# INVESTIGATION OF EXOTIC STATES IN LIGHT NUCLEI 

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## Relevance

## Why ${ }^{13} \mathrm{C}$ ?

Until now, the neutron halo has been observed almost exclusively in the ground states of some radioactive nuclei.
${ }^{13} \mathrm{C}$ nucleus is a good example of a "normal" core, which is well described in the framework of the shell model. The ${ }^{13} \mathrm{C}$ level diagram was reliably determined up to excitation energies of $\sim 10$ MeV.

|  | $\begin{aligned} & 1 / 37 \\ & \text { 的发2-) } \end{aligned}$ |
| :---: | :---: |
| 10460.0 keV |  |
| 9897.0 keV | 3/2- |
| 9499.8 keV | 9/2+ |
| 8860.0 keV | 1/2- |
| 8200.0 keV | 3/2+ |
|  | $3 / 3_{2}^{+}+$ |
| 6864.0 keV | 5/2+ |

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## Search for a neutron halo in the excited state of the ${ }^{13} \mathrm{C}$ nucleus

The halo is most likely to be detected in nuclear states in which the valence neutron occupies an s-orbit, since the absence of a centrifugal barrier can lead to a significant increase in the radius. In the ${ }^{13} \mathrm{C}$ and ${ }^{9} \mathrm{Be}$ nuclei, such states are known near neutron thresholds $\left(1 / 2^{+}, 3.09 \mathrm{MeV}\right.$ and 1.68 MeV , respectively).


## Three types of neutron halo

The neutron is in the s-state

The discrete spectrum


In the continuum

Excited state



## Relevance

## Analogues of Hoyle's state

For nuclei with a pronounced cluster structure, the shell model does not reflect a number of their characteristic features. In the paper Milin M. and von Oertzen W // Eur. Phys. J. A-2002.-Vol. 14, it is assumed that similar Hoyld states can be detected in some neighboring nuclei, for example, the excited state of $8.86 \mathrm{MeV}\left(1 / 2^{-}\right)$in the ${ }^{13} \mathrm{C}$ nucleus.


## Experiment

A series of collaborative experiments were conducted with the scientists of the Kurchatov Institute (Moscow, Russia) and the University of Jyväskylä (Jyvaskyla, Finland) in order to study the "exotic" state in neutron-rich nuclei $\left({ }^{9} \mathrm{Be},{ }^{11} \mathrm{~B}\right.$, and $\left.{ }^{13} \mathrm{C}\right)$,

Experiments were carried out on the cyclotrons U150M (INP, Kazakhstan) and K-130 (UY, Finland) at energies of accelerated atpha particles $E(\alpha)=29$ and 65 MeV , respectively.


## The parameters of optical and double folding potentials of $\alpha+{ }^{13} \mathrm{C}$ system at $26.6-65 \mathrm{MeV}$

| $\begin{gathered} \mathrm{E}, \\ \mathrm{MeV} \end{gathered}$ | Model | $\begin{gathered} \mathrm{V}, \\ \mathrm{MeV} \end{gathered}$ | $\mathrm{f}, \mathrm{fn}$ | $\mathrm{a}_{\mathrm{v}}, \mathrm{fm}$ | $\mathrm{N}_{\mathrm{r}}$ | W, MeV | $r_{w}, \mathrm{fn}$ | $\mathrm{a}_{\mathrm{W}}, \mathrm{fm}$ | $\underset{\mathrm{MeVfim}}{ } \mathrm{~J}_{\mathrm{V}}$ | $\underset{\mathrm{MeVfm}^{3}}{\mathrm{~J}_{\mathrm{w}}}$ | $\mathrm{r}_{\mathrm{c}}$, fm |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 26.6 | OM | 148.32 | 1.112 | 0.736 |  | 12.844 | 1.6 | 0.267 | 321 | 66 | 1.28 |
|  | DF |  |  |  | 0.97 | 18.635 | 1.6 | 0.267 | 399.3 |  | 1.28 |
| 29 | OM | 147.22 | 1.112 | 0.736 |  | 12.844 | 1.6 | 0.731 | 318.5 | 73 | 1.28 |
|  | DF |  |  |  | 0.98 | 12.844 | 1.6 | 0.731 | 402.5 |  | 1.28 |
| 35 | OM | 138.95 | 1.112 | 0.8 |  | 14.125 | 1.6 | 0.15 | 312 | 71 | 1.28 |
|  | DF |  |  |  | 0.99 | 14.125 | 1.6 | 0.15 | 404.6 |  | 1.28 |
| 48.7 | OM | 133 | 1.112 | 0.79 |  | 14.79 | 1.6 | 0.639 | 297.3 | 82 | 1.28 |
|  | DF |  |  |  | 0.97 | 14.79 | 1.6 | 0.639 | 392.6 |  | 1.28 |
| 50 | OM | 131.8 | 1.112 | 0.79 |  | 14.41 | 1.6 | 0.761 | 294.7 | 83 | 1.28 |
|  | DF |  |  |  | 0.97 | 14.41 | 1.6 | 0.761 | 391.6 |  | 1.28 |
| 54.1 | OM | 129.5 | 1.112 | 0.795 |  | 14.21 | 1.6 | 0.8 | 290.5 | 82 | 1.28 |
|  | DF |  |  |  | 0.96 | 14.21 | 1.6 | 0.8 | 386.6 |  | 1.28 |
| 65 | OM | 123 | 1.112 | 0.8 |  | 14.97 | 1.6 | 0.76 | 277 | 86 | 1.28 |
|  | DF |  | $\checkmark$ |  | 0.98 | 14.97 | 7.6 | 0.76 | 390.7 |  | 1.28 |




## Results



## Experimental data at energies 29 and 65 MeV .



Excited states of ${ }^{13} \mathrm{C}$


Excited states of ${ }^{13} \mathrm{C}$ nuclei: $3 / 2,3.68 \mathrm{MeV}$ $5 / 2^{+}, 6.86 \mathrm{MeV}$
$5 / 2,7.5 \mathrm{MeV}$
at $\mathbf{E}=29 \mathrm{MeV}$,
calculated within CC method.
Deformation parameters:

| $\mathbf{E}, \mathbf{M e V}$ | Model | $\boldsymbol{\beta} 5 / \mathbf{2}^{-}$ | $\boldsymbol{\beta}$ 5/2 | $\boldsymbol{\beta}$ 3/2- |
| :---: | :--- | :---: | :---: | :---: |
| $\mathbf{2 9}$ | WS | 0.54 | 0.54 | 0.54 |
|  | DF | 0.51 | 0.51 | 0.51 |




## The measurement of the radii of the excited state



Asymptotic Normalization
Coefficients

Modified
Diffraction
Model


Inelastic
Nuclear
Rainbow
Scattering
Ohkubo S. PRC 70, 041602 (2004)

## Modified Diffraction Model

The modified diffraction model assumes that the root-mean-square radius $\mathrm{R}_{\mathrm{rms}}$ (ex.st) of the excited state can be determined through the difference of the diffraction radii of the excited and ground states:

## $\mathbf{R}_{\text {rms }}($ ex.st. $)=\mathbf{R}_{\text {rms }}($ g.s. $)+\left[R_{d i f}(\right.$ el $)-R_{d i f}($ inel $\left.)\right]$

$\mathbf{R}_{\text {dif }}(\mathbf{e l})$ and $\mathbf{R}_{\text {dif }}($ inel $)$ are the diffraction radii, are determined from the positions of the minima and maxima of the experimental angular distributions of the inelastic and elastic scattering, respectively.

## The neutron halo , the excited state $1 / 2^{+}, 3.09 \mathrm{MeV}$

The first excited state of ${ }^{13} \mathrm{C}\left(3.09 \mathrm{MeV}\left(1 / 2^{+}\right)\right)$is located at 1.86 MeV below the ${ }^{13} \mathrm{C}$ $\rightarrow{ }^{12} \mathrm{C}+\mathrm{n}$ threshold. Measuring of its radius using MDM is of particular importance for the method as a whole, since it allows us to compare the obtained results with the other independent approaches of the INRS and ANC.


According to the ruls of the Blair phases, the angular distributions of the elastic and inelastic with the excitation of the $1 / 2^{+}$scattering state should coincide in phase. However, there is a clear systematic shift of the minima and maxima of the inelastic scattering cross section toward smaller angles in comparison with the elastic scattering curve. This type of shift is an indicator of the increase in the diffraction radius of the state $1 / 2^{+}, 3.09 \mathrm{MeV}$.


Diffraction and root-mean square radii of $\mathbf{3 . 0 9} \mathbf{~ M e V}\left(\mathbf{1} / \mathbf{2}^{+}\right)$excited state of ${ }^{13} \mathrm{C}$ nuclei calculated within MDM

| $\mathrm{E}^{*}, \mathrm{MaB}, \mathrm{I}^{\text {r }}$ | $\mathrm{R}_{\mathrm{dif}} \mathrm{fm}$ | $\mathrm{R}_{\mathrm{rms}}$, fm | $\mathrm{E}_{\mathrm{a}}, \mathrm{MeV}$ |
| :---: | :---: | :---: | :---: |
| 0.00, 1/2 ${ }^{-}$ | $5.31 \pm 0.07$ | 2.31 | 65 |
| 3.09, 1/2+ | $5.75 \pm 0.07$ | $2.73 \pm 0.07$ | 29 |
| 3.09,1/2+ | $5.96 \pm 0.06$ | $2.92 \pm 0.07$ | 65 |

Austern N. and Blair J.S. Calculation of inelastic scattering in terms of elastic scattering // Annals of Physics. - 1965. Vol.33. - P.15-64.

## Analogues of Hoyle's state in ${ }^{13} \mathrm{C}$

Observation of the exotic structure and the anomalously large radius of the Hoyle $0^{+}$ (7.65 MeV) state in the ${ }^{12} \mathrm{C}$ nucleus prompted a number of assumptions that a similar situation can occur in neighboring ${ }^{13} \mathrm{C}$ and ${ }^{11} \mathrm{~B}$ nuclei, which differ from ${ }^{12} \mathrm{C}$, by adding a neutron or by removing the proton, respectively.



For a state of $8.86 \mathrm{MeV}\left(1 / 2^{-}\right)$, the minimum of the rainbow (Airy) in comparison with the Hoyle state of 7.65 $\mathrm{MeV}\left(0^{+}{ }_{2}\right)$ is located at a smaller angle. The observed shifts in the positions of airy minima from large angles in inelastic scattering, with respect to elastic scattering, indicate an increase of the radius in this excited state.


The rms radii of the excited state of the $8.86 \mathrm{MeV}\left(1 / 2^{-}\right){ }^{13} \mathrm{C}$ nucleus obtained in the framework of MDM in comparison with the $0^{+}$state of the ${ }^{12} \mathrm{C}$ nucleus

| $\mathrm{E}^{*}$, MəB, $\mathrm{I}^{\boldsymbol{T}}$ | $\mathbf{R}_{\text {dif }}, \mathrm{fm}$ | $\mathbf{R}_{\text {rms }}, \mathrm{fm}$ | $\mathrm{E}_{\mathrm{a}}$, MeV |
| :---: | :---: | :---: | :---: |
| $\mathbf{7 . 6 5}, \mathbf{0}^{+}\left({ }^{12} \mathbf{C}\right)$ | $5.71 \pm 0.04$ | $\mathbf{2 . 8 9} \pm \mathbf{0 . 0 4}$ |  |
| $8.86, \mathbf{1 / 2}^{-}$ | $5.64 \pm 0.09$ | $\mathbf{2 . 6 7} \pm \mathbf{0 . 0 6}$ | 29 |
| $8.86, \mathbf{1 / 2}^{-}$ | $5.66 \pm 0.10$ | $\mathbf{2 . 6 8} \pm \mathbf{0 . 1 2}$ | 65 |

Danilov A.N. et. al. PRC 80, 054603 (2009).

## "Supercompact" state

The ${ }^{13} \mathrm{C}$ nucleus seems to be unique in the sense that several different structures coexist in its spectrum. In addition to the usual levels of the shell model, there are two "diluted" states of various types: one of them contains a neutron halo ( 3.09 MeV ), and the other is an analog of the Hoyle state $(8.86 \mathrm{MeV})$.


It can not be ruled out that there may exist structures, even more exotic, "supercompact". For example, the rms radius of the excited state is $3 / 2^{-}(9.90 \mathbf{M e V})$, obtained from an inelastic scattering analysis within the MDM framework is less than the ground-state radius ( 2.0 fm and 2.3 fm , respectively).


| $\mathbf{E}^{*}, \mathbf{M a B}, \mathrm{I}^{\pi}$ | $\mathrm{R}_{\text {dif }} \mathrm{fm}$ | $\mathrm{R}_{\mathrm{rms}}, \mathrm{fm}$ | $\mathrm{E}_{\boldsymbol{\alpha}}, \mathrm{MeV}$ |
| :---: | :---: | :---: | :---: |
| $\mathbf{0 . 0 0}, \mathbf{1 / 2 ^ { - }}$ | $5.31 \pm 0.07$ | $\mathbf{2 . 3 1}$ | 65 |
| $\mathbf{9 . 9 , 3 / 2 ^ { - }}$ | $4.99 \pm 0.07$ | $\mathbf{1 . 9 7} \pm \mathbf{0 . 0 7}$ | 29 |
| $\mathbf{9 . 9 , 3 / 2 ^ { - }}$ | $5.00 \pm 0.12$ | $2.02 \pm \mathbf{0 . 1 4}$ | 65 |

## COMPARISON OF RECEIVED RESULTS WITH MDM, WITH OTHER METHODS OF DETERMINATION OF RADIUS OF EXCITED STATES OF NUCLEI



Ohkubo S. and Nirabayashi. PRC 70, $041602(\mathrm{R})$ (2004). Liu Z.H. et al. PRC 64, 034312(2001).

## Conclusion

$>$ The radii of the "exotic" excited levels $\left(1 / 2^{+}, 3.09 \mathrm{MeV}\right.$ and $\left.1 / 2^{-}, 8.86 \mathrm{MeV}\right)$ of the ${ }^{13} \mathrm{C}$ nucleus at 29 and 65 MeV are calculated.
$>$ A "supercompact" excited state is found at ${ }^{13} \mathrm{C}$ nuclei

Thus, the uniqueness of the ${ }^{13} \mathbf{C}$ nucleus is shown. The several different structures coexist in its spectrum.

## Thank you for attention!

