybrid-emulsion experiment. It is identified uniquely as the sequential decay of ${}_{\Lambda\Lambda}^6\text{He}$ emitted from a Ξ^- hyperon nuclear capture at rest. The mass of ${}_{\Lambda\Lambda}^6\text{He}$ and the Λ - Λ interaction energy $\Delta B_{\Lambda\Lambda}$ have been measured for the first time devoid of the ambiguities due to the possibilities of excited states. The value of $\Delta B_{\Lambda\Lambda}$ is $1.01 \pm 0.20_{-0.11}^{+0.18}$ MeV. This demonstrates that the Λ - Λ interaction is weakly attractive.

“Nagara” event



Observation of a ${}_{\Lambda\Lambda}^6\text{He}$ double-hypernucleus

Binding energy:

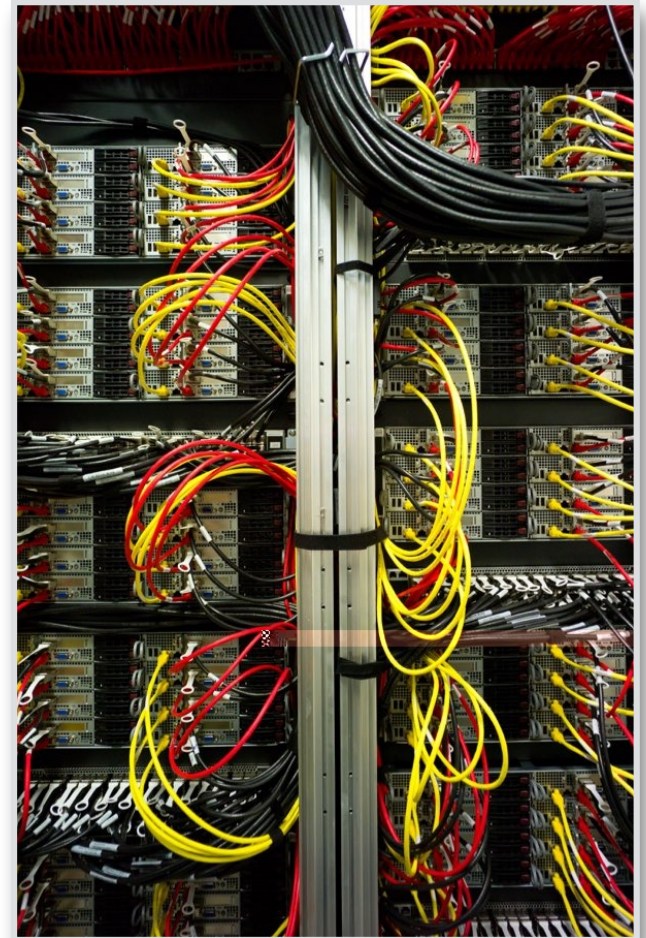
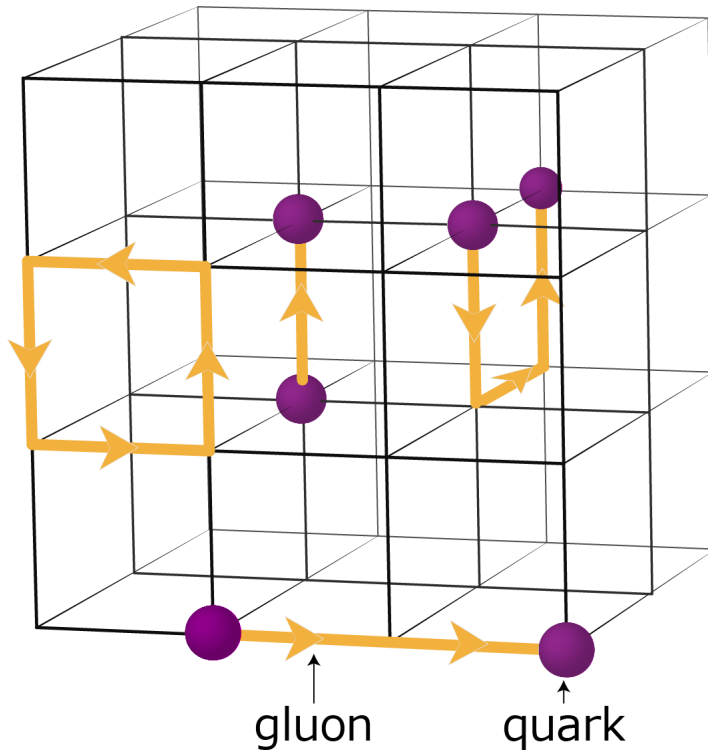
$$B_{\Lambda\Lambda} = 6.91 \pm 0.16 \text{ MeV}$$

Absence of a strong decay ${}_{\Lambda\Lambda}^6\text{He} \longrightarrow {}^4\text{He} + H$

$$\implies m_H > 2m_\Lambda - B_{\Lambda\Lambda}$$

Current Status

- ❖ H-Dibaryon not firmly established
- ❖ Is a bound H-Dibaryon a consequence of QCD?
- ❖ Try “ab initio” technique: Lattice QCD



“Clover” @ HIM

Outline

I. The H-Dibaryon in Lattice QCD

II. Recent lattice calculations

III. Summary & Conclusions

Beyond Perturbation Theory: Lattice QCD

- ❖ Non-perturbative treatment; regularised Euclidean functional integrals

Lattice spacing: $a, \quad x_\mu = n_\mu a, \quad a^{-1} = \Lambda_{UV}$

Finite volume: $L^3 \cdot T$

$$\langle \Omega \rangle = \frac{1}{Z} \int \prod_{x,\mu} dU_\mu(x) \Omega \prod_f \det (\not{D}^{\text{lat}} + m_f) e^{-S_G[U]}$$

- ❖ **Stochastic** evaluation of $\langle \Omega \rangle$ via **Markov process**

Strong growth of numerical cost near physical m_u, m_d

- ❖ Pion mass, i.e. lightest mass in pseudoscalar channel:

$\approx 500 \text{ MeV}$
(2001)



$\approx 130 - 200 \text{ MeV}$
(2015)

Systematic effects

❖ Lattice artefacts:

$$\left\langle \frac{m_N}{f_\pi} \right\rangle^{\text{lat}} = \left\langle \frac{m_N}{f_\pi} \right\rangle^{\text{cont}} + O(a^p), \quad p \geq 1$$

→ extrapolate to continuum limit from $a \approx 0.05 - 0.12$ fm

❖ Finite volume effects

- Empirically: $m_\pi L \geq 4$ sufficient for many purposes
- Could be more severe for multi-baryon systems
- Provide information on **scattering phase shifts**

❖ Unphysical quark masses

- Chiral extrapolation to physical values of m_u, m_d becomes obsolete

❖ Inefficient sampling of SU(3) group manifold

- Simulations become trapped in topological sectors as $a \rightarrow 0$
- Use **open boundary conditions** in time direction *[Lüscher & Schaefer, 2012]*

The H-Dibaryon in Lattice QCD

❖ Spectrum extracted from correlation functions:

$$\sum_{\vec{x}, \vec{y}} \langle O_{\text{had}}(y) O_{\text{had}}^\dagger(x) \rangle = Z_0 e^{-E_0(y_0 - x_0)} + Z_1 e^{-E_1(y_0 - x_0)} + \dots$$

$O_{\text{had}}(x)$: **interpolating operator** for a given hadron;
projects on **all** states with the same quantum numbers

nucleon : $O_N = \epsilon_{abc} (u^a C \gamma_5 d^b) u^c$

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ground state energy

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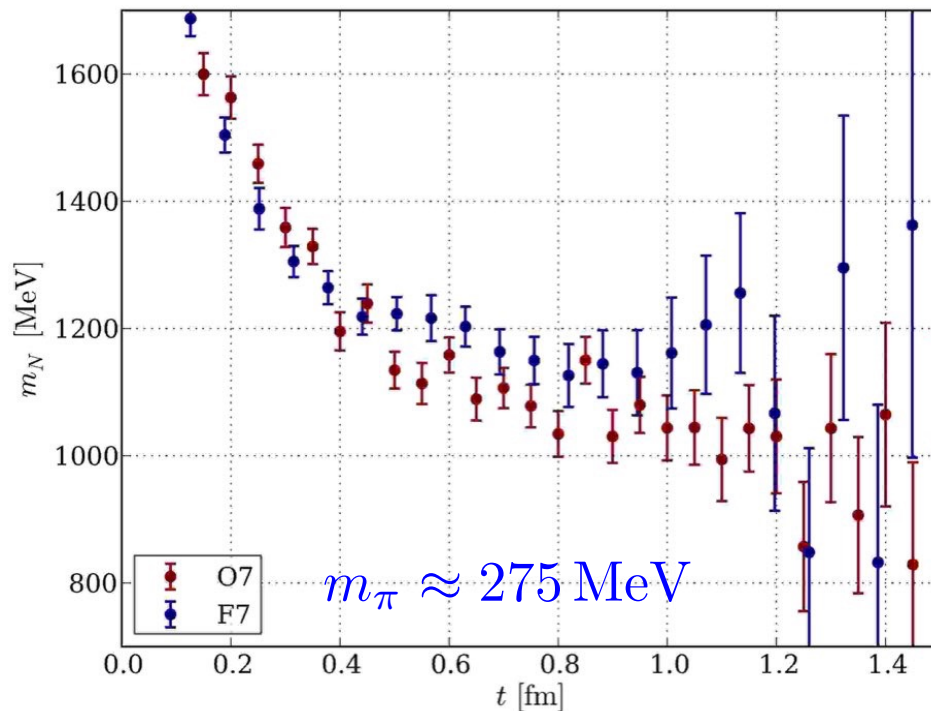
- ❖ Ground state dominates at large Euclidean times: $y_0 - x_0 \rightarrow \infty$
- ❖ Excited states are **sub-leading** contributions

The H-Dibaryon in Lattice QCD

- ❖ Noise problem of baryonic correlation functions:

Exponential growth of noise-to-signal ratio

Nucleon at rest: $R_{NS}(x_0) = e^{(m_N - \frac{3}{2}m_\pi)x_0}$



- ❖ Excited state contributions die out slowly
- ❖ Ground state dominates only for $a \geq 0.5$ fm
- ❖ Precise calculations require very large statistics

[Capitani et al., arXiv:1504.04628]

The H-Dibaryon in Lattice QCD

❖ **Lüscher Method:** Finite-volume effects provide physical information

❖ Two-particle binding momentum:

$$p^2 = \frac{1}{4}(E^2 - \vec{P} \cdot \vec{P}) - m_\Lambda^2 \quad E, m_\Lambda : \text{determined in finite volume}$$

❖ Relation to scattering phase shifts in **infinite volume:**

$$p \cot \delta_0(p) = \frac{2}{\gamma L \sqrt{\pi}} \mathcal{Z}_{0,0}(1, q^2), \quad q = \frac{pL}{2\pi} \quad [\text{Lüscher 1991, Rummukainen \& Gottlieb 1995}]$$

$$\mathcal{Z}_{0,0}(1, q^2) = \frac{1}{\sqrt{4\pi}} \left\{ \sum_{q^2 \neq n^2}^{\Lambda_n} \frac{1}{q^2 - n^2} - 4\pi \Lambda_n \right\}$$

❖ Scattering amplitude: $A \propto \frac{1}{p \cot \delta_0(p) - ip}$

Pole corresponds to binding energy in infinite volume

Recent lattice calculations: Overview

Collaboration	Method	m_π [MeV]	N_f	References
HALQCD	Baryon-baryon potential; Nambu-Bethe-Salpeter wave function	470 — 1170	3	<i>Phys Rev Lett 106 (2011) 162002</i> <i>Nucl Phys A881 (2012) 28</i>
NPLQCD	Two-point correlation functions	806	3	<i>Phys Rev D87 (2013) 034506</i>
	Two-point correlation functions	230, 390	2+1	<i>Phys Rev Lett 106 (2011) 162001</i> <i>Mod Phys Lett A26 (2011) 2587</i>
Mainz	Two-point correlation functions	450, 1000	2	<i>PoS LATTICE2013 (2014) 440</i> <i>PoS LATTICE2014 (2015) 107</i>

❖ **NPLQCD** and **HALQCD** find bound H-dibaryon for $m_\pi \geq 400$ MeV

HALQCD Collaboration

- ❖ Obtain baryon-baryon potential from **Nambu-Bethe-Salpeter** amplitude computed on the lattice

$$G_4(\vec{r}, t - t_0) = \langle 0 | (BB)^{(\alpha)}(\vec{r}, t) (\overline{BB})^{(\alpha)}(t_0) | 0 \rangle = \phi(\vec{r}, t) e^{-2M(t-t_0)}$$

$(BB)^{(\alpha)}(\vec{r}, t)$: 2-baryon interpolating operator; flavour irrep. α

$\phi(\vec{r}, t)$: NBS wave function

M : single baryon mass

- ❖ Determine potential via
$$V(r) = \frac{[-H_0 - (\partial / \partial t)] \phi(\vec{r}, t)}{\phi(\vec{r}, t)}$$

- ❖ Solve the Schrödinger equation to determine binding energies and scattering phase shifts

HALQCD Collaboration

❖ Details of the calculation:

$N_f = 3$ i.e. mass-degenerate u, d, s quarks

Single lattice spacing: $a = 0.121(2)$ fm

5 pion masses in the range: $m_\pi = 469 - 1171$ MeV

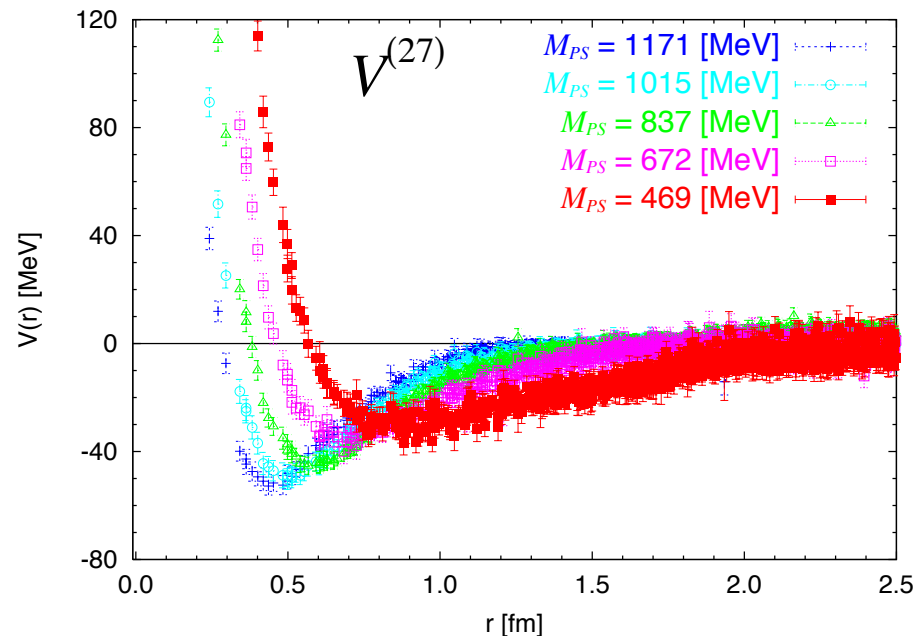
Volumes: $L = 3.87, 2.90, 1.97$ fm

Statistics: $O(500)$ gauge configurations per ensemble

$O(8000)$ “measurements” for each pion mass

HALQCD Collaboration

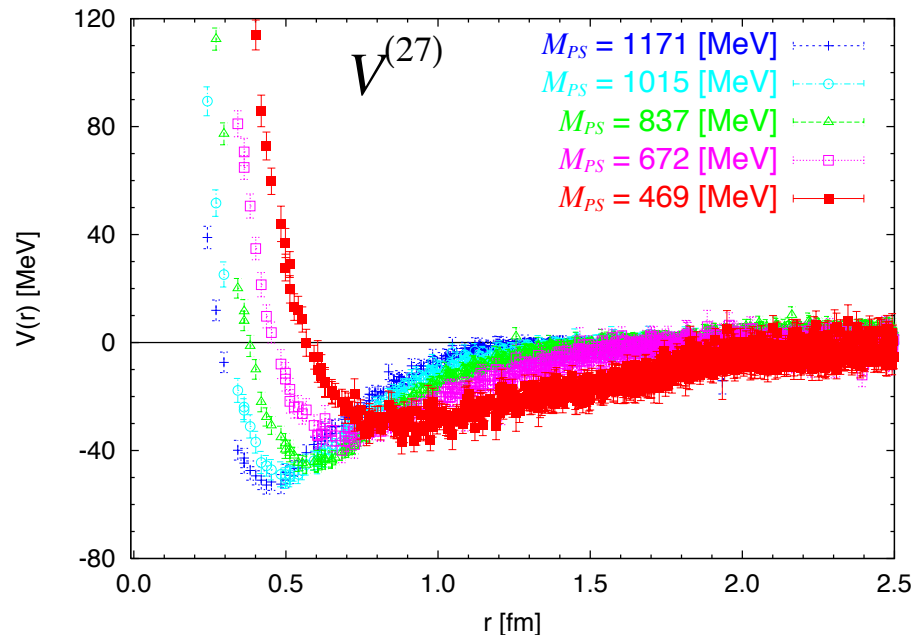
NN potential (27-plet)



- Repulsive core
- Attractive well at intermediate and long distances

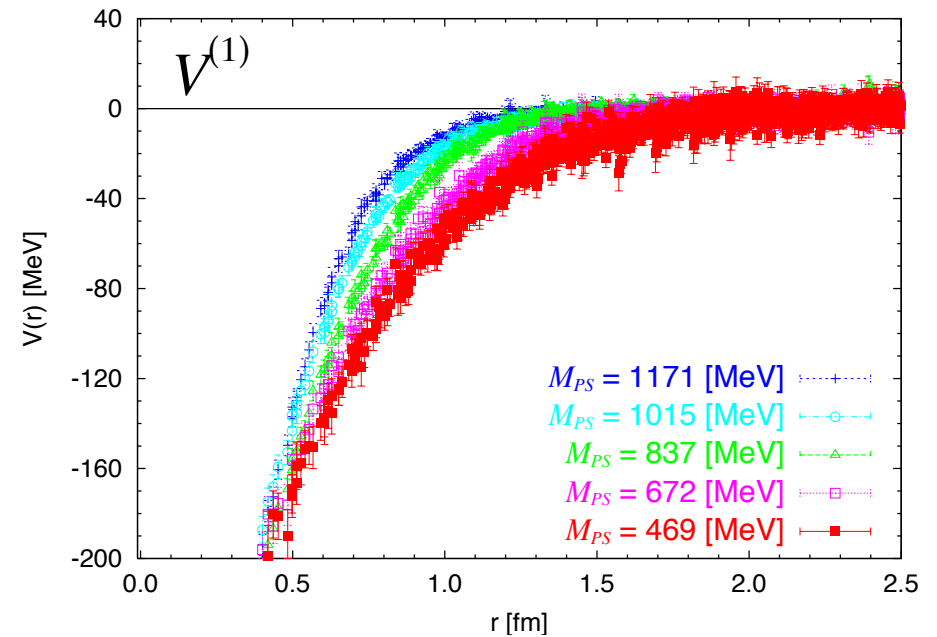
HALQCD Collaboration

NN potential (27-plet)



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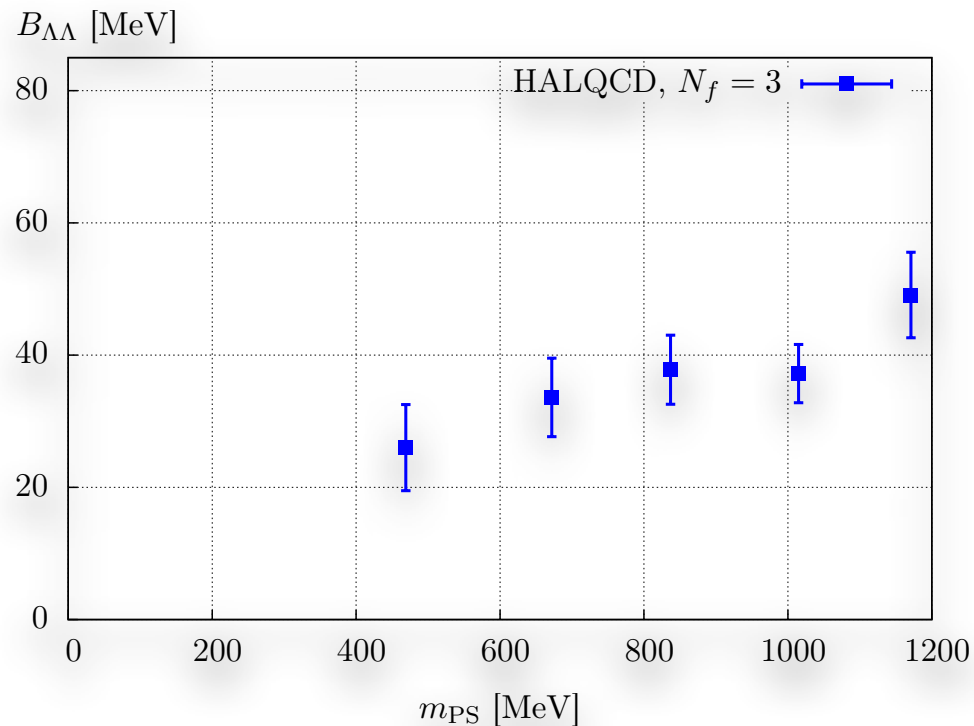
Flavour-singlet potential



- Attractive core
- Mild pion mass dependence
- Parameterise potential as input for Schrödinger equation

HALQCD Collaboration

- ❖ Chiral behaviour of H-dibaryon binding energies in SU(3) limit:



(statistical and systematic errors combined)

- ❖ Stronger attraction at smaller quark mass compensated by larger kinetic energy
- ❖ SU(3) breaking effects not accounted for

Variational method

❖ Consider set of interpolating operators: $O_i, \quad i = 1, \dots, N$

Matrix correlator:
$$C_{ij}(t) = \sum_{\vec{x}} \left\langle O_i(\vec{x}, t + t_0) O_j^\dagger(\vec{x}_0, t_0) \right\rangle$$

❖ Solve Generalised Eigenvalue Problem (GEVP):

$$\mathbf{C}(t + \Delta t)v_n(t) = \lambda_n(t)\mathbf{C}(t)v_n(t) \quad E_{\text{eff}}(t) = -\partial_t \log \lambda_n(t)$$

❖ Operator basis:

- Apply “smearing” to interpolating operators at source and sink
- Use 6-quark and two-baryon operators
- Project onto irreducible representations in flavour space

Variational basis in the dibaryon channel

❖ Six-quark operators:

$$[rstuvw] = \epsilon_{ijk}\epsilon_{lmn} (s^a C \gamma_5 P_+ t^b) (v^l C \gamma_5 P_+ w^m) (r^k C \gamma_5 P_+ u^n)$$

$$H^{(1)} = \frac{1}{48} ([sudsud] - [udusds] - [dudsus])$$

$$H^{(27)} = \frac{1}{48\sqrt{3}} (2[sudsud] + [udusds] - [dudsus])$$

❖ Momentum-projected two-baryon operators:

$$B_\alpha \equiv [rst]_\alpha = \epsilon_{ijk} (s^i C \gamma_5 t^j) r_\alpha^k$$

$$(BB)(\vec{p}_1, \vec{p}_2; t) = \sum_{\vec{x}, \vec{y}} e^{i\vec{p}_1 \cdot \vec{x}} e^{i\vec{p}_2 \cdot \vec{y}} B_1(\vec{x}, t) C \gamma_5 P_+ B_2(\vec{y}, t)$$

Construct $(BB)^{(1)}$, $(BB)^{(8)}$, $(BB)^{(27)}$

NPLQCD Collaboration

- ❖ Compute energy levels in dibaryon channel from two-point correlation functions of multi-baryon operators:

$$O_{\text{H}}^{2\text{baryon}} = \epsilon_{ijk} \{ (r^i C \gamma_5 P_+ s^j) t^k \} (\vec{x}, t) \epsilon_{lmn} \{ (u^l C \gamma_5 P_+ v^m) w^n \} (\vec{y}, t)$$

- ❖ Binding energy: $B_{\Lambda\Lambda} = 2m_{\Lambda} - E_{\Lambda\Lambda}$

- ❖ Details of the calculation: *[Beane et al, Phys Rev D87 (2013) 034506]*

$N_f = 3$ i.e. mass-degenerate u, d, s quarks

Single lattice spacing: $a = 0.145(2)$ fm

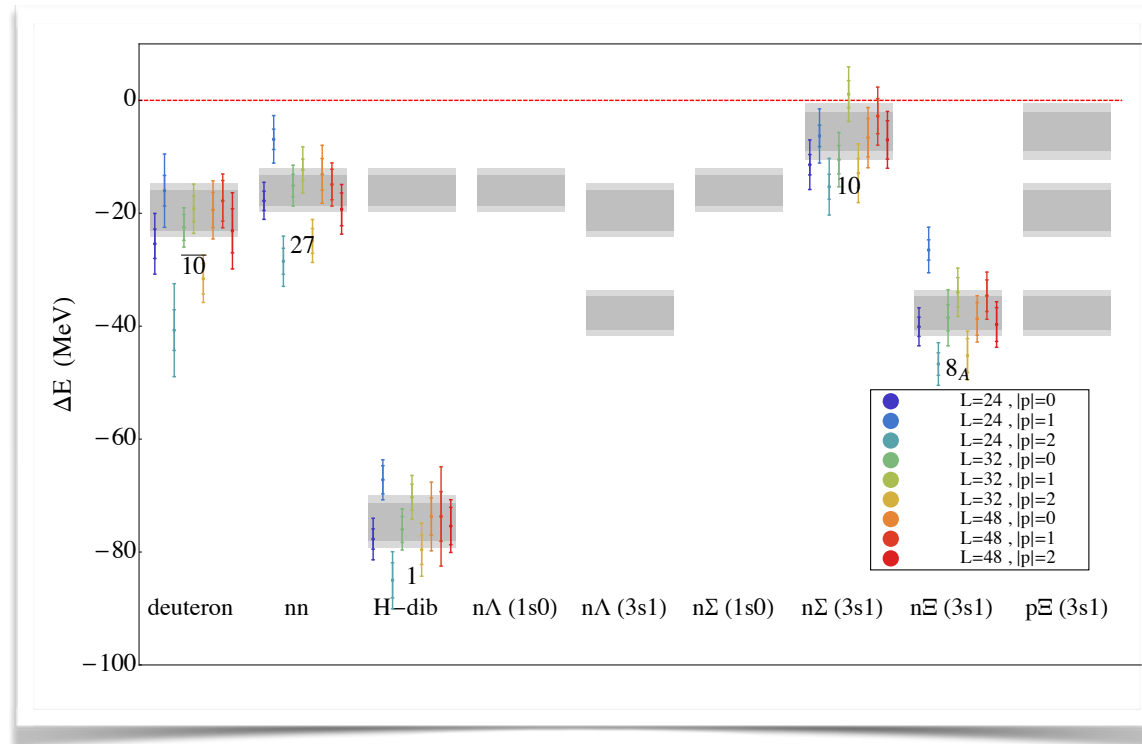
Single pion mass: $m_{\pi} = 807$ MeV

Volumes: $L = 3.4, 4.5, 6.7$ fm

Statistics: $O(10^5)$ “measurements” per ensemble

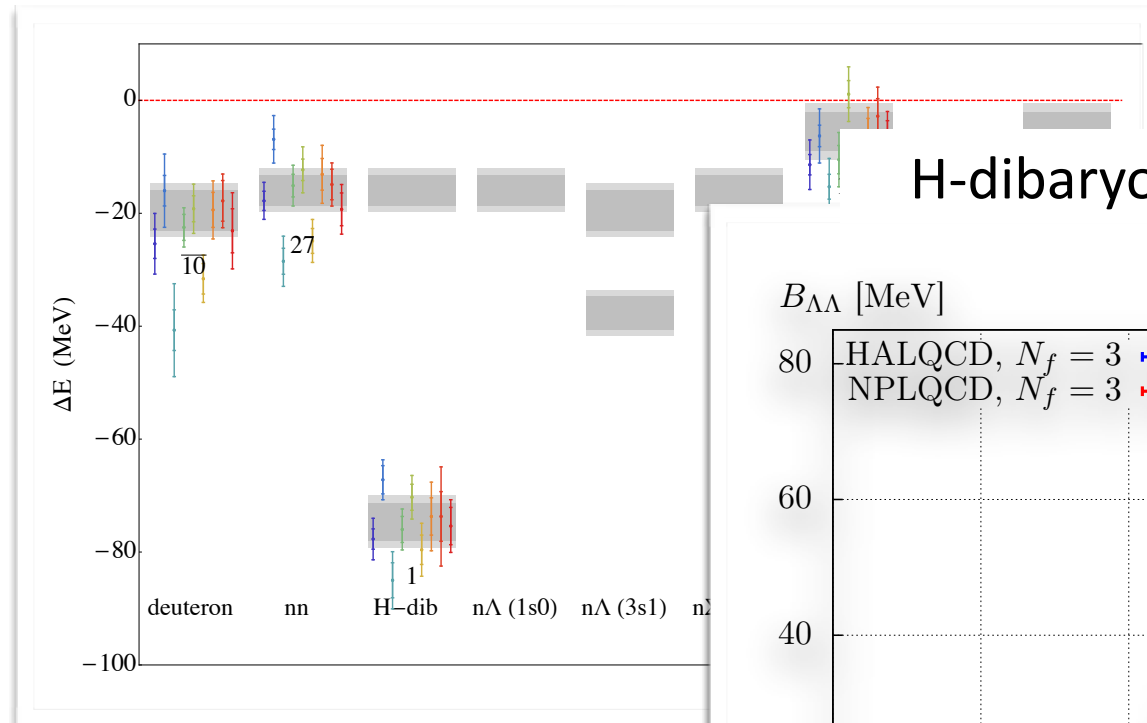
NPLQCD Collaboration

❖ Binding energies of various two-baryon systems:

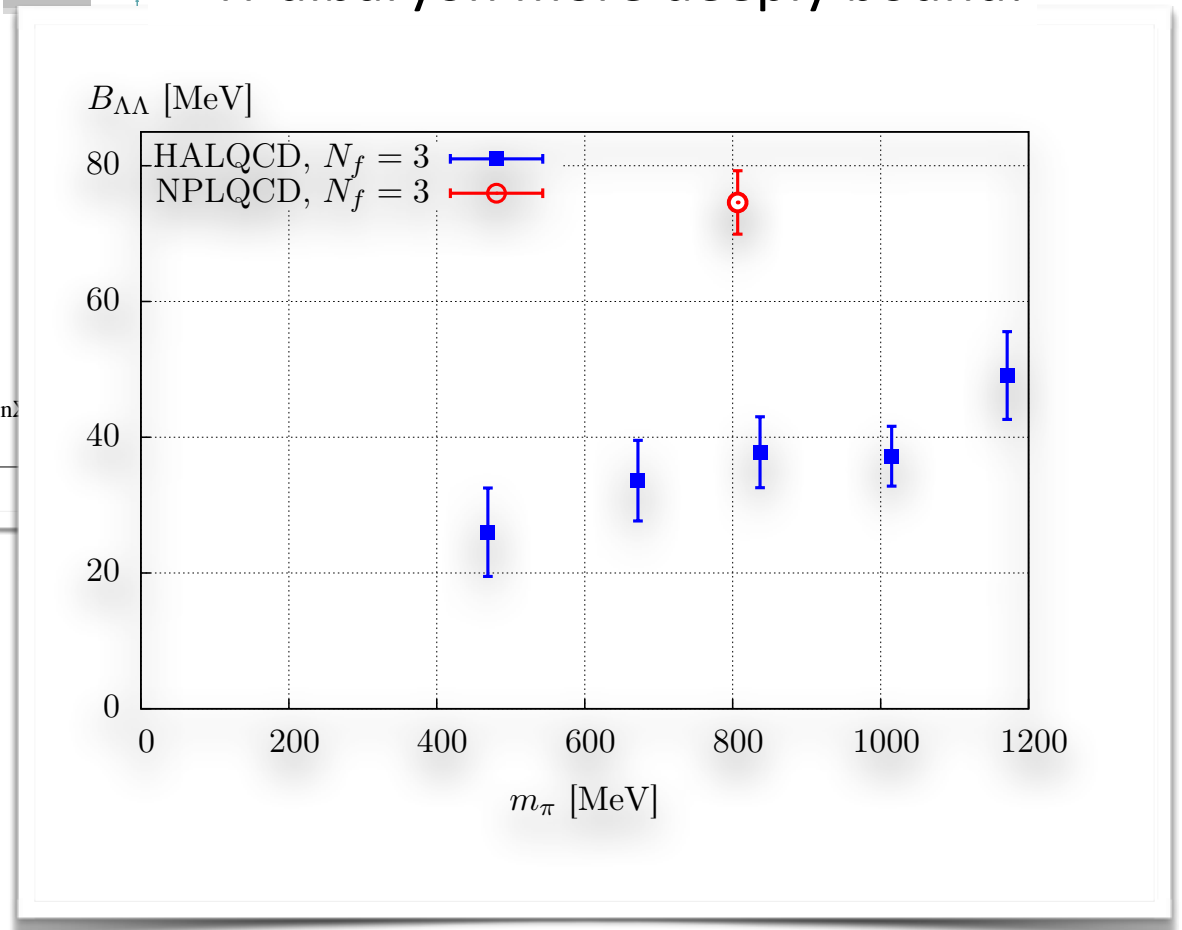


NPLQCD Collaboration

- ❖ Binding energies of various two-baryon systems:



H-dibaryon more deeply bound:



NPLQCD Collaboration

❖ Details of the calculation:

[Beane et al, Phys Rev Lett (2011) 162001,
Mod Phys Lett A26 (2011) 2587]

$N_f = 2 + 1$, anisotropic action: $a_s/a_t = 3.50(3)$

Single lattice spacing: $a_s = 0.123(1)$ fm

Pion masses: $m_\pi = 389, 230$ MeV

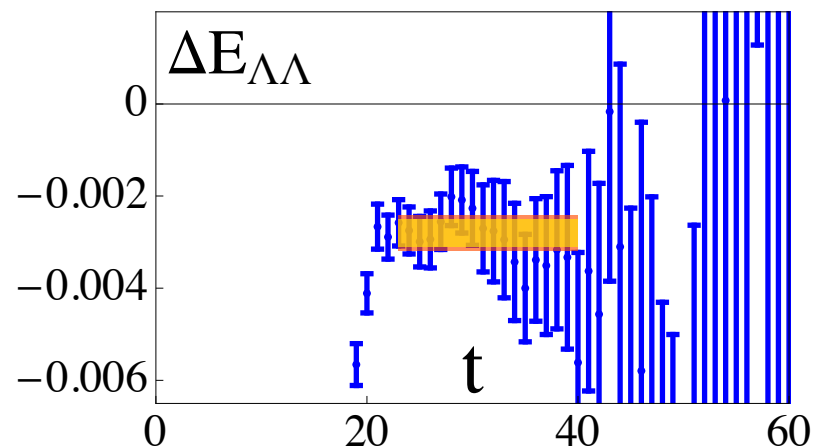
Volumes: $L = 3.9, 3.0, 2.5, 2.0$ fm

Statistics: $O(10^5)$ “measurements” per ensemble

❖ Variational approach:

source: six-quark operators

sink: two-baryon operators



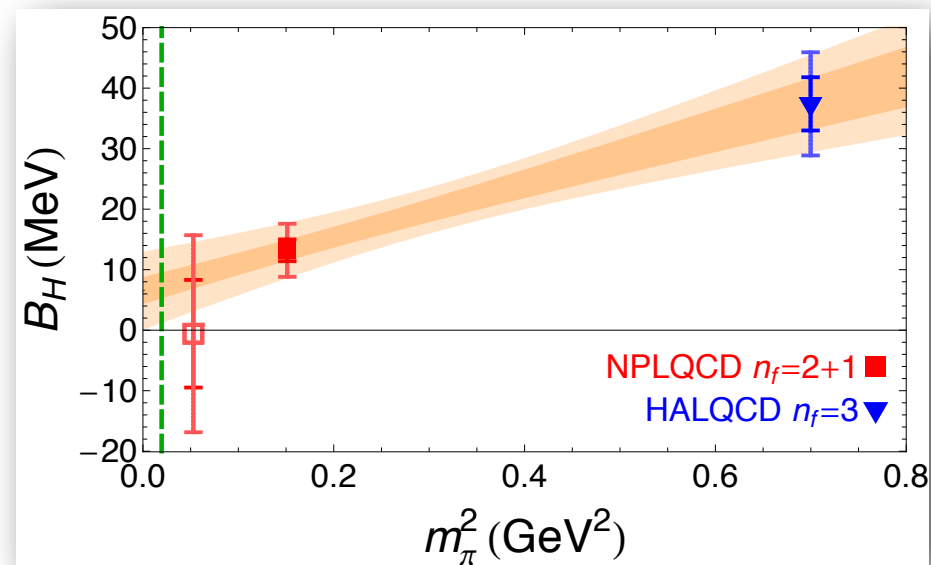
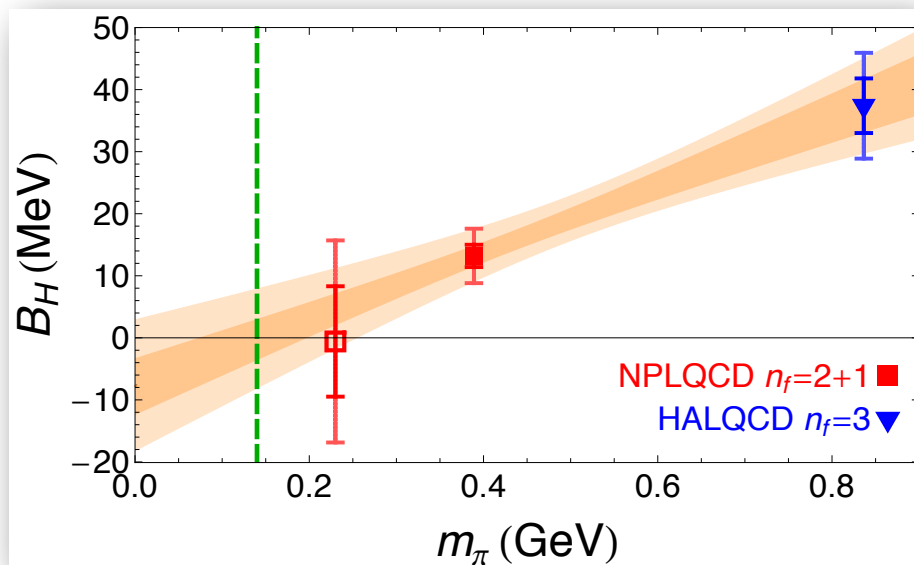
NPLQCD Collaboration

❖ Binding energy in infinite volume limit:

[Beane et al, Phys Rev Lett (2011) 162001,
Mod Phys Lett A26 (2011) 2587]

$$B_{\Lambda\Lambda}^{\infty} = \begin{cases} 13.2(1.8)_{\text{stat}} (4.0)_{\text{syst}} \text{ MeV}, & m_{\pi} = 389 \text{ MeV} \\ 0.6(8.9)_{\text{stat}} (10.3)_{\text{syst}} \text{ MeV}, & m_{\pi} = 230 \text{ MeV} \end{cases}$$

❖ Chiral behaviour:



Mainz / CLS

❖ In collaboration with:

A. Francis, J. Green, P. Junnarkar, Ch. Miao, T. Rae

A. Francis et al., PoS LATTICE2013 (2014) 440, J. Green et al., PoS LATTICE2014 (2015) 107

❖ Details of the calculation:

$N_f = 2$ i.e. quenched strange quark; $O(a)$ improved Wilson fermions

Single lattice spacing: $a = 0.063$ fm

Pion masses: $m_\pi = 457, 1000$ MeV

Single volume: $L = 2.02$ fm

Statistics: 43k, 125k “measurements”

Mainz / CLS

❖ All-mode-averaging

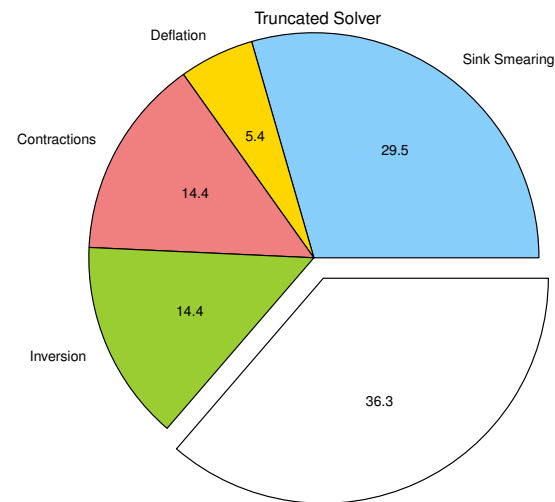
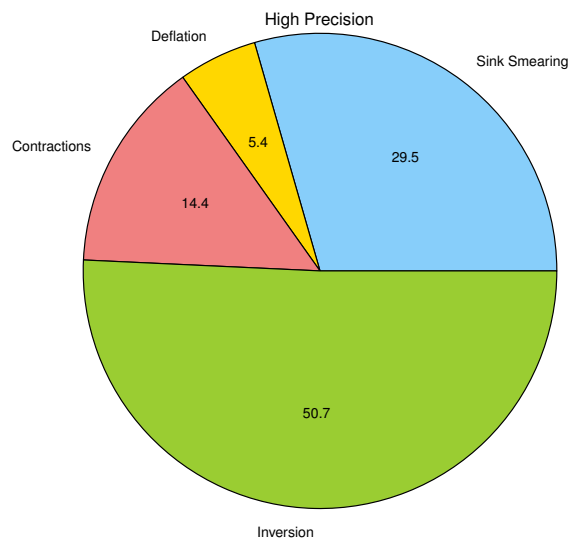
[Blum et al, Phys Rev D88 (2013) 094503]

Combine low-precision solver with bias correction:

$$O = O_{\vec{x}_0} - O_{\vec{x}_0}^{(\text{apprx})} + \frac{1}{N_{\Delta\vec{x}}} \sum_{\Delta\vec{x}} O_{\vec{x}_0 + \Delta\vec{x}}^{(\text{apprx})}$$

Variance reduction:

$$2(1 - r) + \frac{1}{N_{\Delta\vec{x}}}, \quad r = \text{Corr} \left(O_{\vec{x} + \Delta\vec{x}}^{(\text{apprx})}, O_{\vec{x} + \Delta\vec{x}'}^{(\text{apprx})} \right)$$



Mainz / CLS

❖ Ensemble “E1”

$$m_\pi = 1000 \text{ MeV} \approx m_K \quad \text{i.e. } \text{SU}(3)_{\text{flavour}}\text{-symmetric}$$

GEVP set-up:

source: 6-quark operator $H^{(1)}$ with two different smearing levels

sink: 6-quark and multi-baryon operators

❖ Ensemble “E5”

$$m_\pi = 457 \text{ MeV}$$

GEVP set-up:

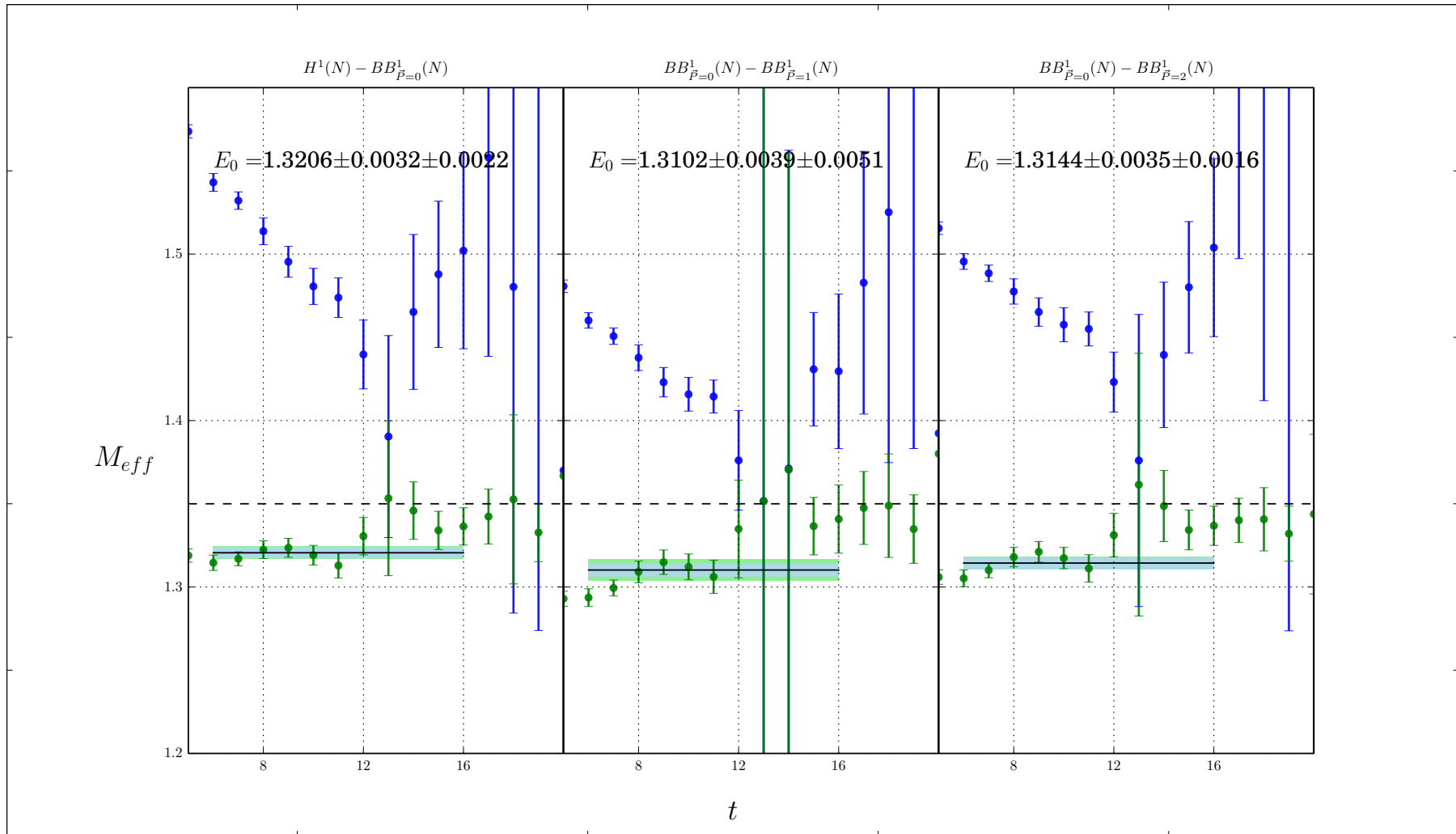
source: 6-quark operators $H^{(1)}$, $H^{(27)}$, different smearing levels

sink: $H^{(1)}$ and $H^{(27)}$ or $(BB)^{(1)}$ and $(BB)^{(27)}$

❖ Consider different total momentum

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❖ Ensemble “E1”

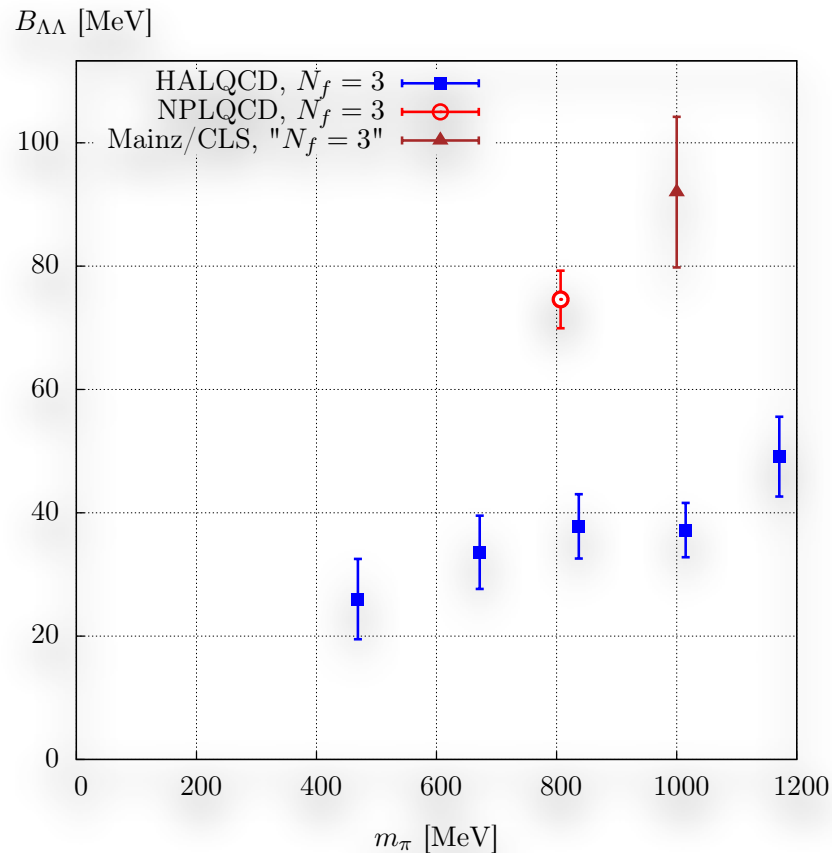


❖ Bound state observed:

$$B_{\Lambda\Lambda} = 92 \pm 10 \pm 7 \text{ MeV}$$

Mainz / CLS

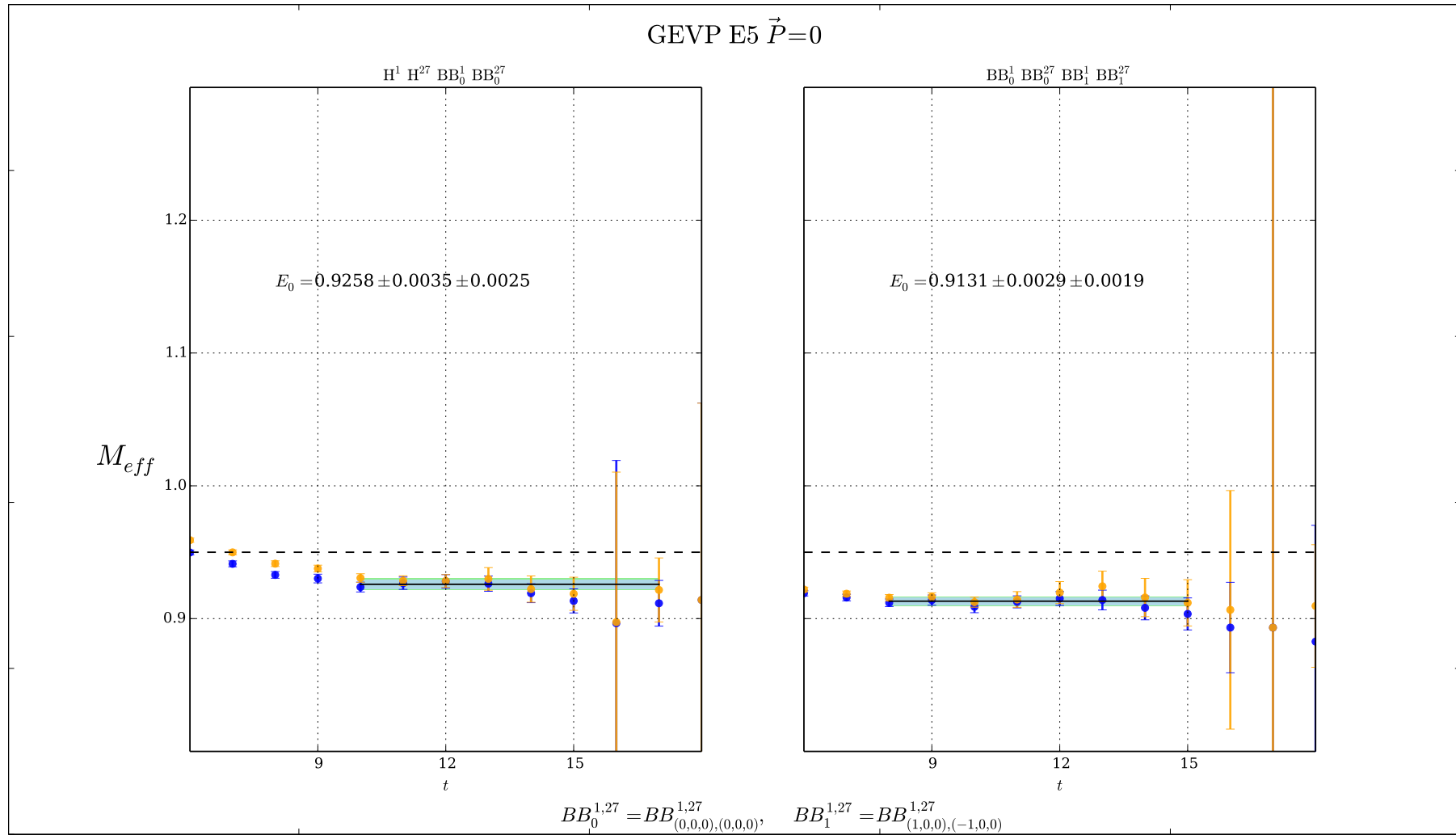
- ❖ Bound state observed: $B_{\Lambda\Lambda} = 92 \pm 10 \pm 7 \text{ MeV}$
- ❖ Multi-baryon operators: better projection onto ground state
- ❖ Chiral behaviour:



- ❖ Caveat: binding energy estimated relative to non-interacting Λ 's

Mainz / CLS

❖ Ensemble “E5”



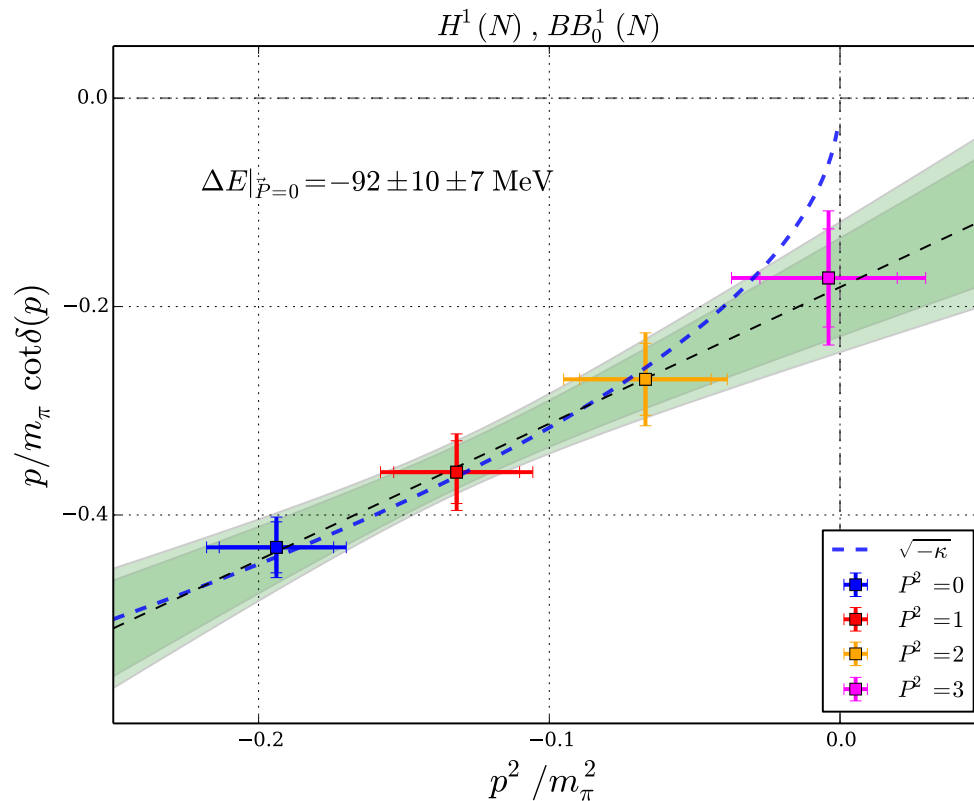
❖ Bound state observed:

$$B_{\Lambda\Lambda} = 77 \pm 11 \pm 7 \text{ MeV}$$

Scattering phase shifts

❖ Ensemble E1:

$$p \cot \delta_0(p) = \frac{2}{\gamma L \sqrt{\pi}} \mathcal{Z}_{0,0}(1, q^2), \quad q = \frac{pL}{2\pi}$$



$$p^2 = \frac{1}{4}(E^2 - \vec{P} \cdot \vec{P}) - m_\Lambda^2$$

$$\mathcal{A} \propto \frac{1}{p \cot \delta_0(p) - ip}$$

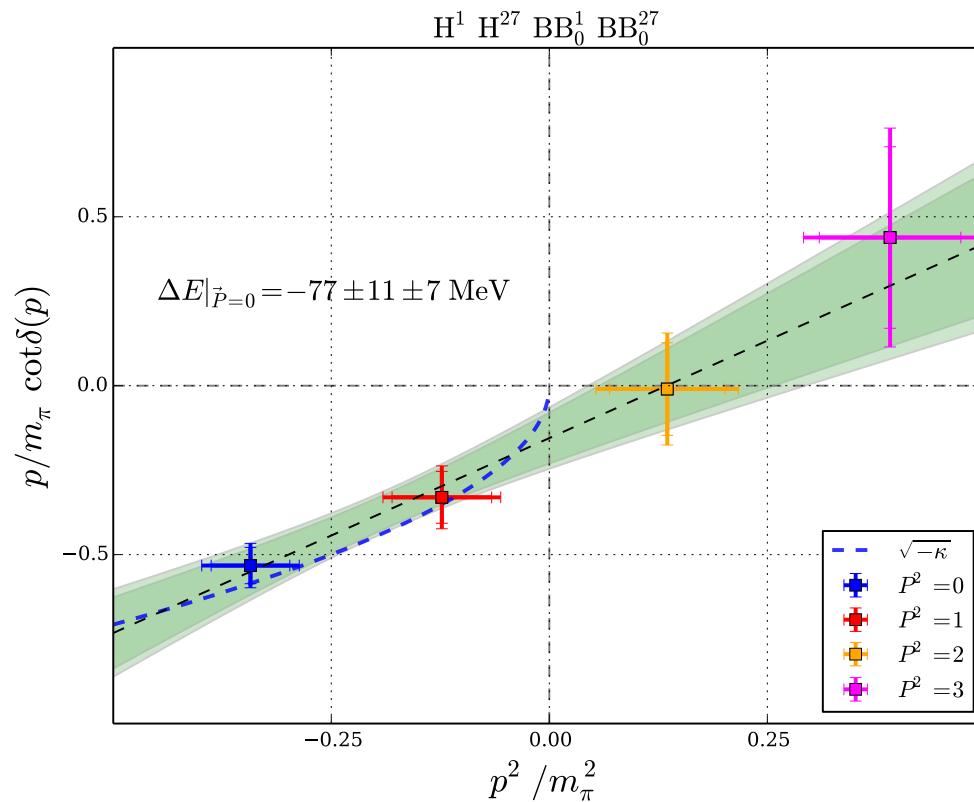
Preliminary!

❖ More data required for reliable determination of the pole

Scattering phase shifts

❖ Ensemble E5:

$$p \cot \delta_0(p) = \frac{2}{\gamma L \sqrt{\pi}} \mathcal{Z}_{0,0}(1, q^2), \quad q = \frac{pL}{2\pi}$$



$$p^2 = \frac{1}{4}(E^2 - \vec{P} \cdot \vec{P}) - m_\Lambda^2$$

$$\mathcal{A} \propto \frac{1}{p \cot \delta_0(p) - ip}$$

Preliminary!

❖ More data required for reliable determination of the pole

Summary and Conclusions

- ❖ Lattice calculations of the H-dibaryon binding energy technically demanding
- ❖ Bound H-dibaryon found for unphysically large pion masses
- ❖ Multi-baryon operators have good overlap onto ground state
- ❖ Chiral behaviour of binding energies to be investigated
- ❖ Need more and better lattice data to make use of finite-volume scattering formalism



Thank you!